

# **Topic 3b,c**

## **Electron Microscopy**

### **1.0 Introduction and History**

- **1.1 Characteristic Information**

### **2.0 Basic Principles**

- **2.1 The Microscope Column**
- **2.2 Signal Detection and Display**
- **2.3 Operating Parameters**

### **3.0 Instrumentation**

- **3.1 Sample Prep**
- **3.2 Handling**

### **4.0 Examples**

### **5.0 Correct Presentation of Results**

# **1.0 Introduction and History**

## **What are they and Where did Electron Microscopes Come From?**

- o Electron Microscopes are scientific instruments that use a beam of highly energetic electrons to examine objects on a very fine scale.**
- o Electron Microscopes were developed due to the limitations of Light Microscopes which are limited by the physics of light to 500x or 1000x magnification and a resolution of 0.2 mm.**

# **1.0 Introduction and History**

## **What are they and Where did Electron Microscopes Come From?**

- o In the early 1930's this theoretical limit had been reached and there was a scientific desire to see the fine details of the interior structures of organic cells (nucleus, mitochondria...etc.).**
- o This required 10,000x plus magnification which was just not possible using current optical microscopes.**

# Dates

- **The Transmission Electron Microscope (TEM) was the first type of Electron Microscope to be developed and is patterned exactly on the Light Transmission Microscope except that a focused beam of electrons is used instead of light to "see through" the specimen. It was developed by Max Knoll and Ernst Ruska in Germany in 1931.**
- **The first Scanning Electron Microscope (SEM) debuted in 1938 ( Von Ardenne) with the first commercial instruments around 1965. Its late development was due to the electronics involved in "scanning" the beam of electrons across the sample.**

# Chronology

**1897: J.J. Thompson** - Discovers the electron

**1924: Louis deBroglie**

Identifies a wavelength to moving electrons -  $\lambda = h / mv$

where:  $\lambda$  = wavelength,  $h$  = Planck's constant  $m$  = mass,

$v$  = velocity (For an electron at 60kV,  $\lambda = 0.005$  nm)

**1926: H. Busch** - Magnetic or electric fields act as lenses for electrons

**1929: E. Ruska** - Ph.D thesis on magnetic lenses

**1931: Knoll & Ruska** - First electron microscope built

**1931: Davisson & Calbrick** - Properties of electrostatic lenses

**1934: Driest & Muller** - Surpass resolution of the OM

**1938: von Borries & Ruska**

First practical EM (Siemens) - 10 nm resolution

**1940: RCA** - Commercial EM with 2.4 nm resolution

# **Microscopy and Related Techniques**

- **Light (optical) microscopy (LM) or (OM)**
- **Scanning electron microscopy (SEM)**  
**Energy dispersive X-ray spectroscopy (EDS)**  
**& Wavelength dispersive X-ray spectroscopy (WDS)**
- **X-ray diffraction (XRD)/X-ray fluorescence (XRF)**
- **Transmission electron microscopy/ Scanning transmission electron microscopy (TEM)/(STEM)**

## **Surface Characterization Techniques**

- **Scanning probe microscopy (AFM & STM)**

# 1.1 Characteristic Information

- **Topography**

**The surface features of an object or "how it looks", its texture; direct relation between these features and materials properties**

- **Morphology**

**The shape and size of the particles making up the object; direct relation between these structures and materials properties**

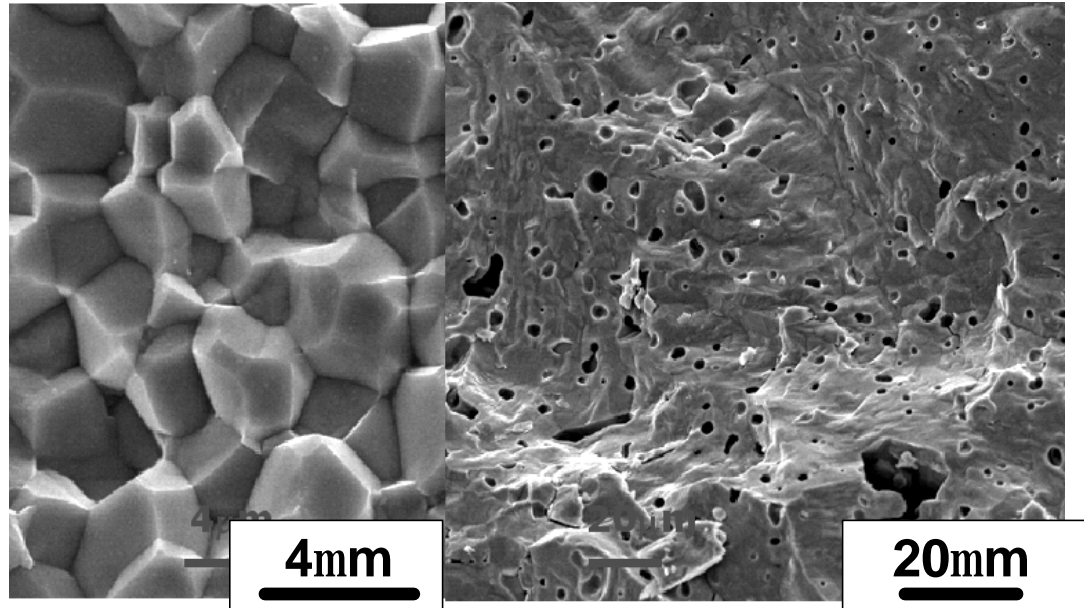
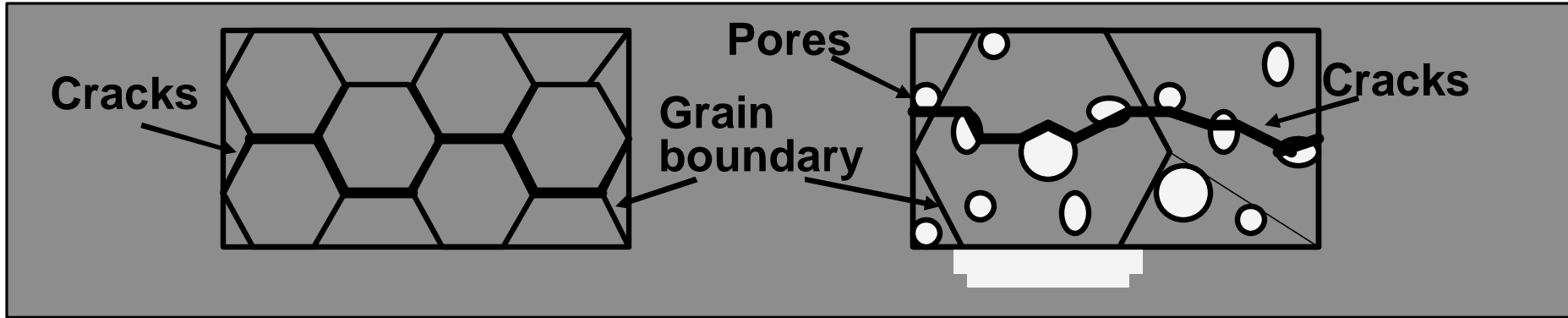
- **Composition**

**The elements and compounds that the object is composed of and the relative amounts of them; direct relationship between composition and materials properties**

- **Crystallographic Information**

**How the atoms are arranged in the object; direct relation between these arrangements and materials properties**

# Identification of Fracture Mode



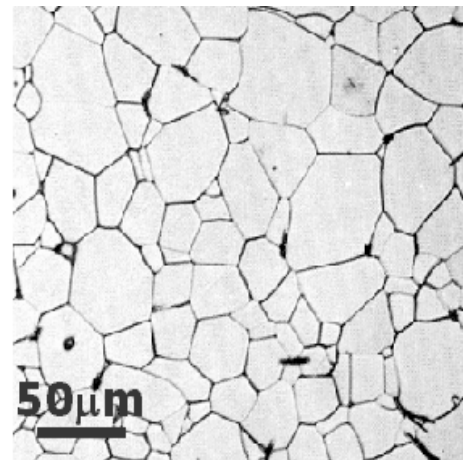
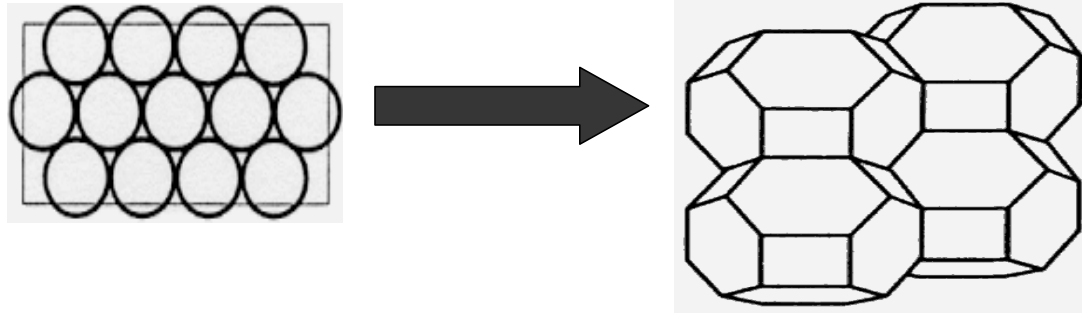
Intergranular fracture

Intragranular fracture

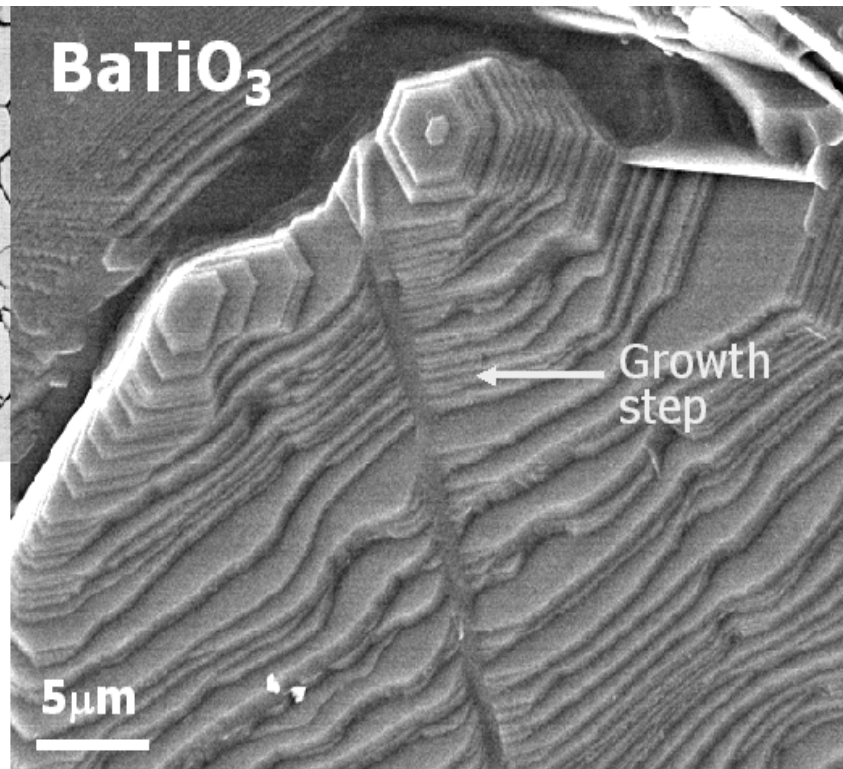
**SEM micrographs of fractured surface of two BaTiO<sub>3</sub> samples.**



# OM and SEM



**OM - 2D**



**SEM - 3D**

# How Fine can You See?

- **Can you see a sugar cube? The thickness of a sewing needle? The thickness of a piece of paper? ...**
- **The resolution of human eyes is of the order of 0.1 mm, 100mm » 4 mils.**
- **However, something vital to human beings are of sizes smaller than 0.1mm, e.g. our cells, bacteria, microstructural details of materials, etc.**

# **Microstructural Features which Concern Us**

- **Grain size: from  $\mu\text{m}$  to the cm regime**
- **Grain shapes**
- **Precipitate size: mostly in the  $\mu\text{m}$  regime**
- **Volume fractions and distributions of various phases**
- **Defects such as cracks and voids:  $\mu\text{m}$  to the cm regime**

Phase contrast microscopy works by introducing an additional phase shift between  $R_0$  and  $R_1$ , creating a difference in amplitude of light at  $C$  and  $D$ . That is, contrast is created. This phase shift is introduced by placing a phase plate in the back focal plane of the objective lens. In positive phase contrast the first-order diffracted beam is shifted an additional  $90^\circ$  relative to the zeroth-order (direct) beam. The result is that  $R_1$  is directed opposite to  $R_0$  so that they subtract at image point  $C$  and give a minimum of intensity. At image point  $D$ ,  $R_1$  is further retarded by  $90^\circ$  with respect to  $R_0$  and is therefore at right

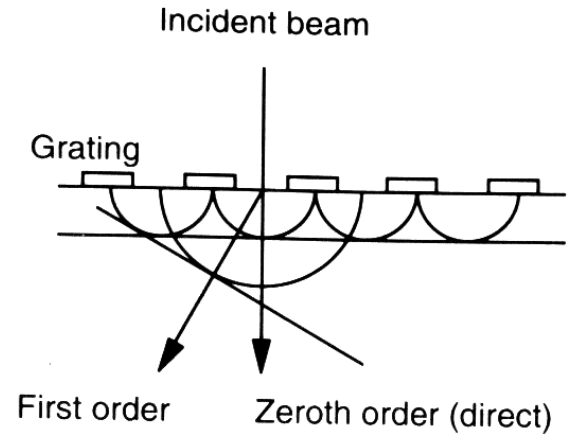
**Phase Contrast Microscopy** - Phase contrast microscopy, first described in 1934 by Dutch physicist Frits Zernike, is a contrast-enhancing optical technique that can be utilized to produce high-contrast images of transparent specimens such as living cells, microorganisms, thin tissue slices, lithographic patterns, and sub-cellular particles (such as nuclei and other organelles). *In effect, the phase contrast technique employs an optical mechanism to translate minute variations in phase into corresponding changes in amplitude, which can be visualized as differences in image contrast.* One of the major advantages of phase contrast microscopy is that living cells can be examined in their natural state without being killed, fixed, and stained. As a result, the dynamics of ongoing biological processes in live cells can be observed and recorded in high contrast with sharp clarity of minute specimen detail.

### 33.5 PHASE CONTRAST MICROSCOPY

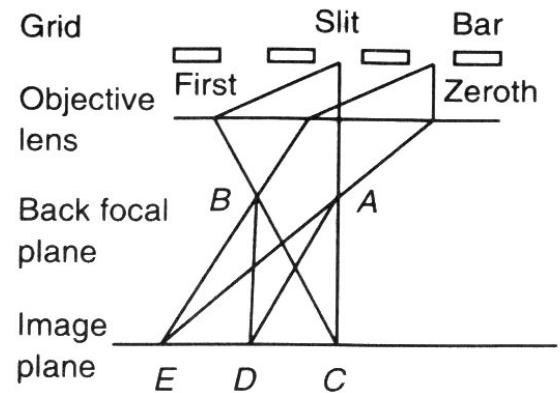
Optical phase contrast microscopy is used with essentially transparent specimens and is probably most widely used with biological materials. It is important to have some degree of understanding of phase contrast microscopy with respect to materials characterization, not only for its direct use but also because of the analogy to phase contrast effects in electron microscopy. The subject is complicated and a full treatment is beyond the scope of this book. An attempt is made to present the basic ideas in enough detail to give a physical understanding of the origin of the enhanced contrast and the techniques for achieving such contrast. The treatment here is an abbreviated version of the simplified vector model given by Goldstein (Goldstein, 1982).

An incident beam of parallel light is shown illuminating a specimen in Figure 33.4. According to Huygens' construction each point of the specimen can be considered to emit a spherical wave and the resulting waves added to give the result of scattering of light by the specimen (Born and Wolf, 1975). The Abbe theory of image formation considers the behavior of a grating in the specimen to treat the formation of an image and its resolution and contrast. The figure shows a grating made up of transmitting portions (slots) and opaque portions (bars). A spherical wave is emitted from the center of each slot. The envelope of these waves gives the zeroth-order diffraction maximum, which is the direct beam. The drawing also shows the formation of the first-order diffraction maximum by interference of light at a distance one wavelength from the first slot with light at a distance two wavelengths from the second slot, and so on. This is the same effect as one-dimensional Bragg diffraction of X-rays or electrons.

Figure 33.5 shows the formation of a diffraction pattern and an image from the diffracted beams. Only the zeroth- and first-order diffracted beams are shown, but these are sufficient to illustrate the principles involved. Each diffracted beam consists of parallel light, and is brought into focus in the back focal plane of the objective lens. The position of focus in the back focal plane will depend on the angle that the diffracted beam makes with the optic axis of the microscope. Thus the direct beam is focused into a



**Figure 33.4** Diffraction from a grating



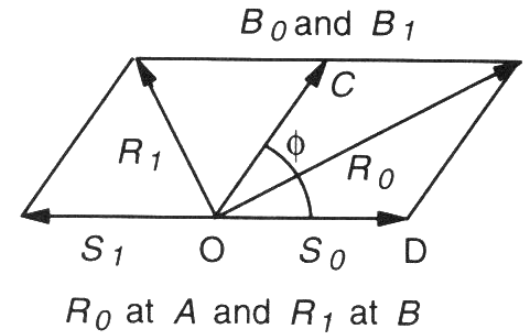
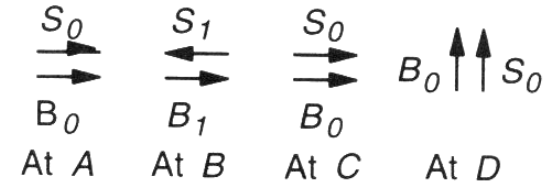
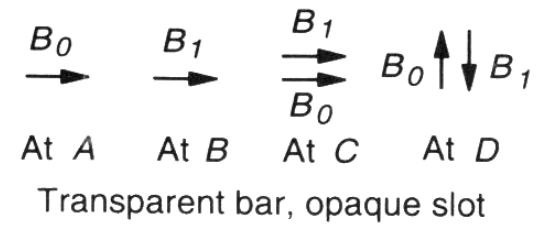
**Figure 33.5** Diffraction pattern and image

diffracted spot at point *A* on the optic axis, while the first-order diffracted beam is focused into a diffraction spot at point *B* off the optic axis. The rays of light continue to form an image of the grating in the image plane. Points *A* and *B* are equidistant from point *C* (the back focal “plane” is approximately, but not strictly, a plane), and the rays from *A* and *B* constructively interfere to form an intensity maximum at *C* and produce an image of a slit. At point *D* the waves are completely out of phase and destructively interfere, forming an image of a bar of the grating. At point *E* the waves are again in phase and interfere constructively to form the image of another slit of the grating.

To understand the principle of phase contrast microscopy one considers the situation of bars and slits that are transparent but that have different indexes of refraction. If the indexes of refraction were the same, each bar would form a spherical wave centered on the bar, as do the slits in Figure 33.4. For the zeroth-order beam all the waves are in phase, giving a strong direct beam. For the first-order beam, the waves from the bars are exactly out of phase with the waves from the slits so that the intensity of the first-order diffracted beam is zero. The Huygens’ construction and diffraction picture thus correctly describes the transmission of light through a specimen that is transparent and has a uniform index of refraction.

The phase relations are shown in more detail in Figure 33.6. The arrows show amplitude (length of arrow) and phase (angle of arrow with retardation taken counterclockwise). At location *A* of Figure 33.5 (the central diffraction spot in the back focal plane of the objective lens), the phase is taken as zero and the arrow is labeled  $B_0$  for bar, zeroth-order. Here, following Goldstein, a transparent bar has been taken to be on the optic axis and the slit has been taken for the moment to be opaque. At location *B*, the first-order diffraction from the bar  $B_1$  is of the same magnitude and phase as  $B_0$ . Point *C* is located an integral number of wavelengths from both *A* and *B* so that  $B_0$  and  $B_1$  are parallel and equal in size. Point *E* is located a half-wavelength farther from *A* than *C*, and a half-wavelength closer to *B* than *C*. At *E* the waves will again constructively interfere. Point *D* is located one-quarter wavelength farther from *A* than *C* and one-quarter wavelength closer to *B* than *C*. Comparing the arrows at *D* with those at *C* it is apparent that  $B_0$  has been rotated  $90^\circ$  clockwise (one-quarter wavelength retardation), and  $B_1$  has been rotated  $90^\circ$  counterclockwise (one-quarter wavelength advance in phase). Vectors  $B_0$  and  $B_1$  cancel, giving zero intensity at *D*.

In the second portion of Figure 33.6 both bar and slit are taken to be transparent and to have equal index of refraction. At *A* the amplitudes from the bar  $B_0$  and from the slit  $S_0$  are of equal amplitude and in phase. At *B* the phase of  $S_1$  is retarded by  $180^\circ$ , so that there is zero total first-order diffracted amplitude. At *C* the contributions are in phase and



Transparent bar, transparent slot, unequal index

**Figure 33.6** Vector model of phase contrast

add. At  $D$  both are retarded by one-quarter wavelength and still add. The result is uniform brightness.

In the third portion of Figure 33.6 the bars are taken to have a higher index of refraction than the slits. This index difference causes both  $B_0$  and  $B_1$  to be retarded in phase by an angle  $\phi$ . At point  $A$ ,  $S_0$  and  $B_0$  add to a resultant  $R_0$  as shown. At point  $B$ ,  $S_1$  and  $B_1$  add to a resultant  $R_1$ . The magnitudes of all vectors are equal and the direction of  $S_1$  is opposite to that of  $S_0$ . Using boldface for vectors and taking the dot product gives

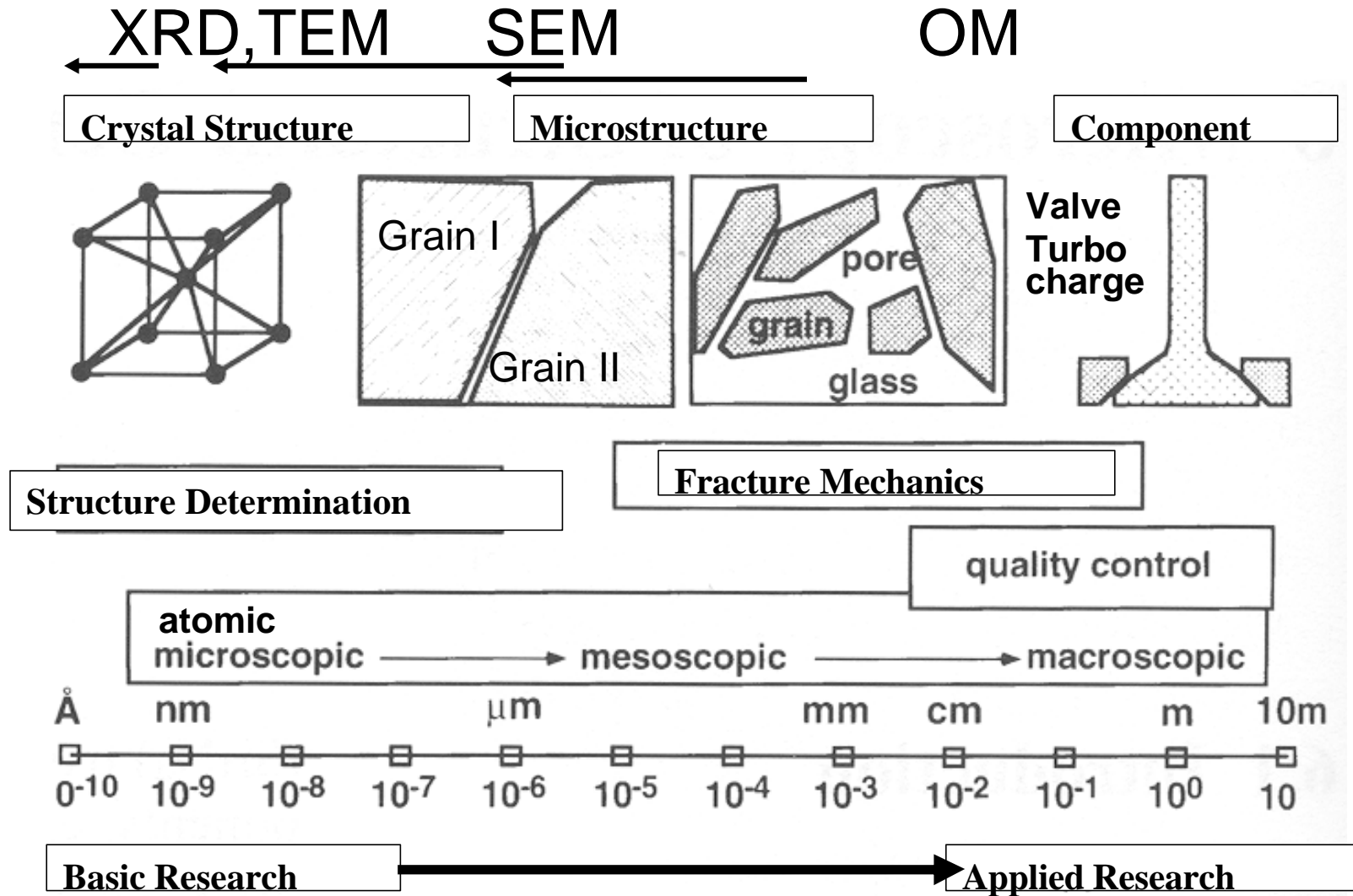
$$\mathbf{R}_0 \cdot \mathbf{R}_1 = (\mathbf{S}_0 + \mathbf{B}_0) \cdot (-\mathbf{S}_0 + \mathbf{B}_0) = 0 \quad (33.9)$$

Thus  $R_0$  and  $R_1$  are perpendicular; i.e., the waves are  $\pi/4$  out of phase.

The image at  $C$  will be formed by interference between  $R_0$  and  $R_1$  with no further change of phase and so will be twice  $OC$  in amplitude and phase. The image at  $D$  will be formed by  $R_0$  and  $R_1$  after a further phase change of  $\lambda/4$  and so will be twice  $OD$ . The eye responds to intensity, which is the square of the amplitude, so that the same intensity will be seen at  $C$  and  $D$  and there will be no contrast in the image despite the difference in phase of the light at  $C$  and  $D$ .

Phase contrast microscopy works by introducing an additional phase shift between  $R_0$  and  $R_1$ , creating a difference in amplitude of light at  $C$  and  $D$ . That is, contrast is created. This phase shift is introduced by placing a phase plate in the back focal plane of the objective lens. In positive phase contrast the first-order diffracted beam is shifted an additional  $90^\circ$  relative to the zeroth-order (direct) beam. The result is that  $R_1$  is directed opposite to  $R_0$  so that they subtract at image point  $C$  and give a minimum of intensity. At image point  $D$ ,  $R_1$  is further retarded by  $90^\circ$  with respect to  $R_0$  and is therefore at right

# Scale and Microscopy Techniques



**Microstructure ranging from crystal structure to Engine components ( $\text{Si}_3\text{N}_4$ )**

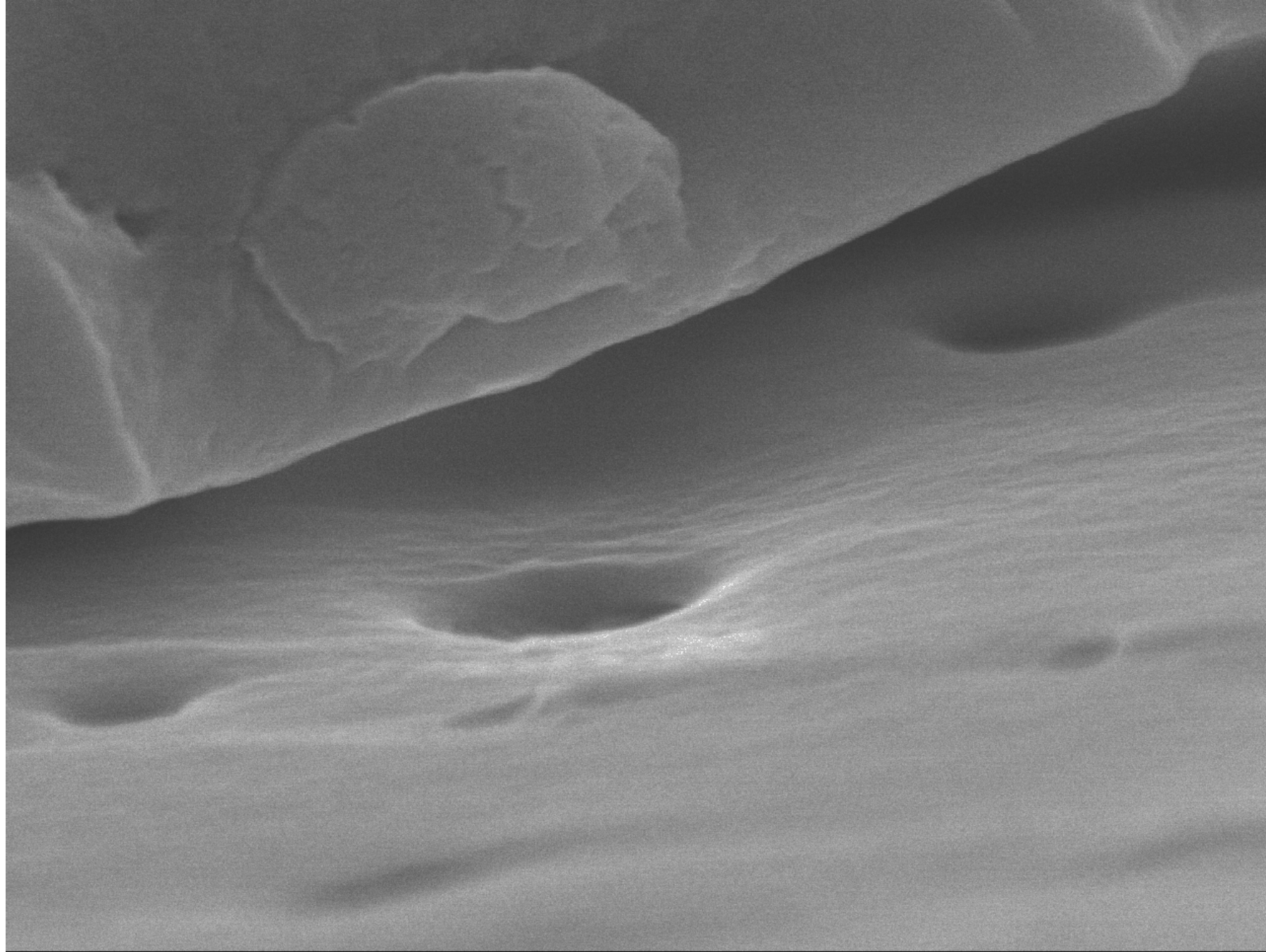


# Advantages of Using SEM over OM

Mag	Depth of Field	Resolution
OM: 4x – 1000x	15.5mm – 0.19mm	~ 0.2mm
SEM: 10x – 500Kx	4mm – 0.4mm	1-10nm

The SEM has a large depth of field, which allows a large amount of the sample to be in focus at one time and produces an image that is a good representation of the three-dimensional sample. The SEM also produces images of high resolution, which means that closely features can be examined at a high magnification.

The combination of higher magnification, larger depth of field, greater resolution and compositional and crystallographic information makes the SEM one of the most heavily used instruments in research areas and industries, especially in semiconductor industry.



NONE

SEI

5.0kV

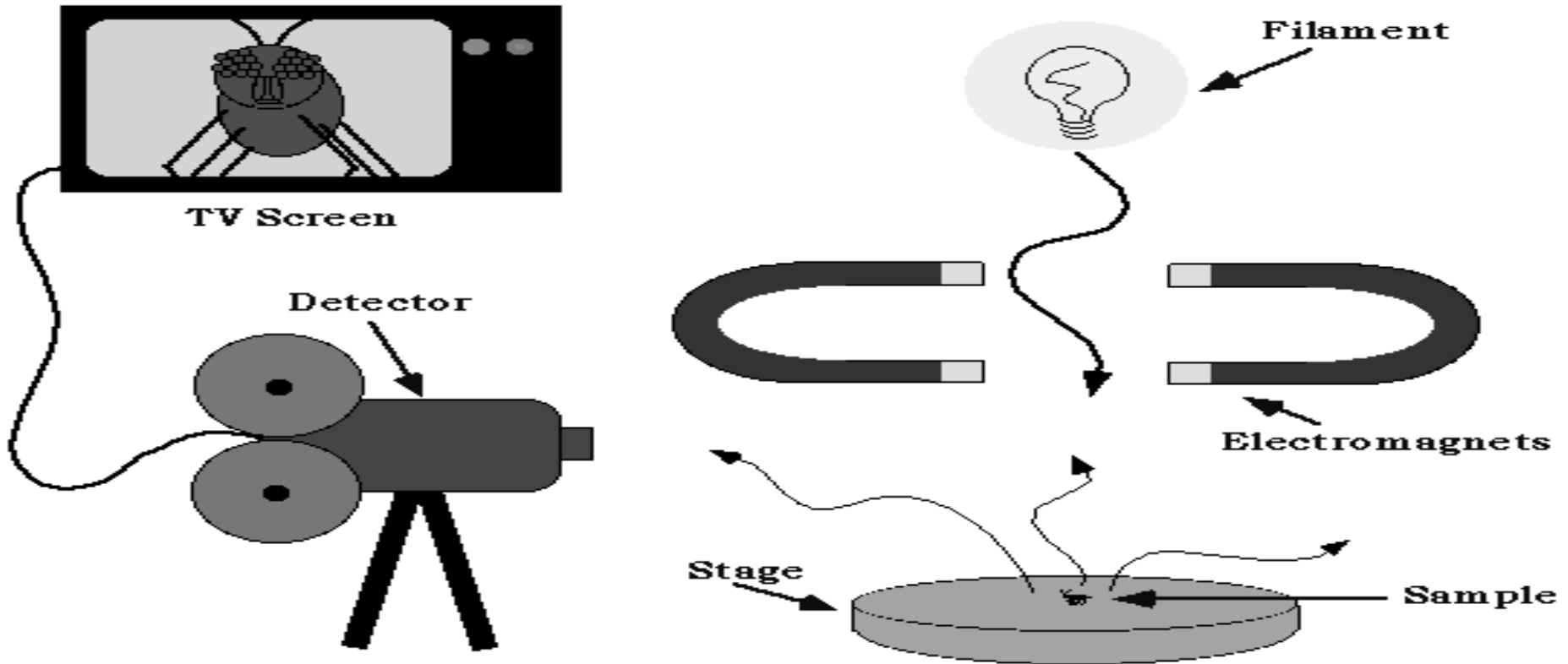
X50,000

100nm

WD 2.9mm

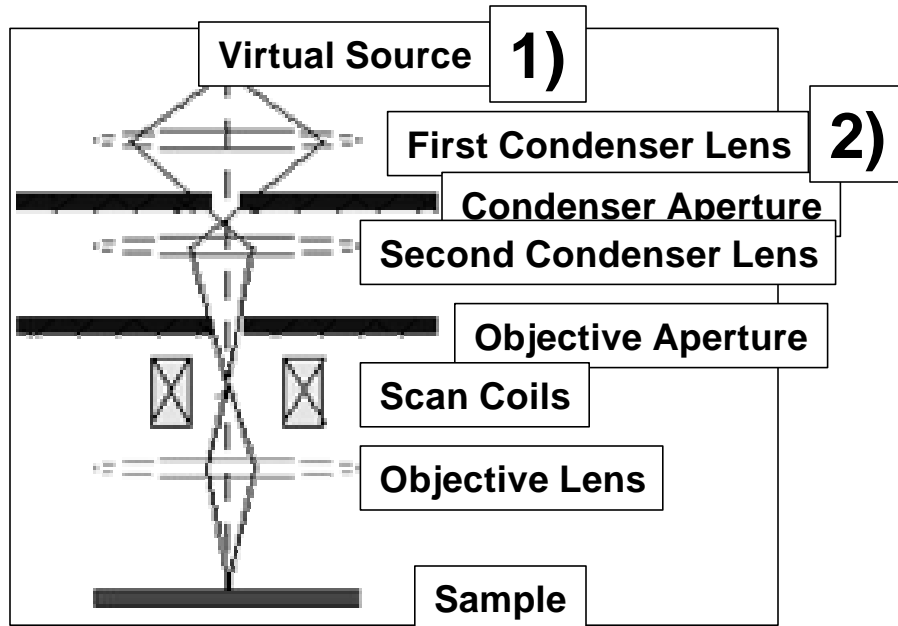
# 2.1 The Microscope Column

## General Layout



In simplest terms, an SEM is really nothing more than a television. We use a filament to get electrons, magnets to move them around, and a detector acts like a camera to produce an image.

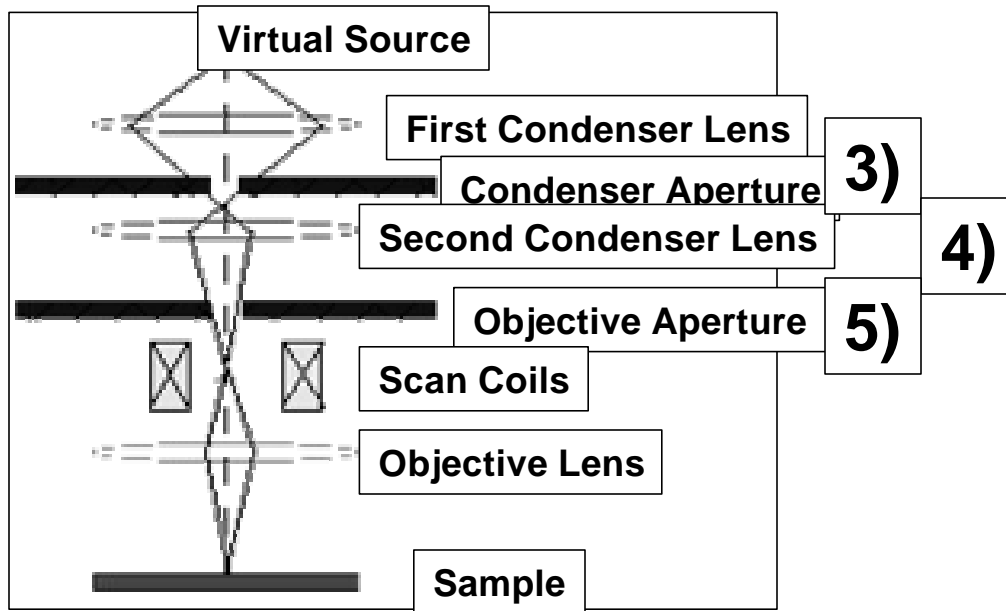
# Scanning Electron Microscope



**1) The "Virtual Source" at the top represents the electron gun, producing a stream of monochromatic electrons.**

**2) The stream is condensed by the first condenser lens (usually controlled by the "coarse probe current knob"). This lens is used to both form the beam and limit the amount of current in the beam. It works in conjunction with the condenser aperture to eliminate the high-angle electrons from the beam.**

# Scanning Electron Microscope

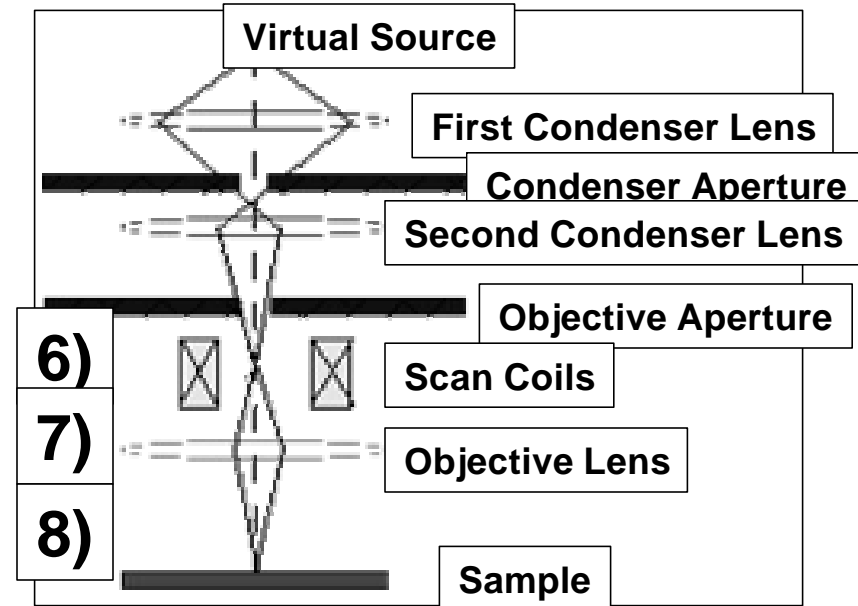


**3) The beam is then constricted by the condenser aperture (usually not user selectable), eliminating some high-angle electrons.**

**4) The second condenser lens forms the electrons into a thin, tight, coherent beam and is usually controlled by the "fine probe current knob".**

**5) A user selectable objective aperture further eliminates high-angle electrons from the beam.**

# Scanning Electron Microscope



**6) A set of coils then "scan" or "sweep" the beam in a grid fashion (like a television), dwelling on points for a period of time determined by the scan speed (usually in the microsecond range).**

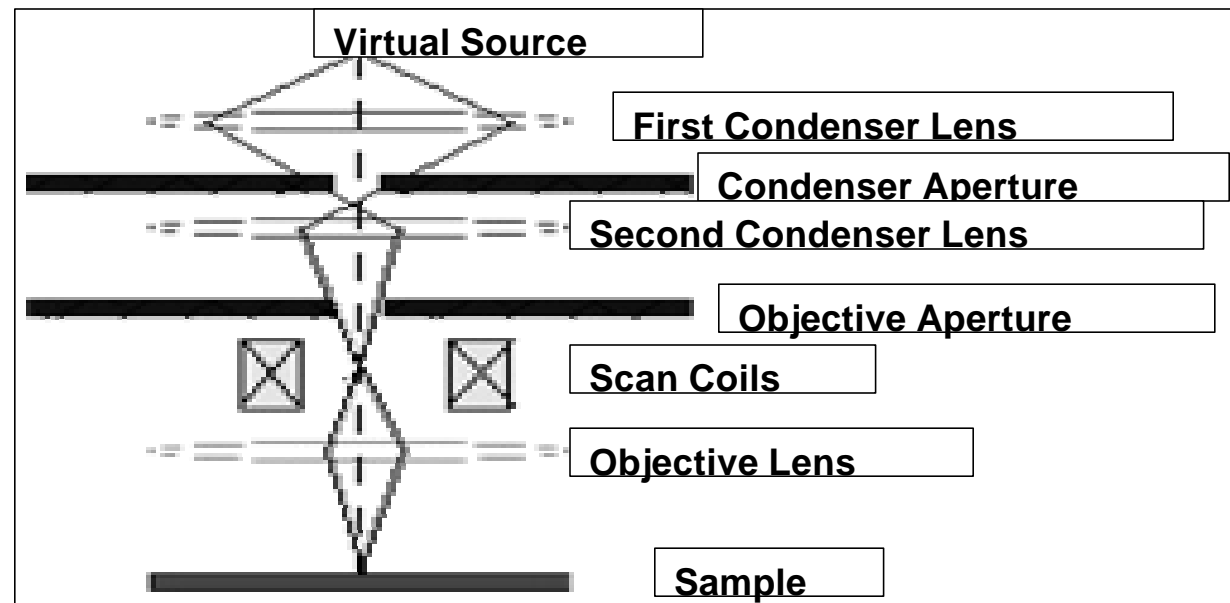
**7) The final lens, the Objective, focuses the scanning beam onto the part of the specimen desired.**

**8) When the beam strikes the sample (and dwells for a few microseconds) interactions occur inside the sample and are detected with various instruments.**

# Scanning Electron Microscope

9) Before the beam moves to its next dwell point these instruments count the number of interactions and display a pixel on a CRT whose intensity is determined by this number (the more reactions the brighter the pixel).

10) This process is repeated until the grid scan is finished and then repeated, the entire pattern can be scanned 30 times/sec.



# Parts of the Microscope

- 1. Electron optical column consists of:**
  - electron source to produce electrons**
  - magnetic lenses to de-magnify the beam**
  - magnetic coils to control and modify the beam**
  - apertures to define the extent of beam, prevent electron spray, etc.**
- 2. Vacuum systems consists of:**
  - chamber which “holds” vacuum**
  - pumps to produce vacuum**
  - valves to control vacuum**
  - gauges to monitor vacuum**



# Parts of the Microscope

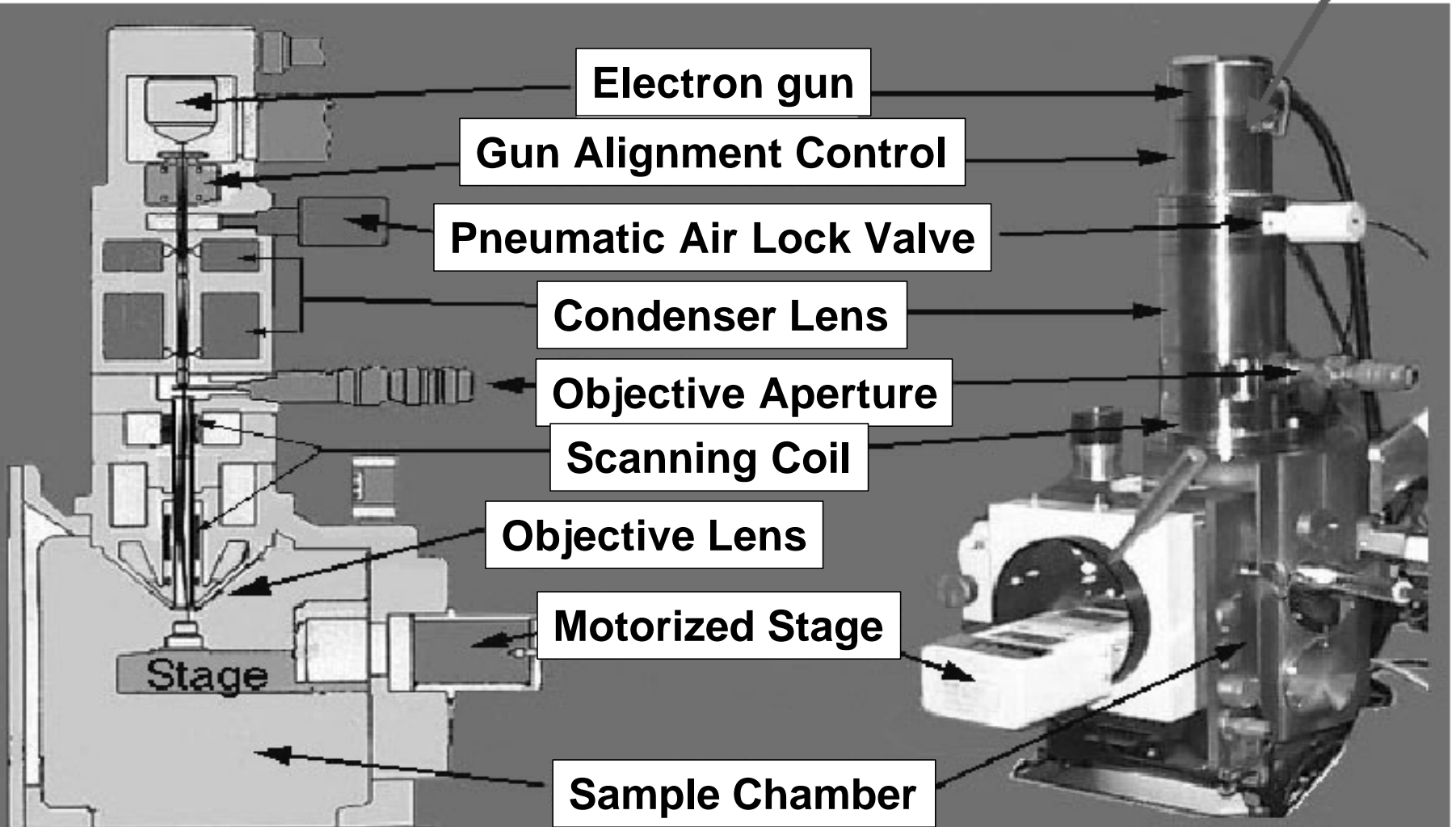
## 3. Signal Detection & Display

consists of:

- detectors which collect the signal
- electronics which produce an image from the signal

# A Look Inside the Column

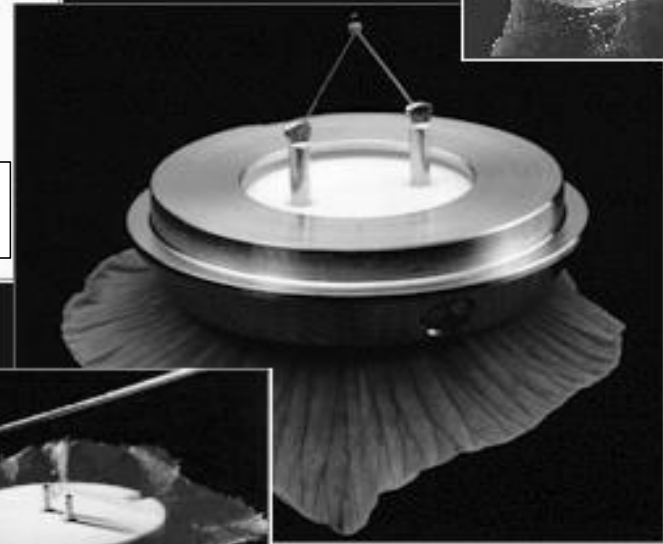
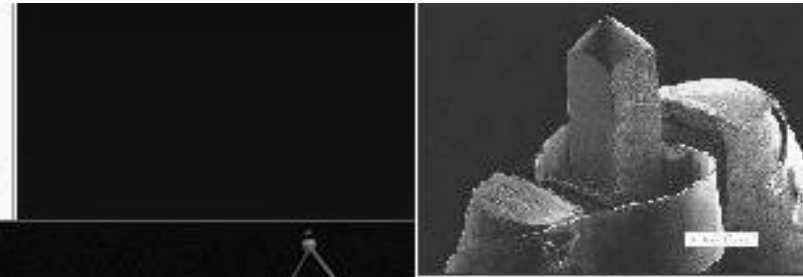
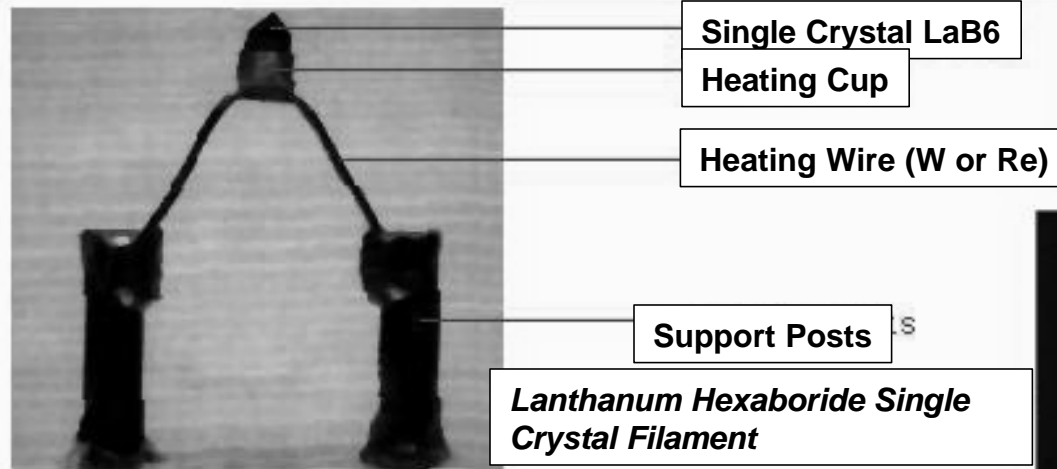
Column



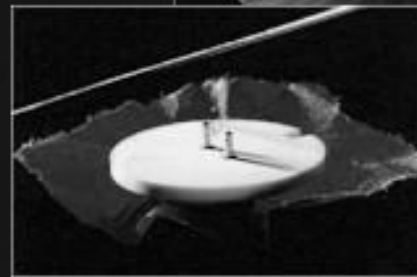
# Getting a Beam

- **A filament is heated to a high T to emit electrons.**
- **A bias (of up to 45 kV) between the filament and a target material draws the electrons to an anode.**
- **Up to this point, the process is identical to that used to form x-rays. However, instead of striking a target to form x-rays, the electrons are drawn towards an anode, where they go through a hole, forming a beam.**
- **This beam is attracted through the anode, condensed by a condenser lens, and focused as a very fine point on the sample by the objective lens.**

# Electron Beam Source



## Electron Beam Source

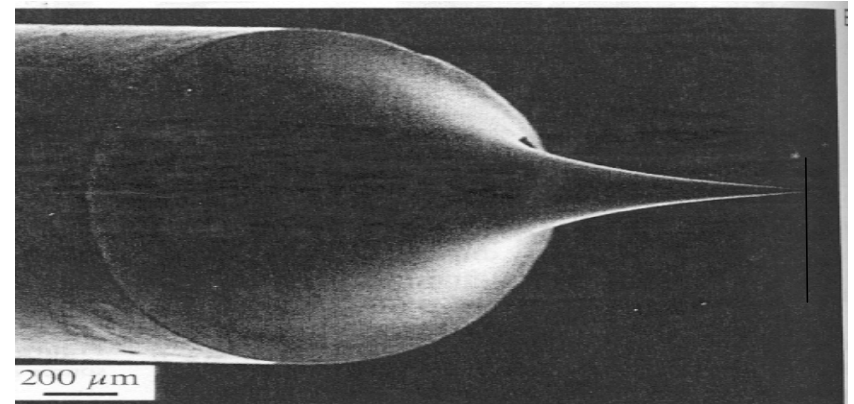
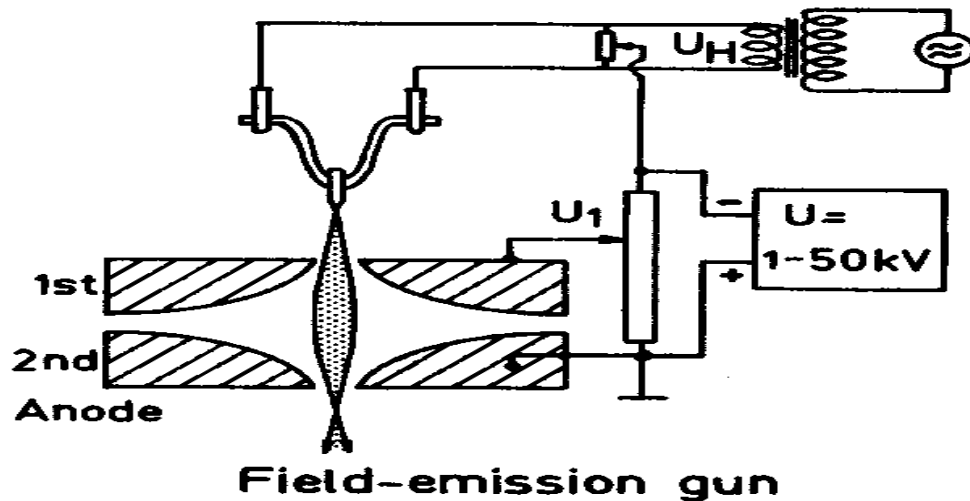


**W / LaB<sub>6</sub> (Filament)**

**Thermionic or Cold Cathode (Field Emission Gun)**

# Field Emission Gun

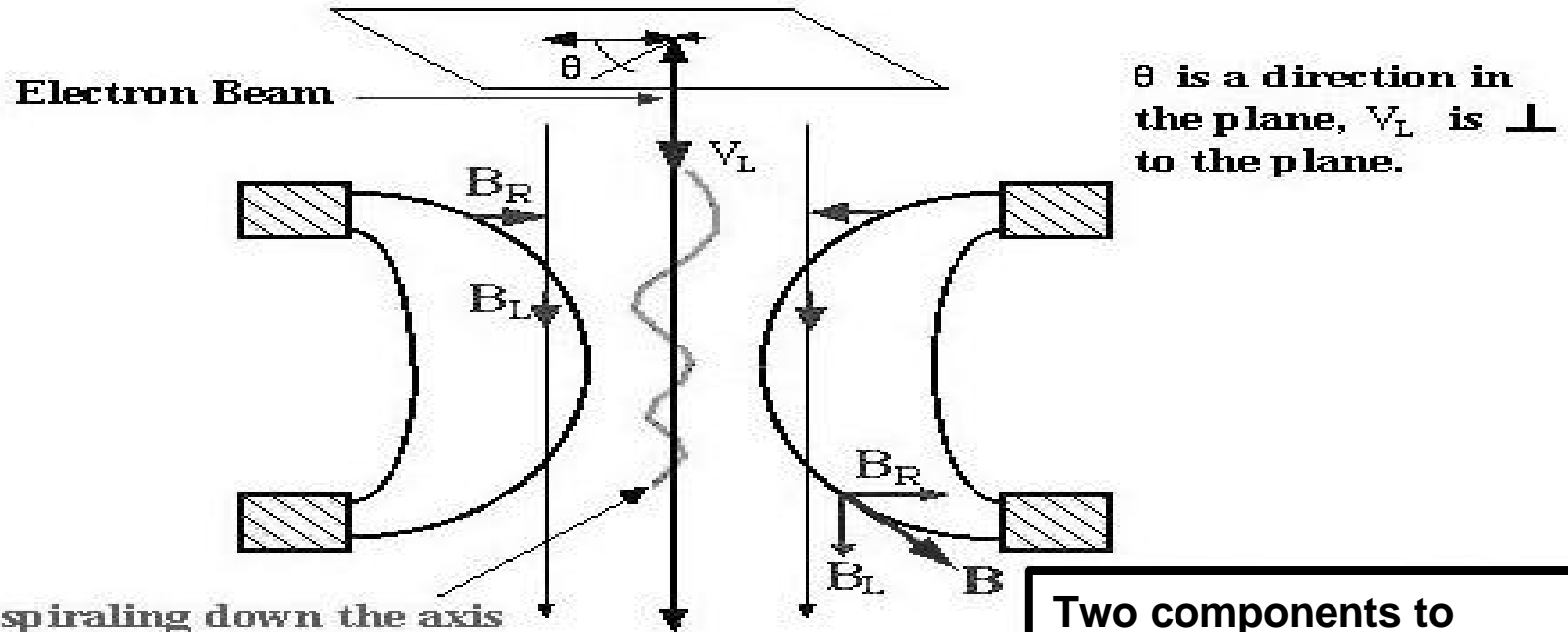
- The tip of a tungsten needle is made very sharp (radius  $< 0.1 \text{ mm}$ )
- The electric field at the tip is very strong ( $> 10^7 \text{ V/cm}$ ) due to the sharp point effect
- Electrons are pulled out from the tip by the strong electric field
- Ultra-high vacuum (better than  $10^{-6} \text{ Pa}$ ) is needed to avoid ion bombardment to the tip from the residual gas.
- Electron probe diameter  $< 1 \text{ nm}$  is possible



# Electromagnetic Lenses

Lorentz force equation:

$$\mathbf{F} = q_0 \mathbf{v} \times \mathbf{B}$$



Nonaxial electrons will experience a force both down the axis and one radial to it. Only electrons traveling down the axis feel equal radial forces from all sides of the lens. The unequal force felt by the off-axis electrons causes spiraling about the optic axis.

Two components to the B field:  
 $B_L$  = longitudinal component (down the axis)  
 $B_R$  = radial component (perpendicular to axis)

after "Electron Beam Analysis of Materials", M.H. Loretto

# Electromagnetic Lenses

- **Condenser lens – focusing determines the beam current which impinges on the sample.**
- **Objective lens – final probe forming**
- **The objective lens determines the final spot size of the electron beam, i.e., the resolution of a SEM.**

# Electromagnetic Lenses

## The Condenser Lens

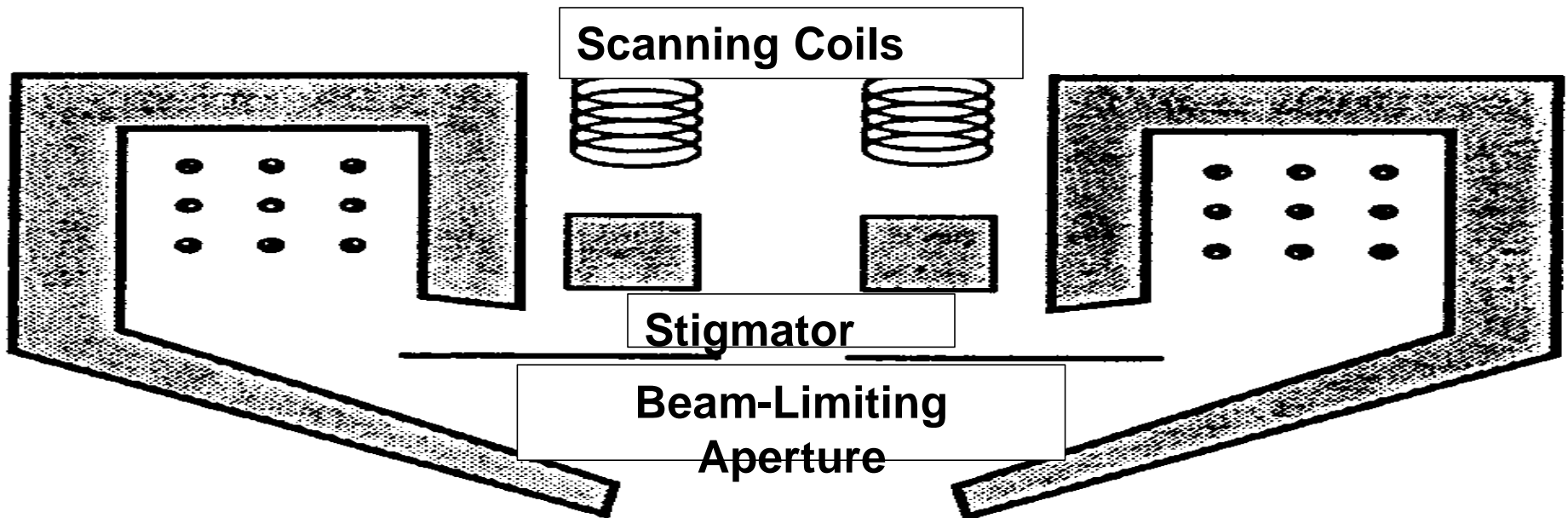
- **For a thermionic gun, the diameter of the first cross-over point  $\sim 20\text{-}50\mu\text{m}$ .**
- **If we want to focus the beam to a size  $< 10\text{ nm}$  on the specimen surface, the magnification should be  $\sim 1/5000$ , which is not easily attained with one lens (say, the objective lens) only.**
- **Therefore, condenser lenses are added to demagnify the cross-over points.**



# Electromagnetic Lenses

## The Objective Lens

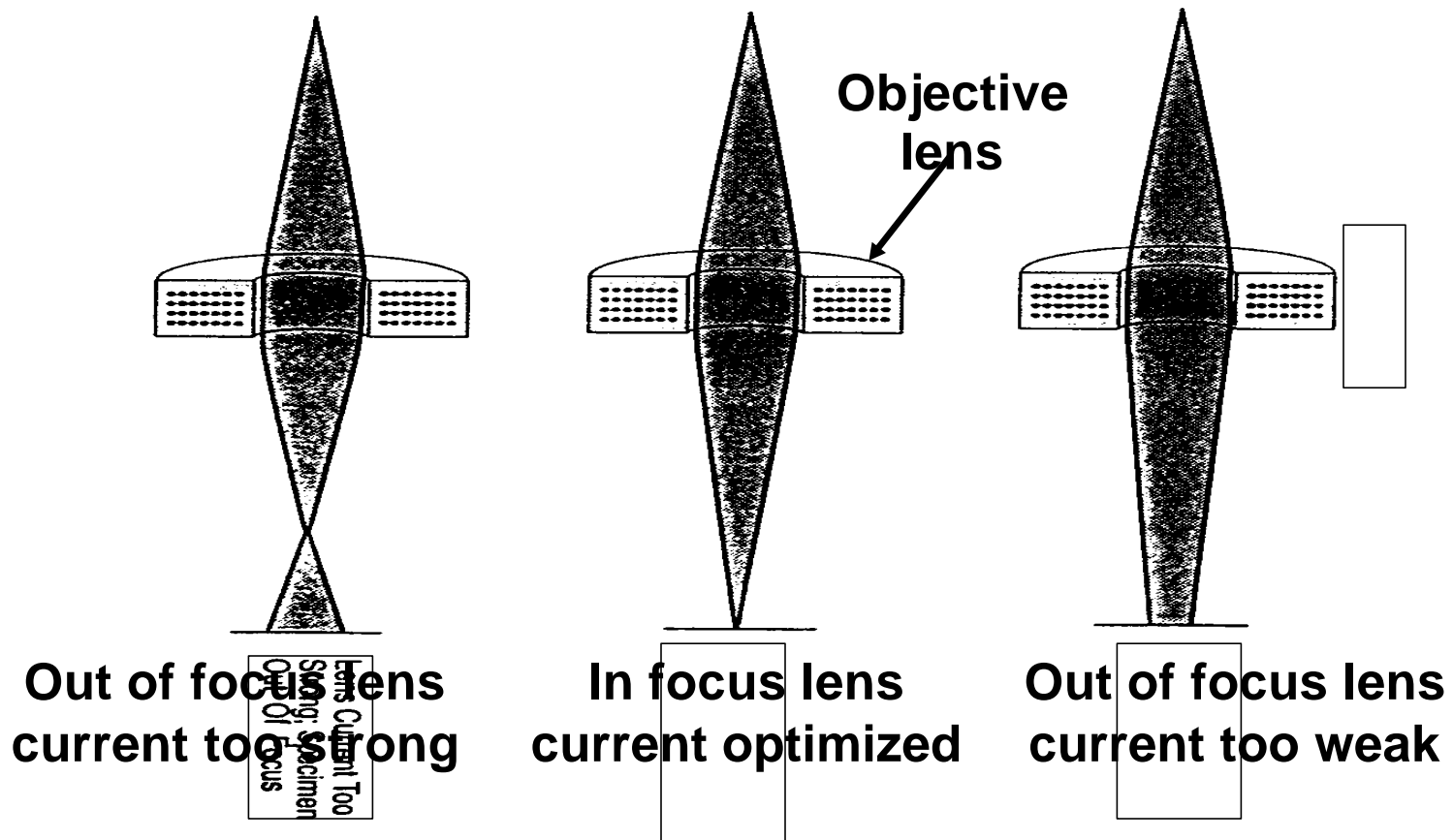
- The objective lens controls the final focus of the electron beam by changing the magnetic field strength
- The cross-over image is finally demagnified to an  $\sim 10\text{nm}$  beam spot which carries a beam current of approximately  $10^{-9} - 10^{-13}$  A.



# Electromagnetic Lenses

## The Objective Lens - Focusing

- By changing the current in the objective lens, the magnetic field strength changes and therefore the focal length of the objective lens is changed.



# Electromagnetic Lenses

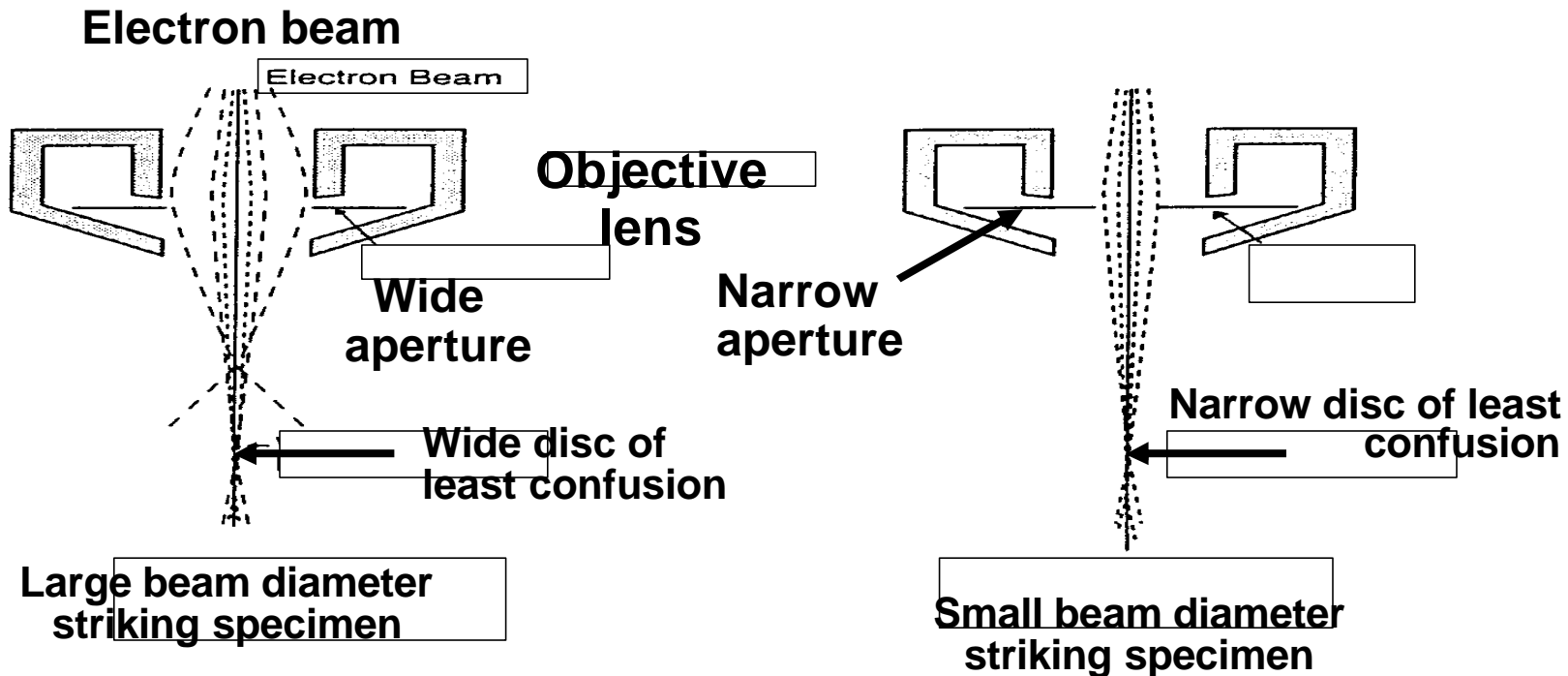
## The Objective Lens - Stigmator

- The objective lens is machined to very high precision and the magnetic field pattern is very carefully designed.
- However, the precision attainable by machining cannot match that required for controlling a beam with a 10 nm diameter.
- The **stigmator**, which consist of two pairs of pole-pieces arranged in the X and Y directions, is added to correct the **minor imperfections** in the objective lens.

# Electromagnetic Lenses

## The Objective Lens - Aperture

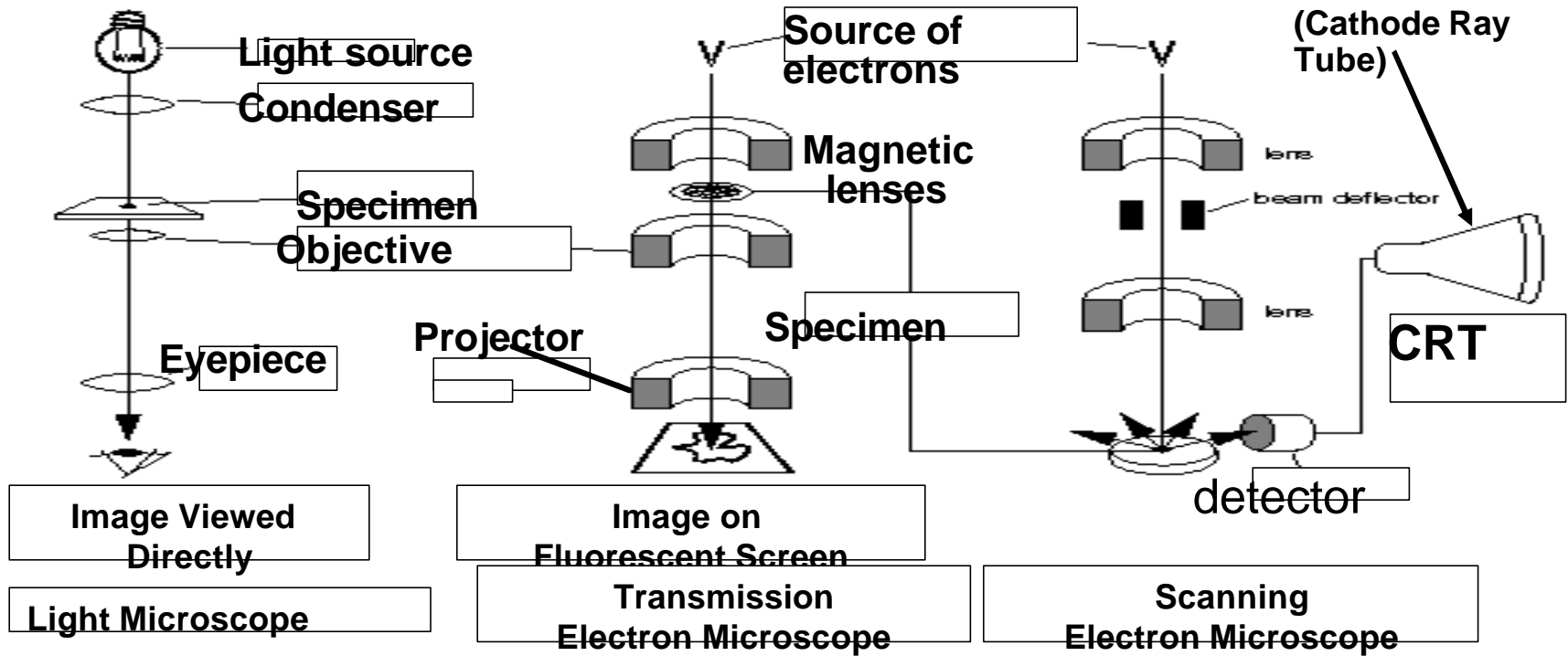
- Since the electrons coming from the electron gun have spread in kinetic energies and directions of movement, they may not be focused to the same plane to form a sharp spot.
- By inserting an aperture, the stray electrons are blocked and the remaining narrow beam will come to a narrow “Disc of Least Confusion”



# Electron vs. Optical Lenses

- **e<sup>-</sup>'s don't actually touch the lens**  
**No definite interface**
- **e<sup>-</sup>'s rotate in the magnetic field**
- **e<sup>-</sup>'s repel each other**
- **$f \propto H \propto I$** 
  - **Focus and magnification controlled electronically**
  - **No physical movements**
- **e<sup>-</sup> lenses can only be positive elements (converging)**
- **Can't correct e<sup>-</sup> lens aberrations like you can with compound optical lenses**
- **e<sup>-</sup> lenses always operate at small apertures**

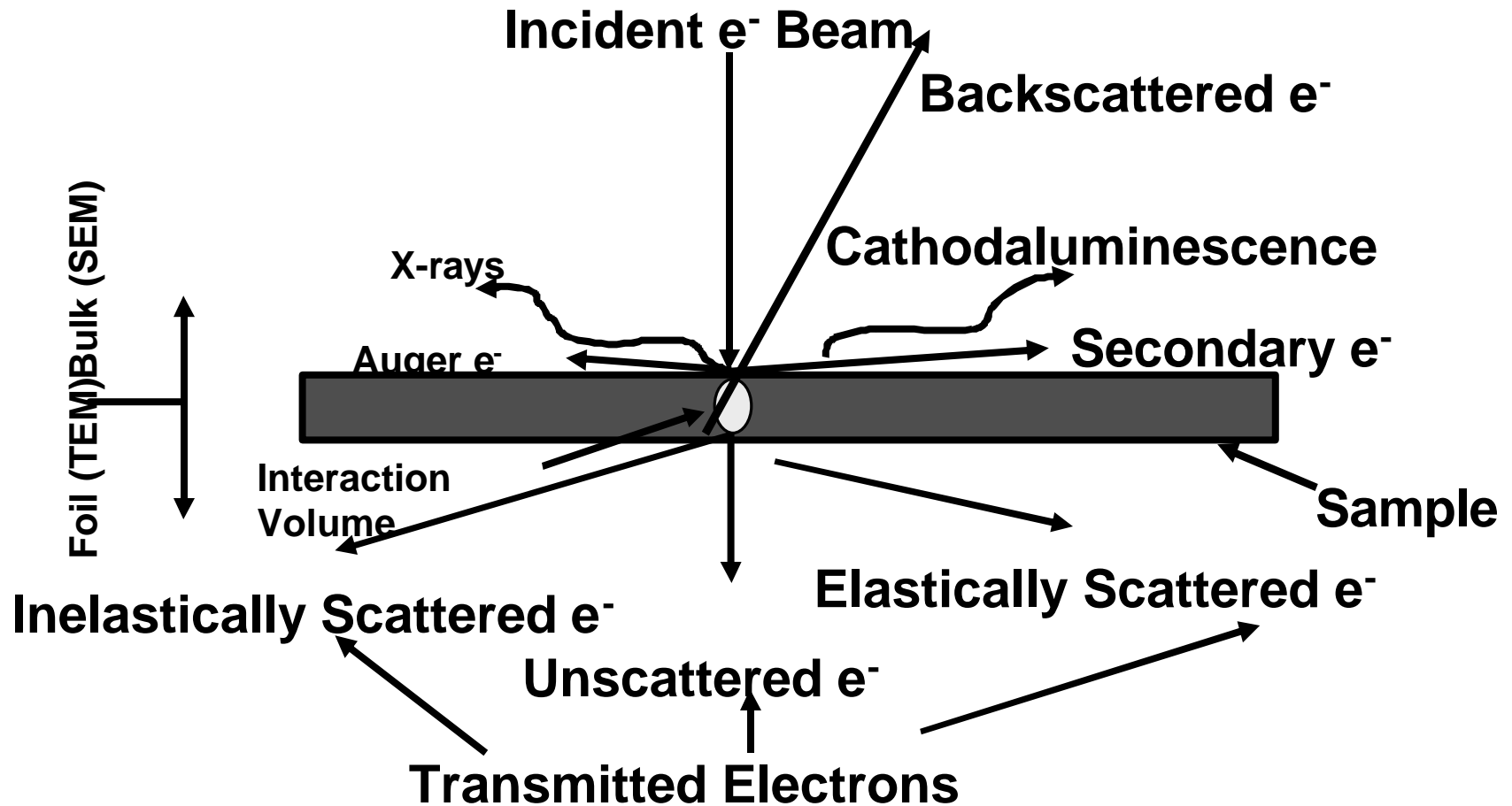
# Comparison of OM, TEM and SEM



**Principal features of an optical microscope, a transmission electron microscope and a scanning electron microscope, drawn to emphasize the similarities of overall design.**

# 2.2 Signal Detection and Display

When an electron beam strikes a sample, a large number of signals are generated.



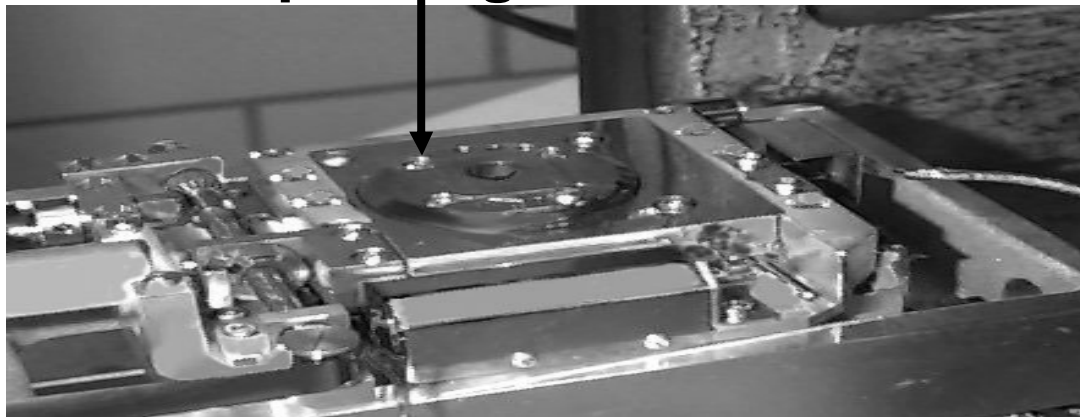
We can divide the signals into two broad categories:

- electron signals, b) photon signals

# Electron Detectors and Sample Stage



Sample stage



Example



# **Specimen Interaction Volume**

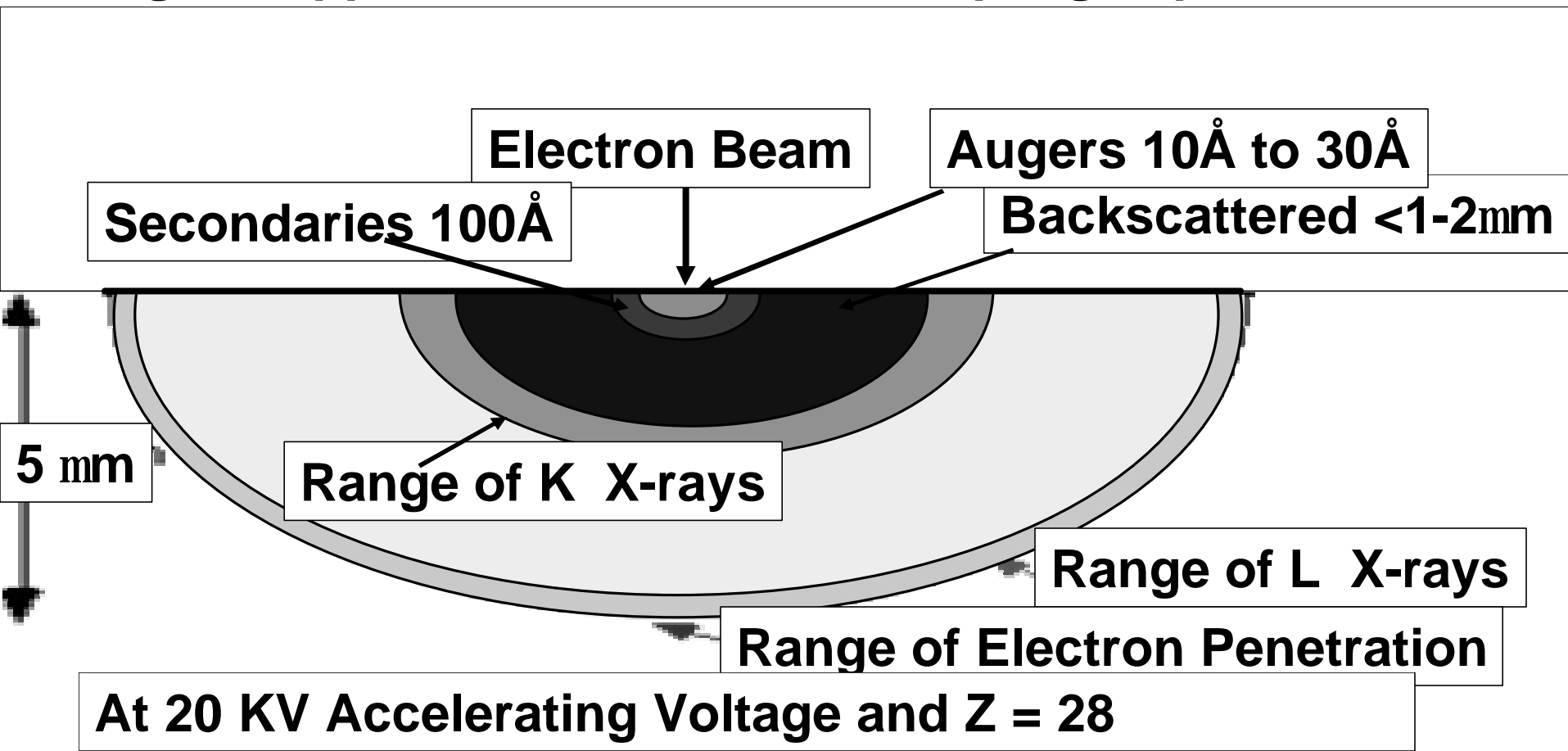
**The volume inside the specimen in which interactions occur while being struck with an electron beam. This volume depends on the following factors:**

- Atomic number of the material being examined; higher atomic number materials absorb or stop more electrons and so have a smaller interaction volume.**
- Accelerating voltage: higher voltages penetrate farther into the sample and generate larger interaction volumes**
- Angle of incidence for the electron beam; the greater the angle (further from normal) the smaller the volume**

# Specimen Interaction Volume

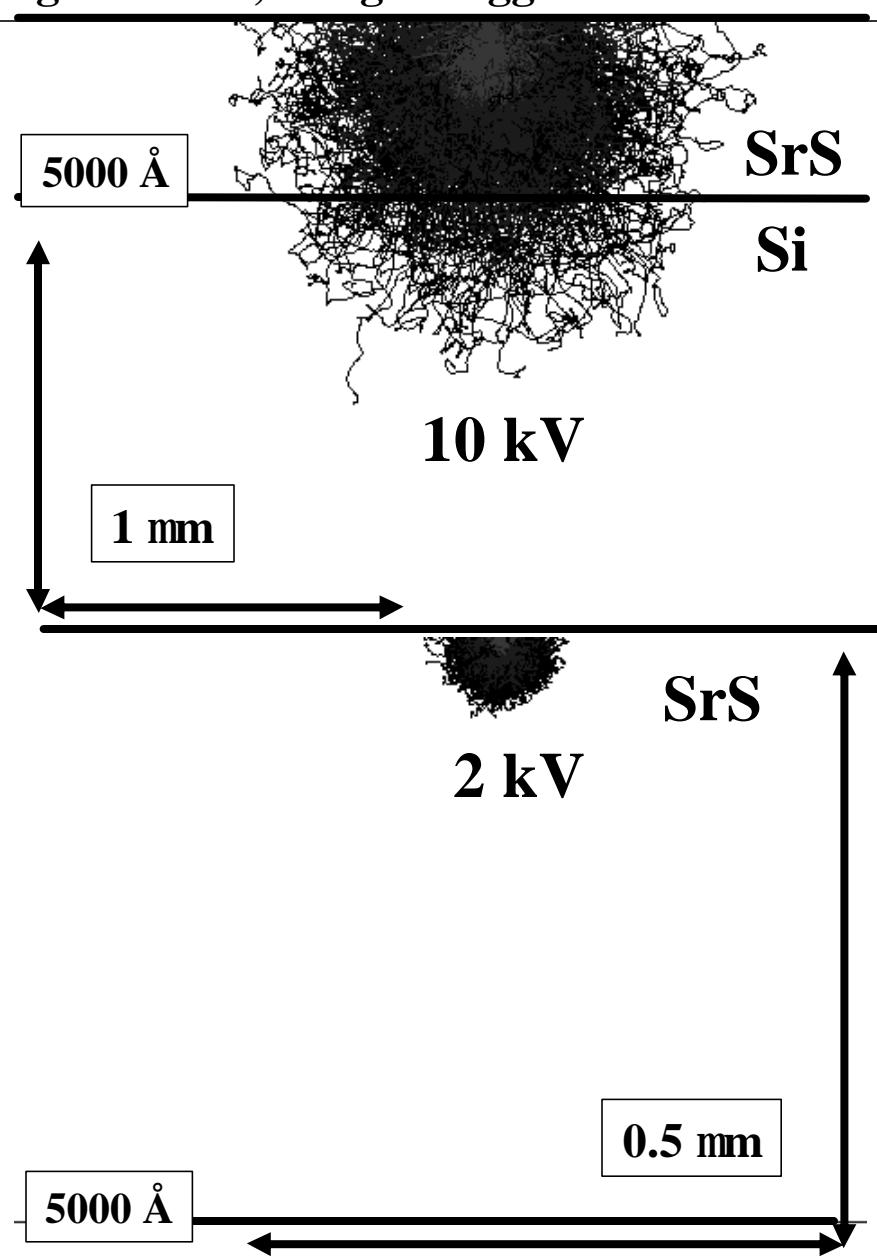
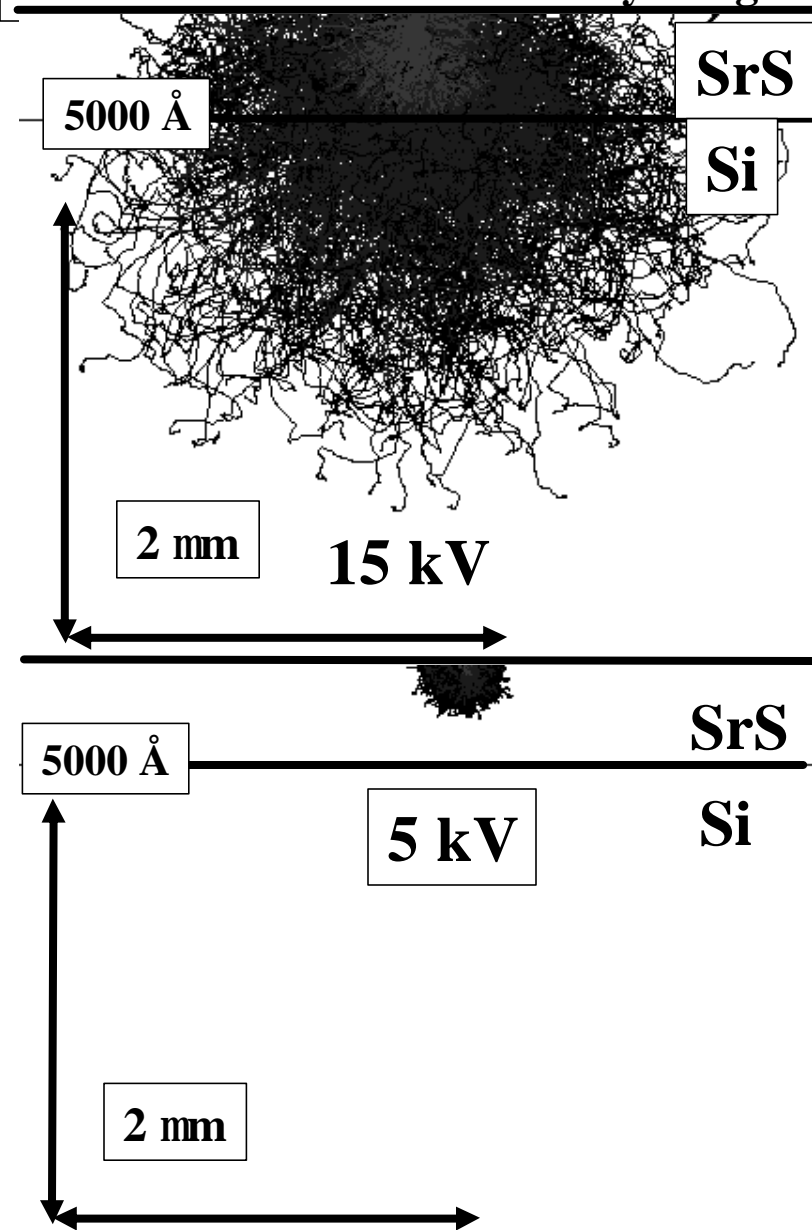
Below is an example of a typical Interaction volume for:

- Specimen with atomic number 28, 20 kV
- 0° degrees tilt, incident beam is normal to specimen surface noting the approximate maximum sampling depths for the



# Beam Interaction Simulations

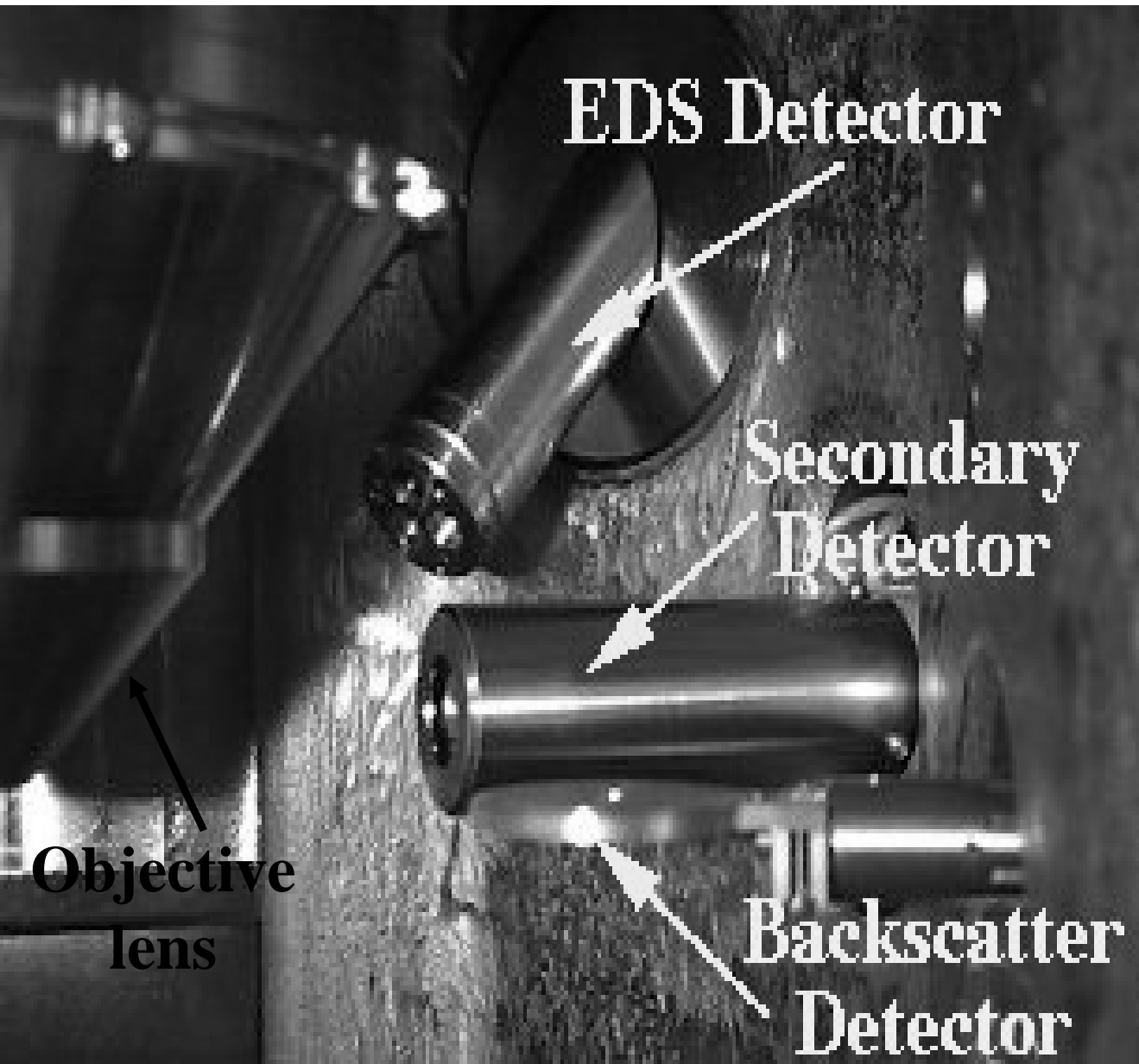
Theoretical SrS density:  $3.7\text{g/cm}^3$ , Long. range:  $1062\text{ \AA}$ , Long. straggle:  $384\text{ \AA}$



# Signal Detection and Display

- If you change the target material, the high and low energy peaks remain (although their intensity may change) while the low intensity peaks change position and are characteristic of the sample.
- The reason we produce this type of profile is because the incident electrons we send into the sample are scattered in different ways.
- There are two broad categories to describe electron scattering:
  - elastic: Backscattered electrons
  - inelastic: Secondary electrons

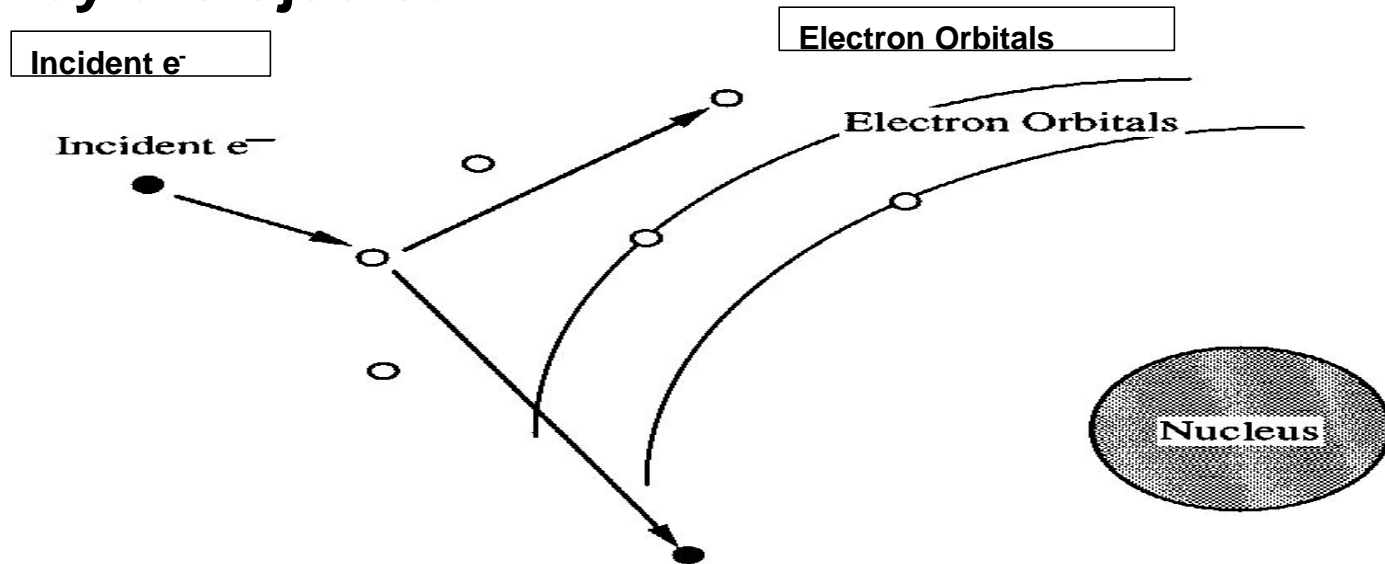
# Electron Detectors



**Example**

# Secondary Electrons

These electrons arise due to inelastic collisions between primary electrons (the beam) and loosely bound electrons of the conduction band (more probable) or tightly bound valence electrons. The energy transferred is sufficient to overcome the work function which binds them to the solid and they are ejected.



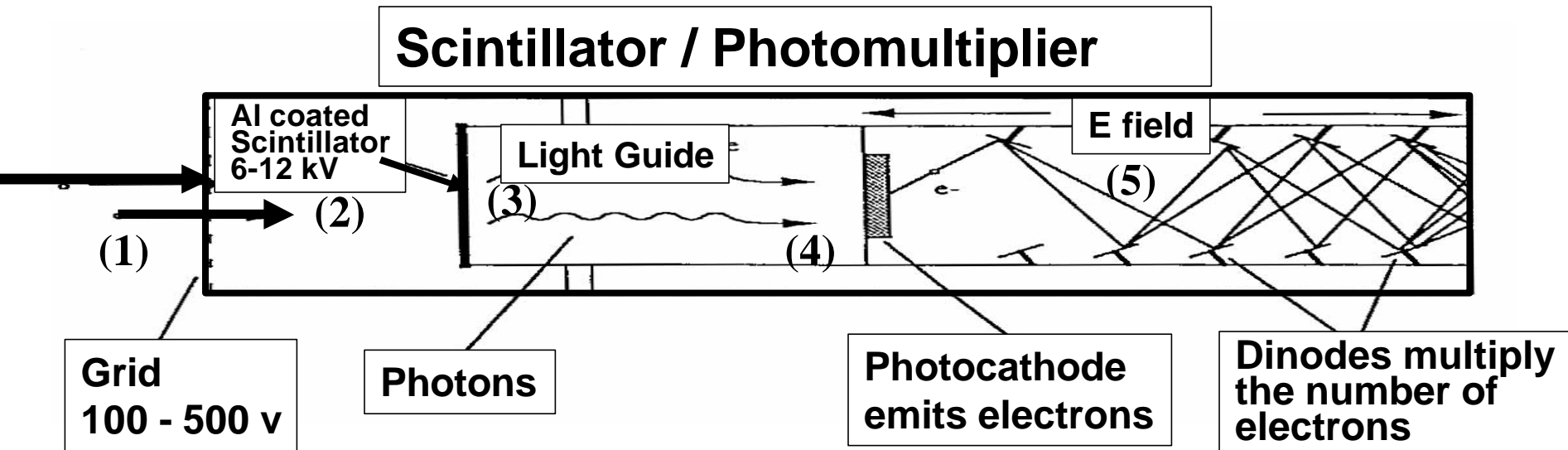
The interaction is Coulombic in nature and the ejected electrons typically have  $\sim 5 - 10$  eV. 50 eV is an arbitrary cut-off below which they are said to be secondary electrons.

# Detection

**Remember, secondary electrons are low energy electrons. We can easily collect them by placing a positive voltage (100 - 300V) on the front of our detector. Since this lets us collect a large number of the secondaries (50 - 100%), we produce a “3D” type of image of the sample with a large depth of field.**

**The type of detector used is called a scintillator / photo-multiplier tube.**

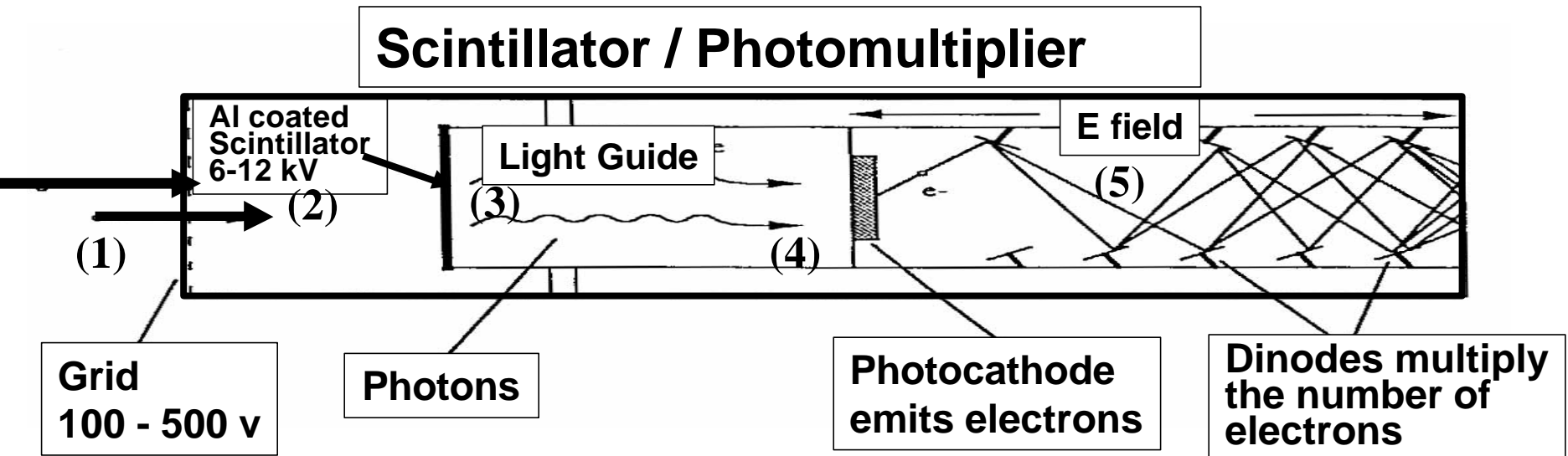
# Detection Sequence



- 1. Secondary electrons (SE) are accelerated to the front of the detector by a bias voltage of 100 - 500 eV.**
- 2. They are then accelerated to the scintillator by a bias of 6 - 12 keV, (10 KeV is normal).**
- 3. Scintillator is doped plastic or glass covered with a fluorescent material (e.g. Europium). A thin (700Å) layer of Al covers it to prevent light from causing fluorescence. The 10keV potential allows the SE to get through the Al and fluoresce.**



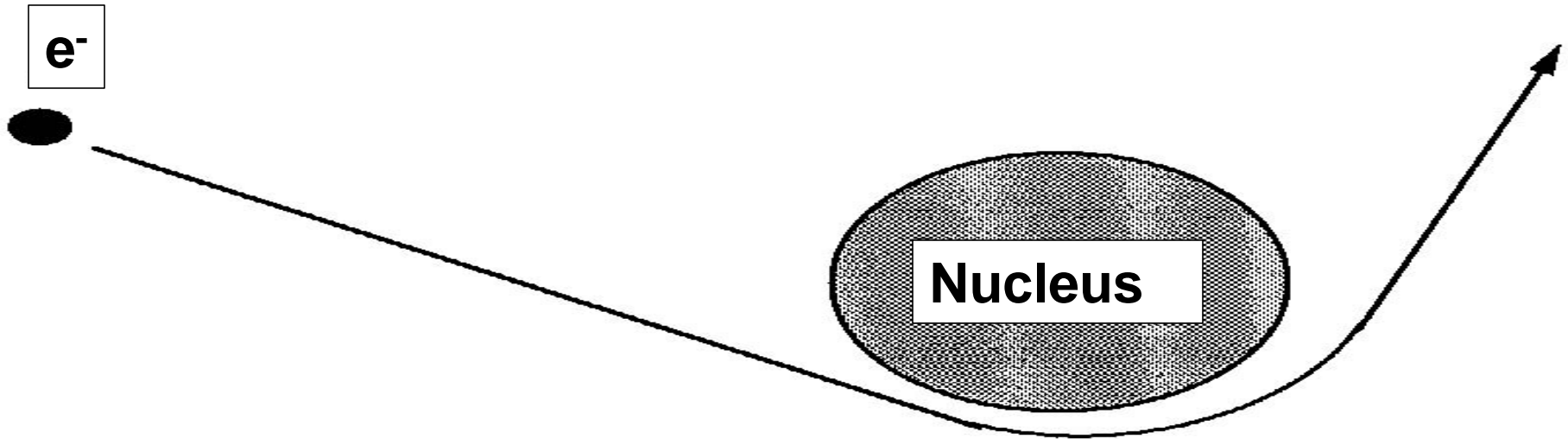
# Detection Sequence



4. The light photons travel down the tube (guide) to a photocathode which converts them into electrons
5. The electrons move through the detector, producing more electrons as they strike dynodes. An output electron pulse is then detected.

# Backscattered Electrons

Backscattered electrons (BSE) arise due to elastic collisions between the incoming electron and the nucleus of the target atom (i.e. Rutherford scattering). Higher  $Z$ , more BSE emitted.

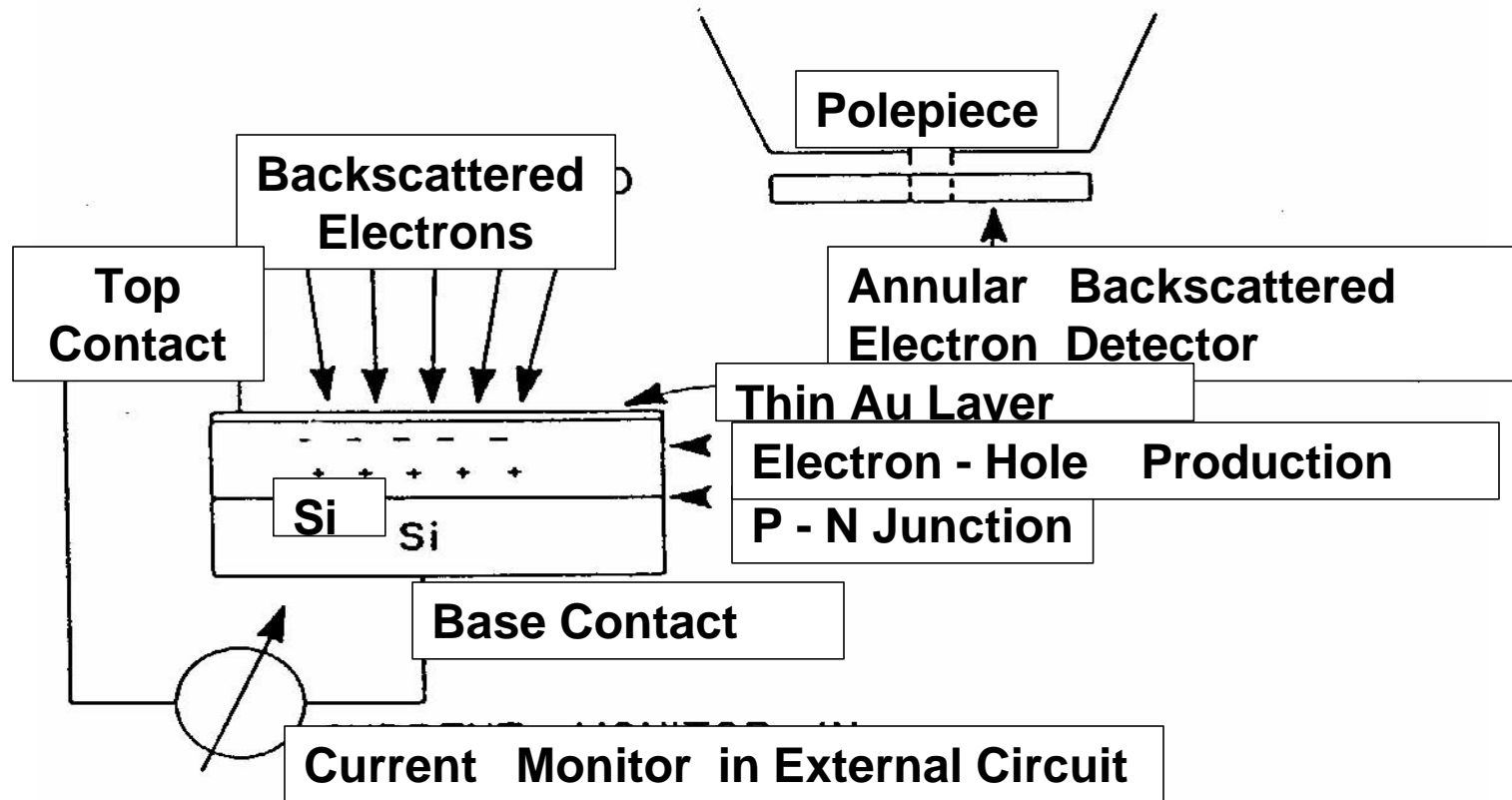


As the name implies, elastic scattering results in little ( $< 1$  eV) or no change in energy of the scattered electrons, although there is a change in momentum ( $p$ ). Since  $p = mv$  and the mass of the electron doesn't change, the direction of the velocity vector must change. The angle of scattering can range from  $0$  to  $180^\circ$ .

# Detection

Since BSE have high energies, they can't be pulled in like secondaries. If you placed a potential on a grid to attract them, you would also attract the incident beam!!

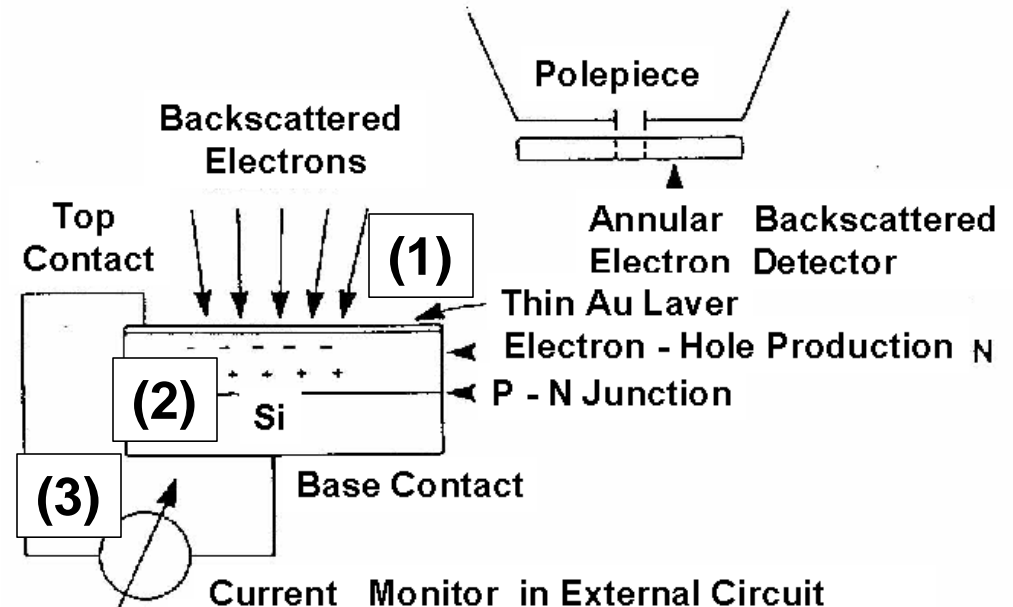
The most common detector used is called a surface barrier detector. It sits above the sample, below the objective lens. BSE which strike it are detected.



# Detection Sequence

Surface barrier detectors are solid state devices made up of semiconducting materials. A semiconducting material has a filled valence band and an empty conduction band- similar to ceramic materials.

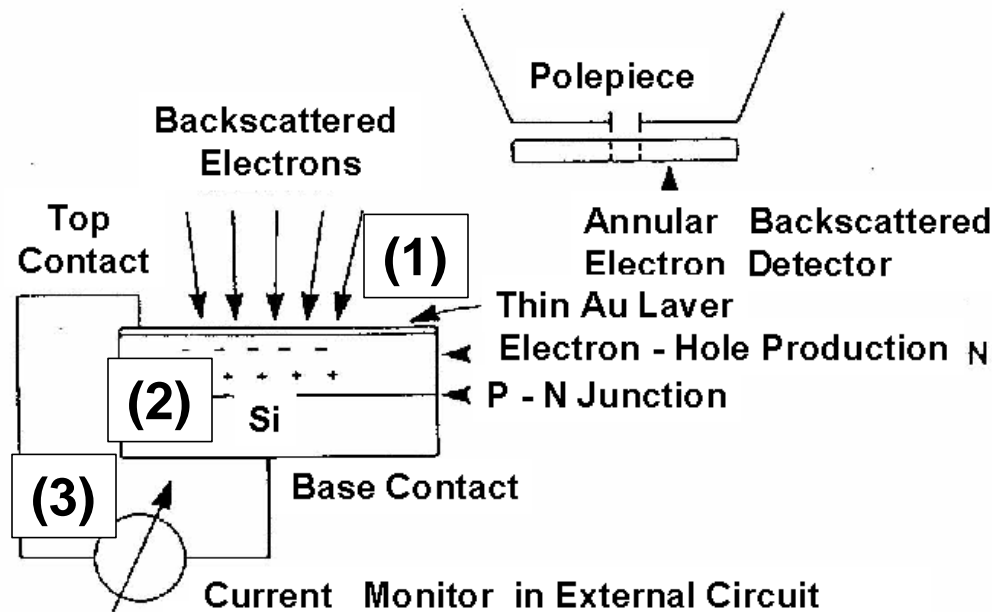
1. When a BSE electron strikes the detector, electrons in the material move from valence to conduction band.
2. The electrons are now free to move in the conduction band or drop back into the valence band.



# Detection Sequence



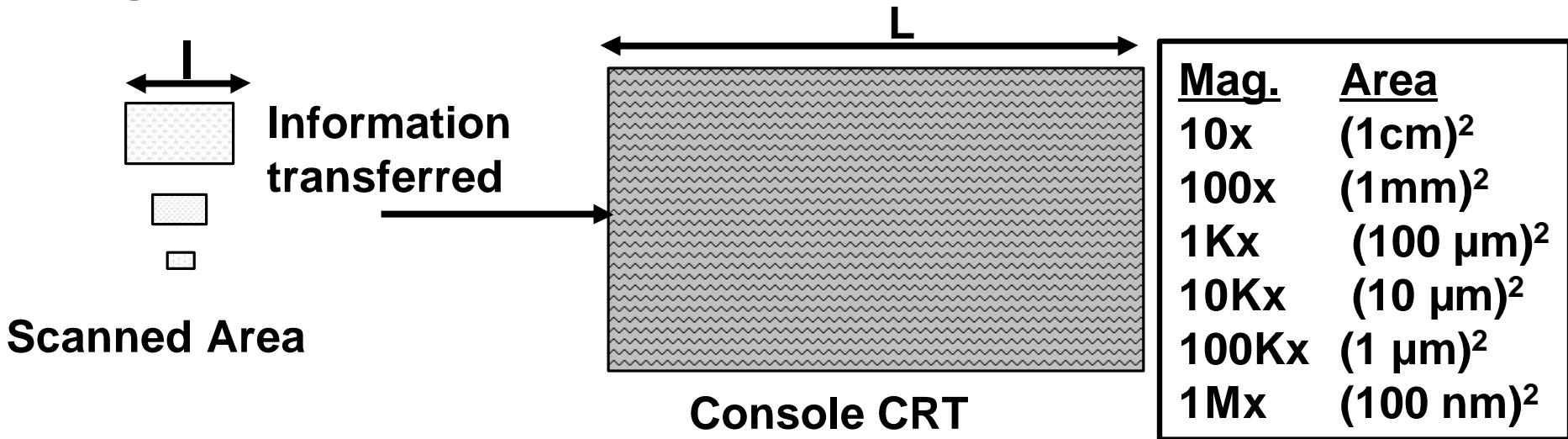
3. If a potential is applied, the  $e^-$  and  $e^+$  can be separated, collected, and the current measured. The strength of the current is proportional to the number of BSE that hit the detector.



# 2.3 Operating Parameters

## Magnification

An image is obtained by taking the signal from the sample and transferring it to a CRT screen. By decreasing the size of the scanned area (from which we get the signal), magnification is produced.



Magnification is determined by taking the ratio of the lengths of the scans:

$$\text{Mag.} = L/l$$

# Resolution

**Resolution is the ability to resolve two closely spaced points. While you may have to be at a high magnification to see small features, resolution is NOT the same as magnification.**

**One way to improve resolution is by reducing the size of the electron beam that strikes the sample:**

$$d_{\min} = 1.29C_s^{1/4}l^{3/4}[7.92(iT/J_c)\times 10^9 + 1]^{3/8}$$

**at low current:**

$$d_{\min} = 1.29C_s^{1/4}l^{3/4}$$

**$J_c$  = current density of the source,  $l$  = electron wavelength**

**$C_s$  = spherical aberration,  $i$  = current,  $T$  = temperature,**

# Resolution

**We can also improve the resolution by:**

- Increasing the strength of the condenser lens**
- Decreasing the size of the objective aperture**
- Decreasing the working distance (WD = the distance the sample is from the objective lens)**



# Depth of Field (I)

The height over which a sample can be clearly focused is called the Depth of Field. The SEM has a large depth of field which produces the images that appear 3-dimensional in nature.

Depth of field is improved by:

- Longer working distance
- Smaller objective apertures
- Lower magnifications

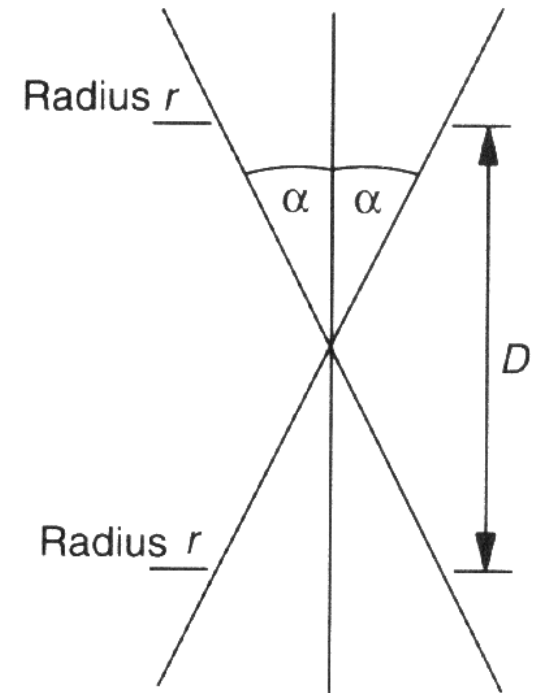
# Depth of Field (II)

A beam having a semi-angle of convergence  $a$  will converge from a radius  $r$  to a focus and diverge again in a vertical distance  $D$ . For small  $a$  :

$$D = 2r / a$$

In a current high-resolution CRT's (spot size = 0.1 mm = 100 microns), the focusing becomes objectionable when two pixels are fully overlapped, where the pixel size on the specimen is  $0.1/M$  mm, where  $M$  is the magnification, which gives us the practical expression for the depth of focus/field:

$$D \gg 0.2 / aM \text{ (mm)}$$



# Depth of Field (III)

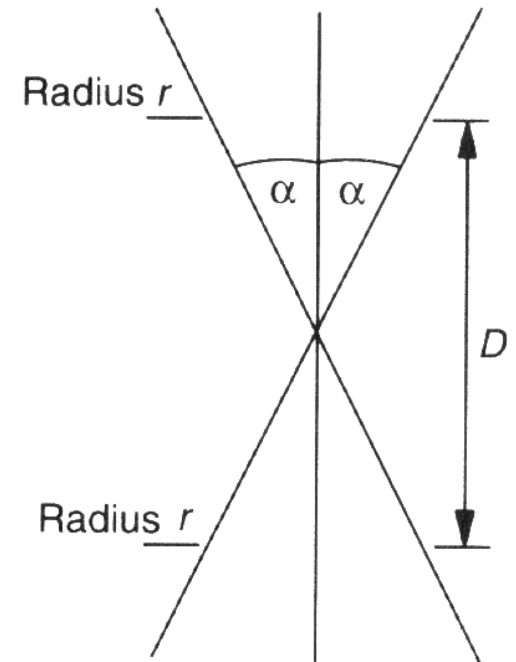
The angle  $\alpha$  is determined by:

$$\alpha = R/WD$$

$R$  = radius of the aperture used, and  $WD$  is the working distance of the aperture from the specimen.

Table 34.2 Depth of field in  $\mu\text{m}$  with 10-mm working distance

Magnification	With 100- $\mu\text{m}$ aperture	With 400- $\mu\text{m}$ aperture
100	400	100
1000	40	10
10000	4	1
100000	0.4	0.1



# Depth of Field vs. Resolution

**As for the optical microscope, the depth of field and resolution have a reciprocal relationship:**

- Improving resolution in conventional SEM's leads to a smaller depth of field**
- While increasing depth of field decreases resolution, useful for each particular sample.**

# 3.0 Instrumentation



**JEOL 6700F Ultra High Resolution  
Scanning Electron Microscope**

# 3.1 Sample Preparation

## Sample Coating

**Q: Why ?**

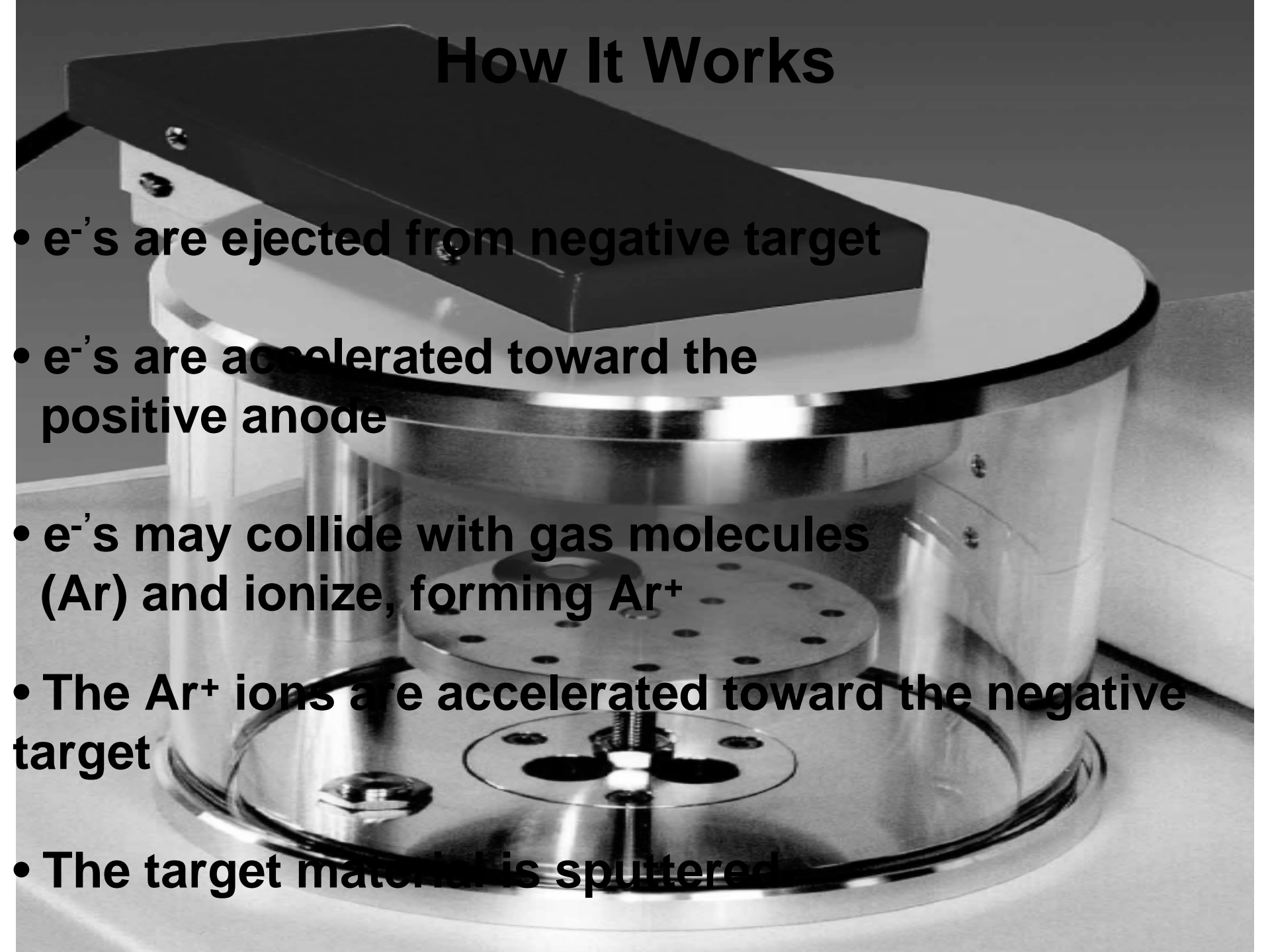
**A: Charging:**

- **Deflection of SE's**
- **Increased emission of SE's in cracks**
- **Periodic SE bursts**
- **Beam deflection**

**Solutions:**

- **Sputter coating with C, Cr, or Au-Pd**
- **Carbon tape, carbon paint, In foil**

# How It Works

- $e^-$ 's are ejected from negative target
  - $e^-$ 's are accelerated toward the positive anode
  - $e^-$ 's may collide with gas molecules (Ar) and ionize, forming  $Ar^+$
  - The  $Ar^+$  ions are accelerated toward the negative target
  - The target material is sputtered
- 

# Sputtering: A Physical Process ?

Vacuum Chamber

Anode

Negative target

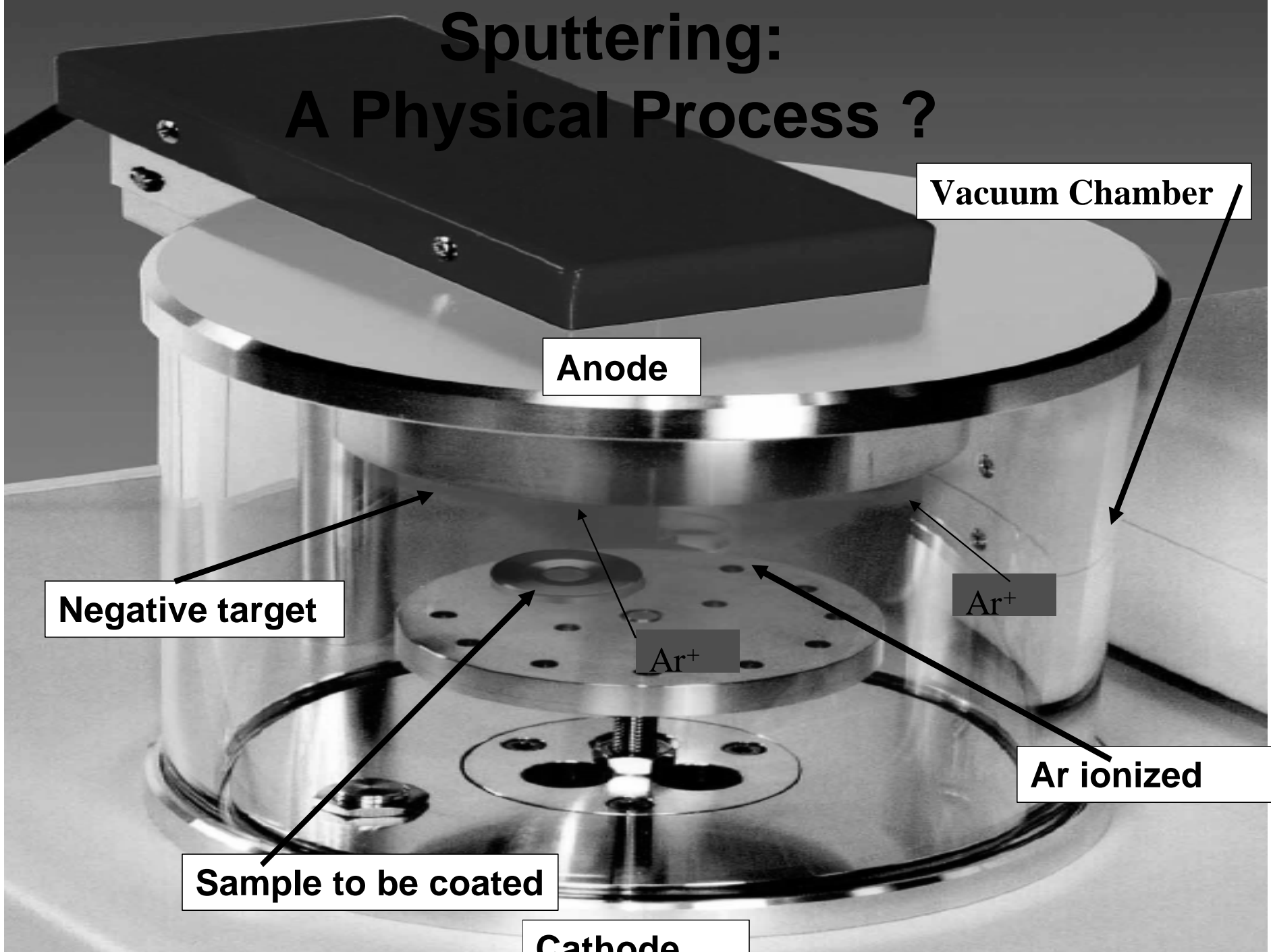
Ar<sup>+</sup>

Ar<sup>+</sup>

Ar ionized

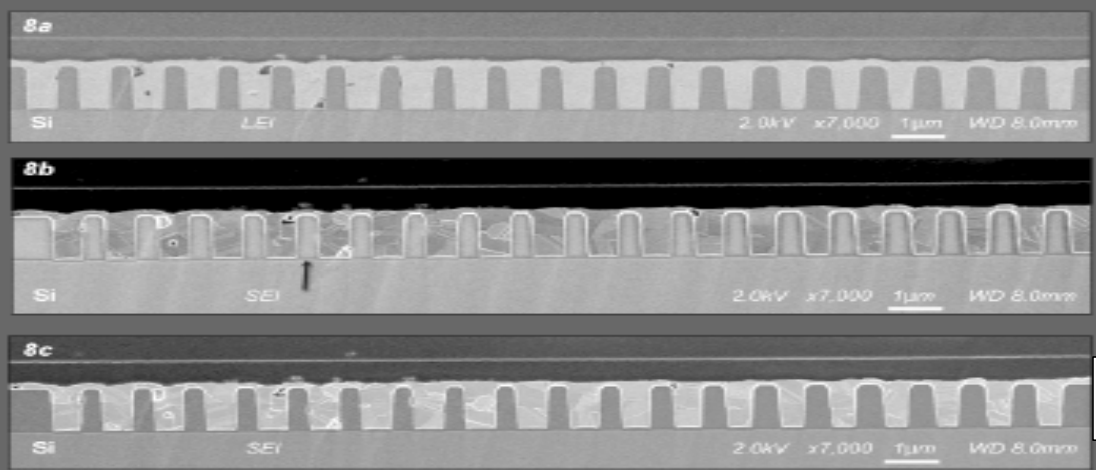
Sample to be coated

Cathode





# Etching and Coating



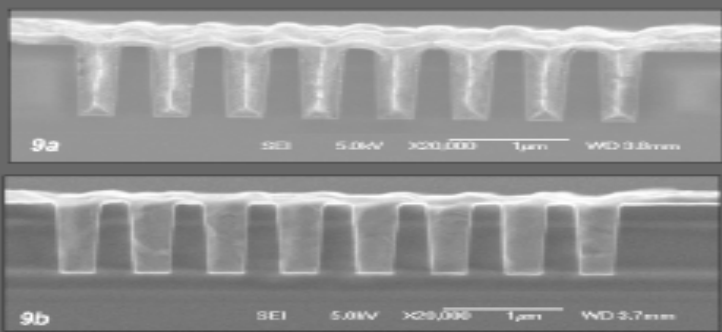
**8.** A sequence of SEM images taken from Cu contacts on a Si substrate after a variety of surface treatments in the PECS.

Figure 8a is taken after ion polishing and before etching. The specimen was coated with 20 Pt to prevent charging.

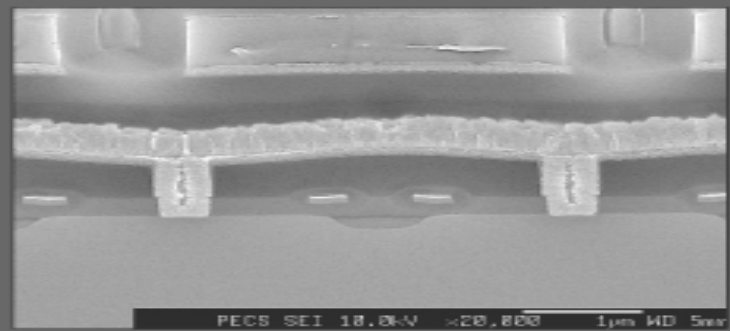
Figure 8b, as Figure 8a, but after ion beam etching only and no conductive coating. It is seen that there are some charging effects on SiO<sub>2</sub> areas between Cu contacts (see arrow).

Figure 8c, as Figure 8b, but after coating with -20 coating. Note charging has been prevented.

## Gatan PECS Model #682



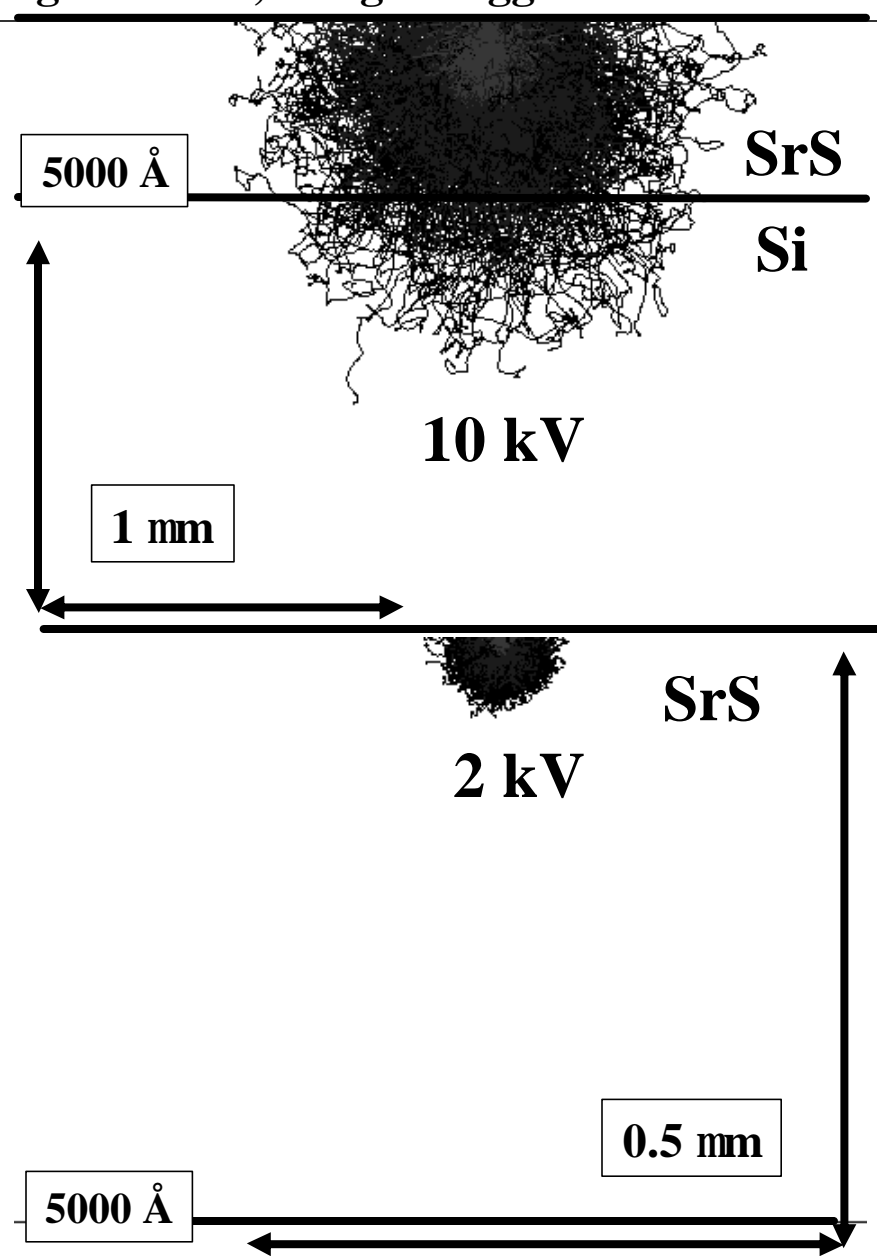
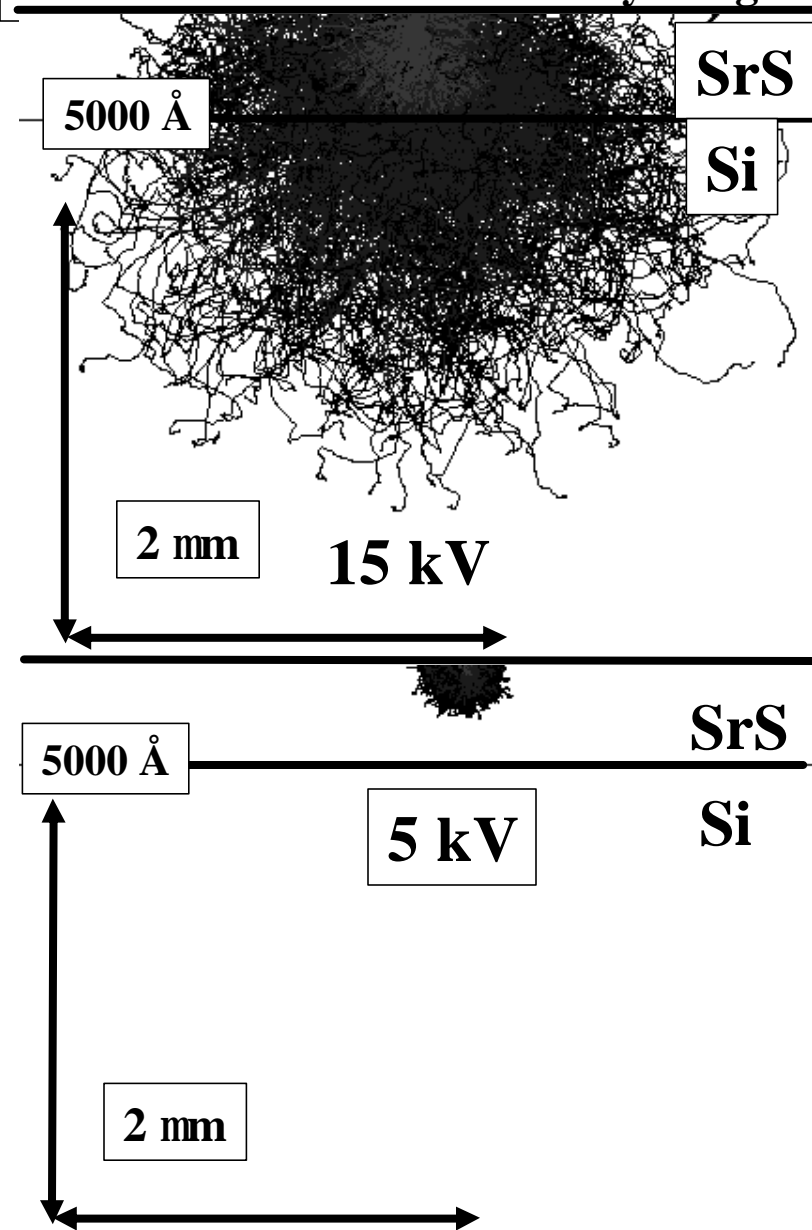
**9.** Before (9a) and after (9b) etching SEM images of cleaved Cu contacts on Si substrate. Note that as all plastically deformed Cu areas are removed, more details of Cu contacts along with barrier layers are visible.



**10.** As cleaved IC device containing tungsten plugs and other features after etching and coating (Pt) in the PECS.

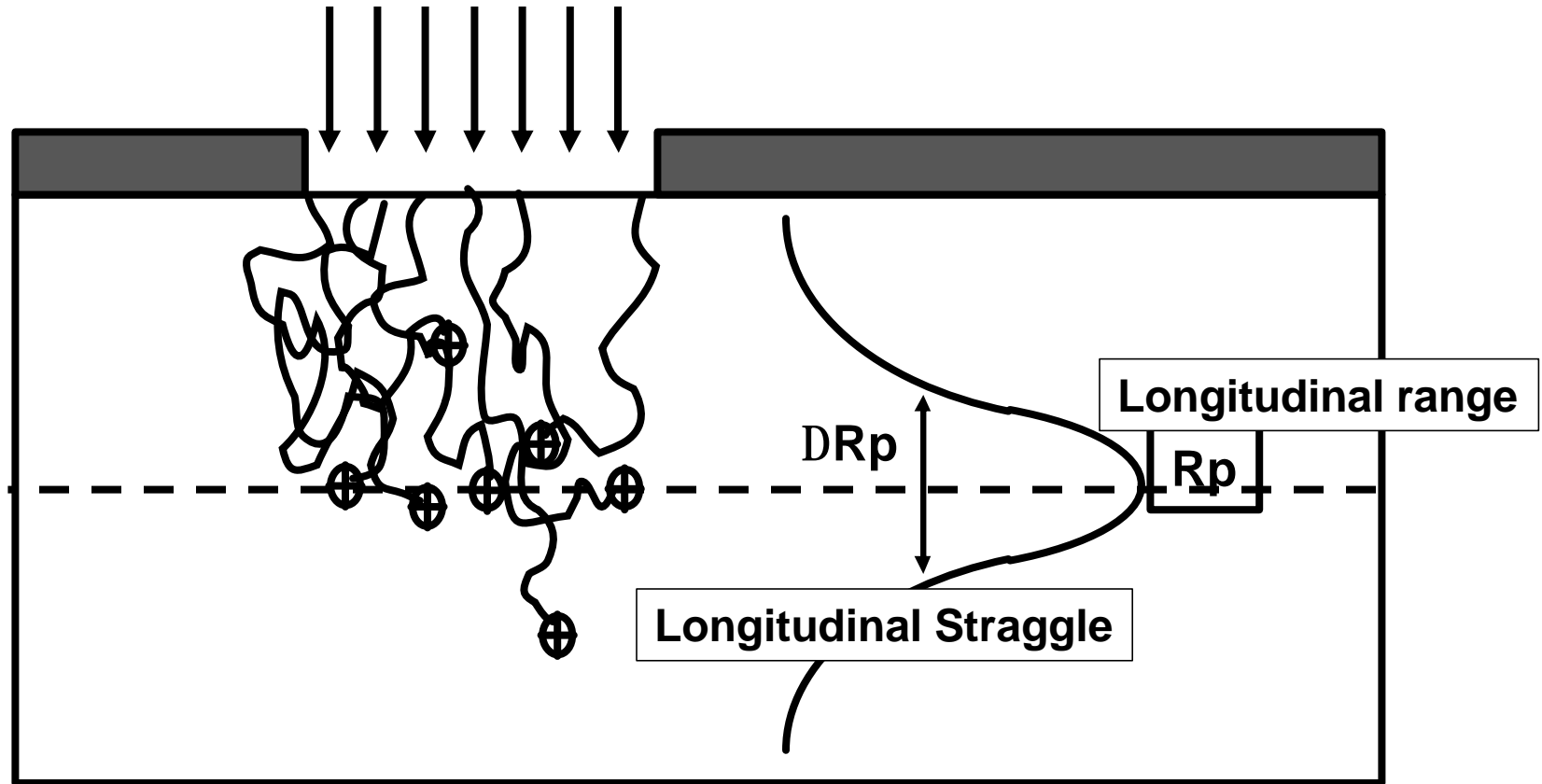
# Beam Interaction Simulations

Theoretical SrS density:  $3.7\text{g/cm}^3$ , Long. range:  $1062\text{ \AA}$ , Long. straggle:  $384\text{ \AA}$



# Random interactions with target atoms

Ion, dose, energy

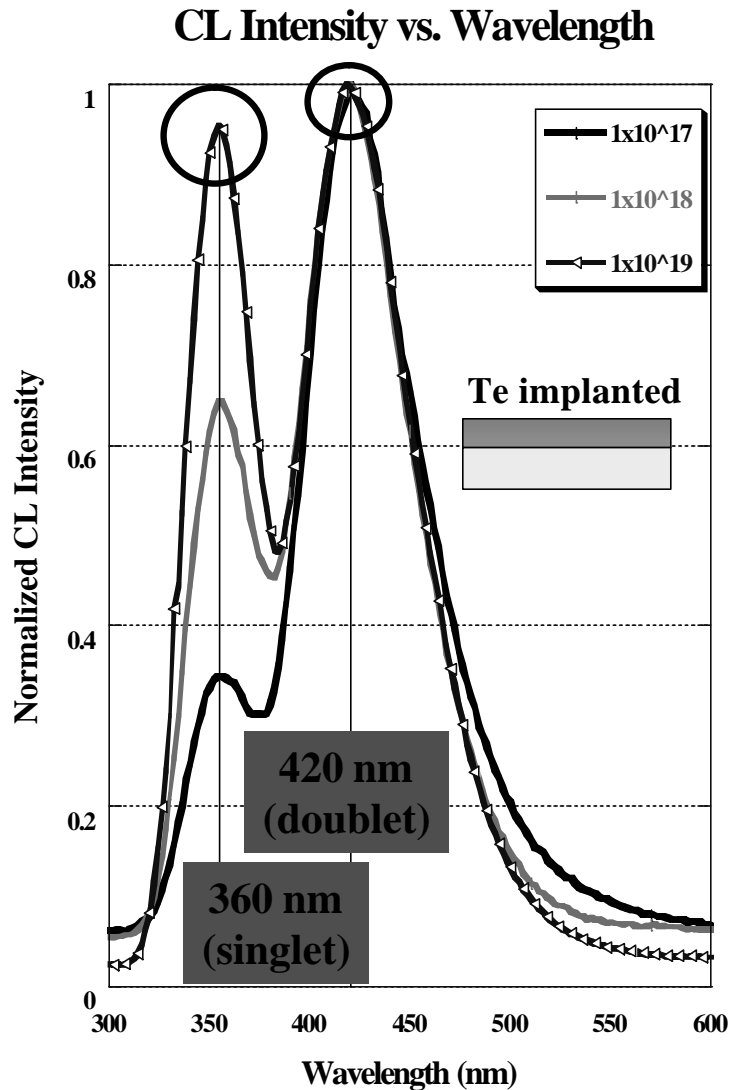


From LSS Theory:

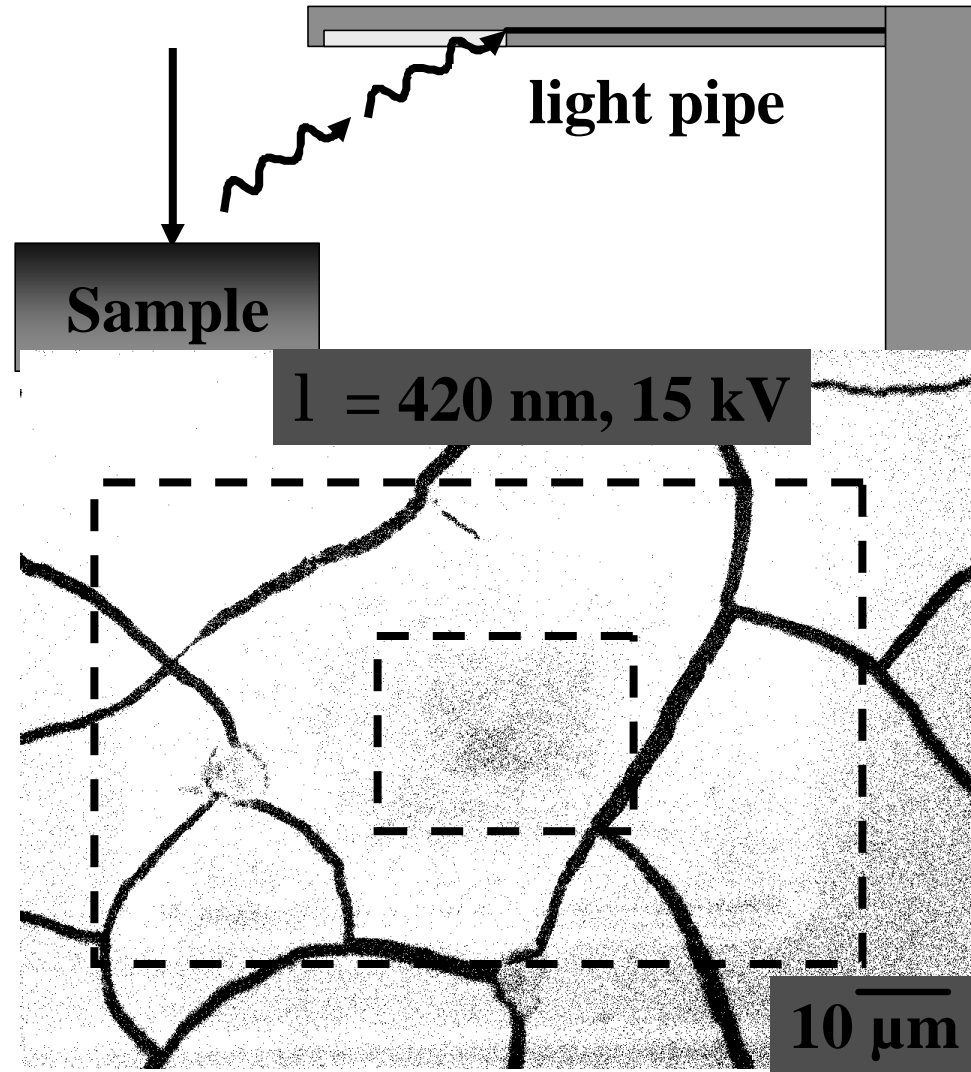
$$R_p \propto \frac{R}{1 + \frac{M_T}{3M_i}}$$

$$DR_p \propto \frac{2R_p}{3} \frac{\sqrt{M_i M_T}}{M_i + M_T}$$

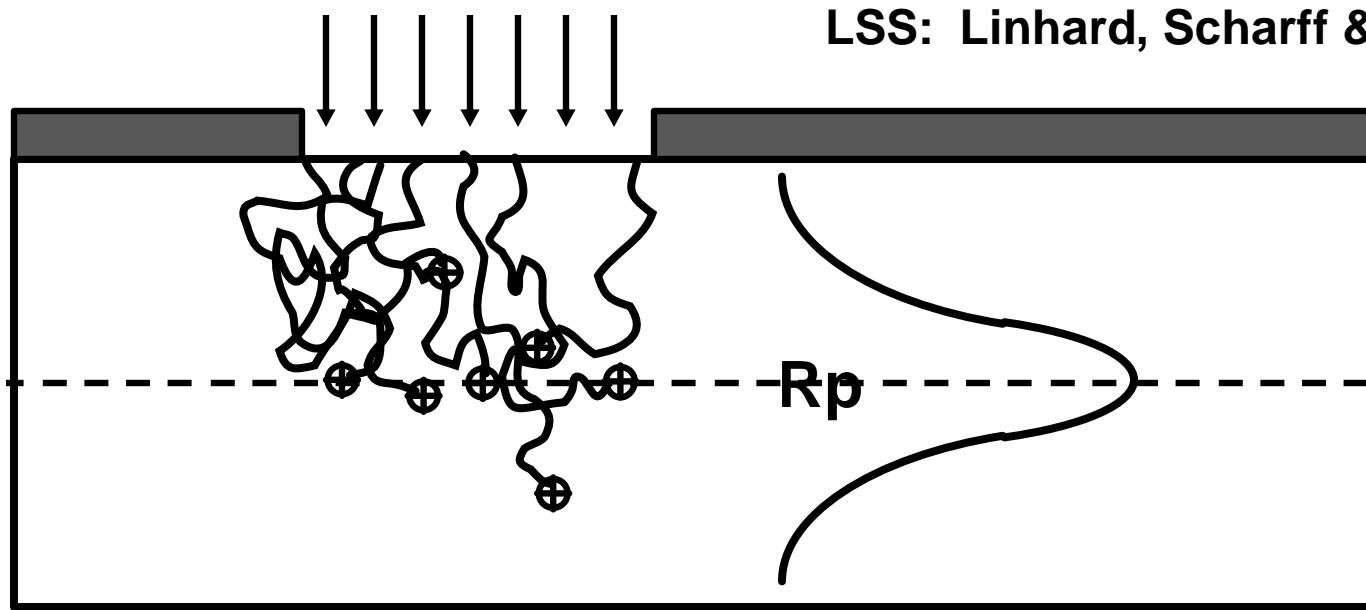
# Cathodoluminescence (CL)



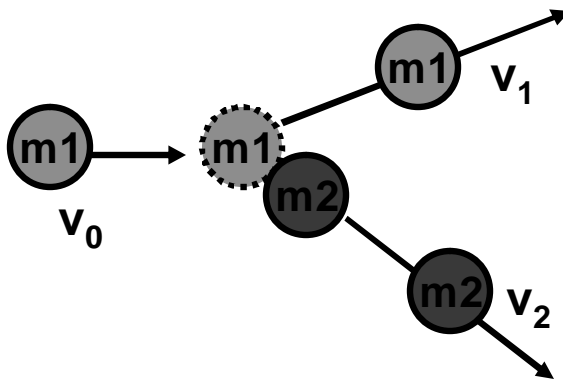
CL Imaging at 15 kV, 10kX - 1kX  
Sample:  $1 \times 10^{19} \text{cm}^{-3}$



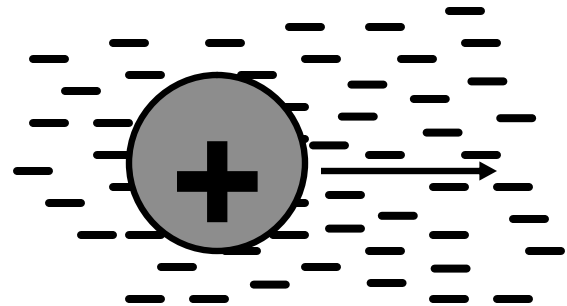
# LSS Theory of Ion Stopping



Nuclear Stopping: Coulombic Scattering



Electronic Stopping



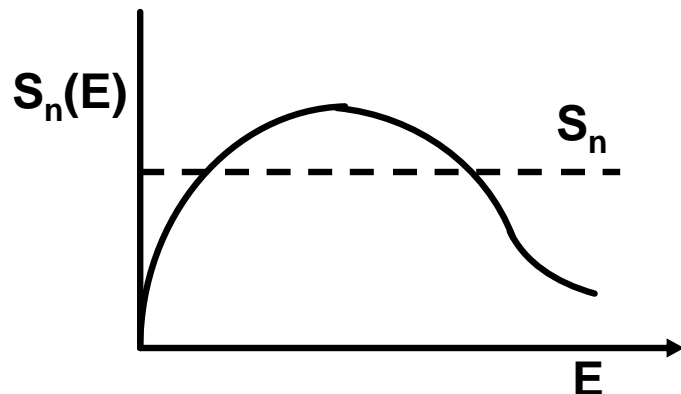
# Nuclear Stopping

Energy loss per distance traveled as a function of energy:

- depends on incident ion and target atom nuclear charge ( $Z = \#$  protons) and atomic masses ( $M$ )

$$S_n(E) = \frac{dE}{dx} \frac{1}{\rho_n}$$

Energy loss due to interactions with atomic nuclei is basically a decreasing function of energy. At high kinetic energy (ion velocity) there is a very short interaction time for any absorption of energy by target atoms.



Approximate nuclear stopping near max of  $S_n(E)$ :

$$\frac{dE}{dx(n)} = S_n = N \frac{p^2}{2} e^2 a \frac{Z_1 Z_2 M_1}{M_1 + M_2}$$

where

$N$  is the atomic density (atoms/Volume)

$a \sim 1.4 \times 10^{-2}$  nm, subscripts 1 and 2 refer to ion and target respectively,  $Z$  is atomic number and  $M$  is mass number

# Electronic Stopping

Due to interactions with electrons in the target material  
- like a drag force that is proportional to the ion velocity

Ion velocity  $\propto (\text{Energy})^{1/2}$

$$S_e(E) = \frac{dE}{dx} = \frac{2\pi Z^2 e^4 n_e}{m_e v} \mu E^{1/2} = k_e E^{1/2}$$

where  $k_e$  is relatively independent of the incident ion

For silicon:  $k_e \sim 10^7 (\text{eV})^{1/2}/\text{cm}$

Total rate of energy loss =  $dE/dx = S_n(E) + S_e(E)$

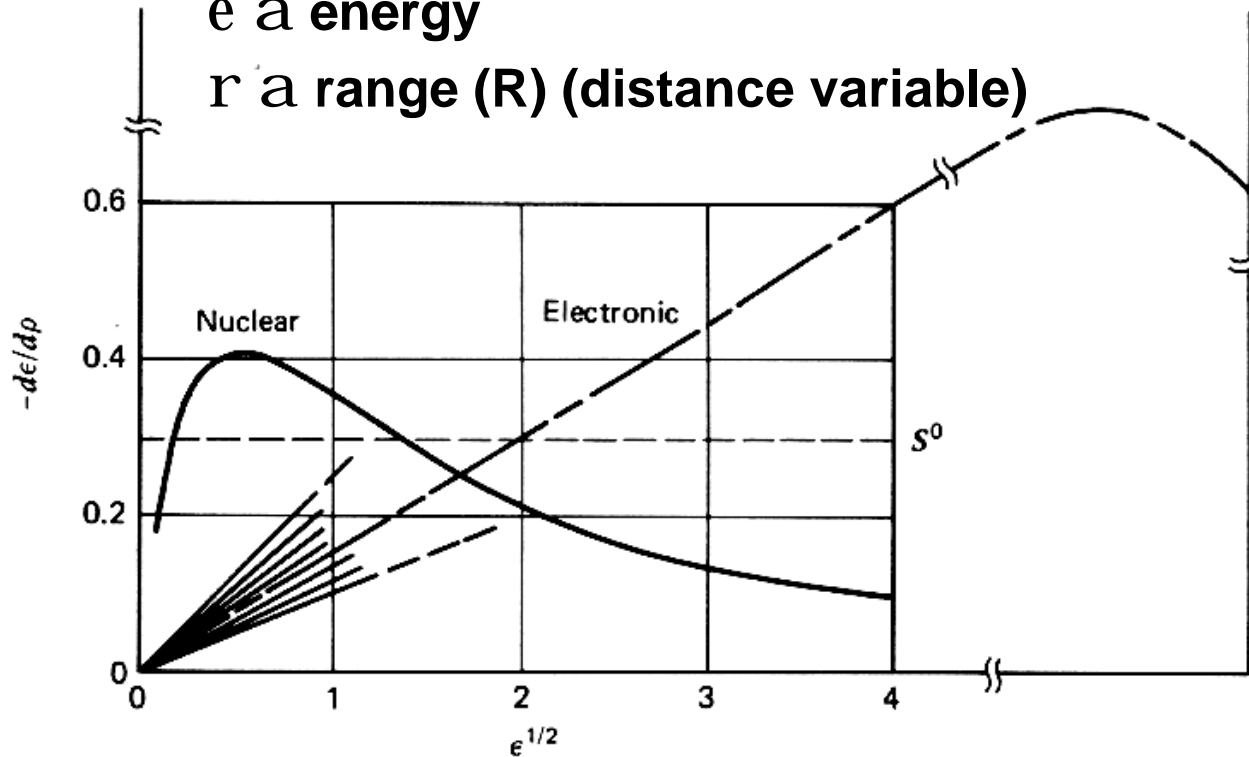
# LSS Calculations

$e$  &  $r$  are dimensionless parameters

$e$  a energy

$r$  a range ( $R$ ) (distance variable)

“like”  
 $dE/dx$



**Fig. 6.8** Nuclear and electronic stopping power curves. From J. W. Mayer, L. Eriksson, and J. A. Davies, *Ion Implantation in Semiconductors* [7], 1970. Used with permission of Academic Press, Inc.



# Nuclear & Electronic Stopping

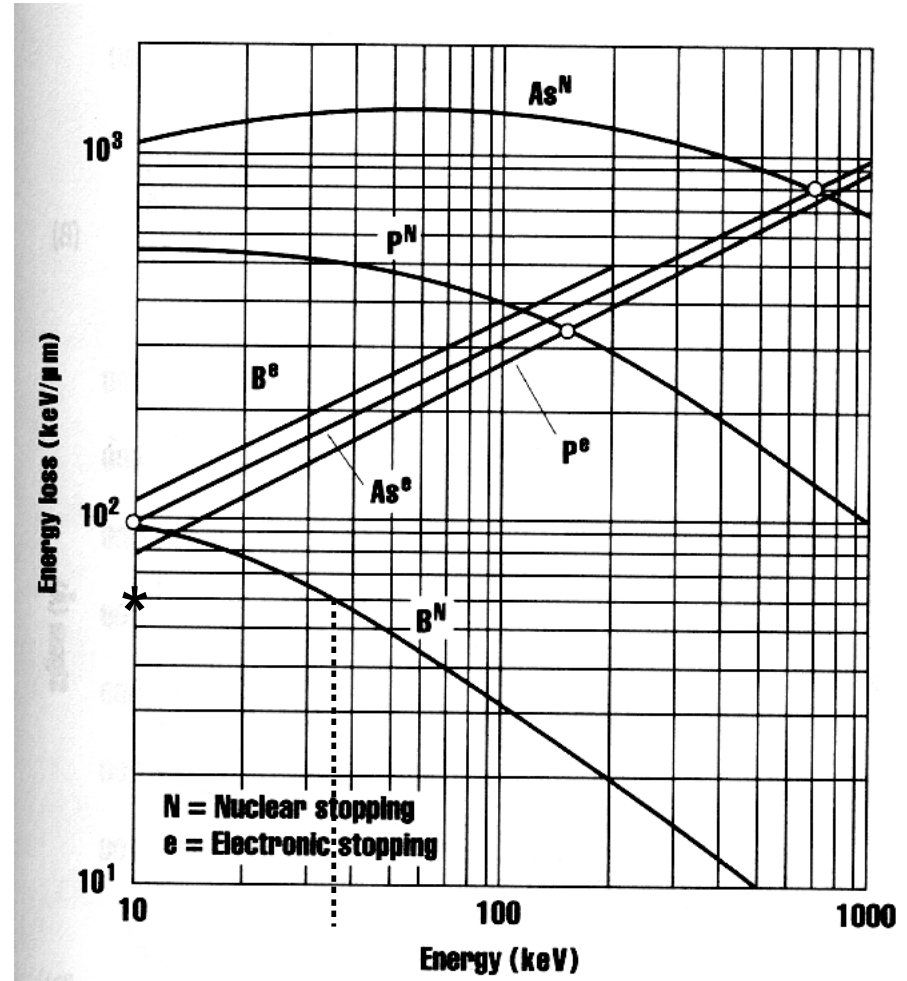
Electronic stopping dominates:

Light ions and high energies

Nuclear stopping dominates:

Heavy ions and low energies

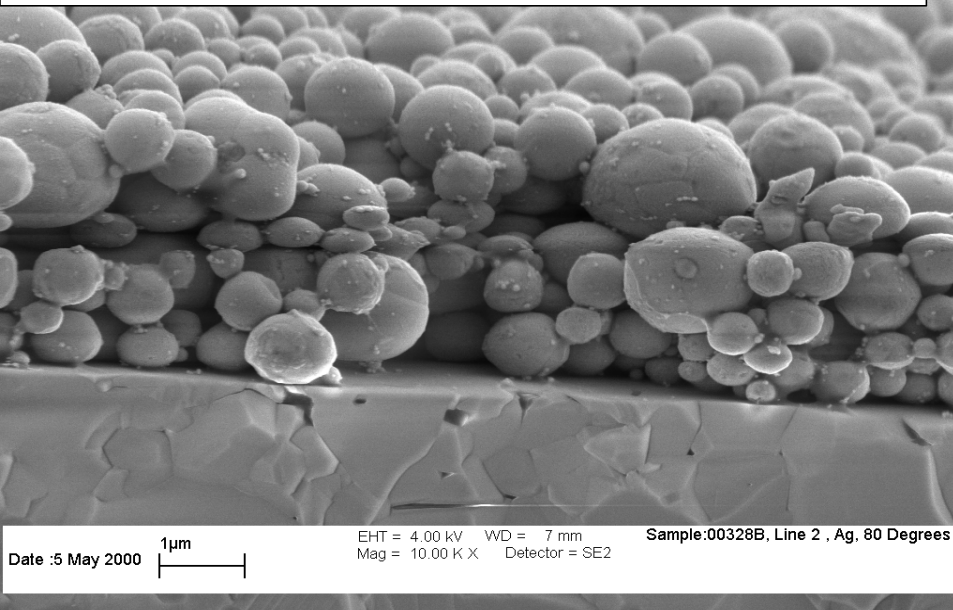
Implant damage occurs due to nuclear interactions. The extent of damage depends on  $S_n(E)$ .



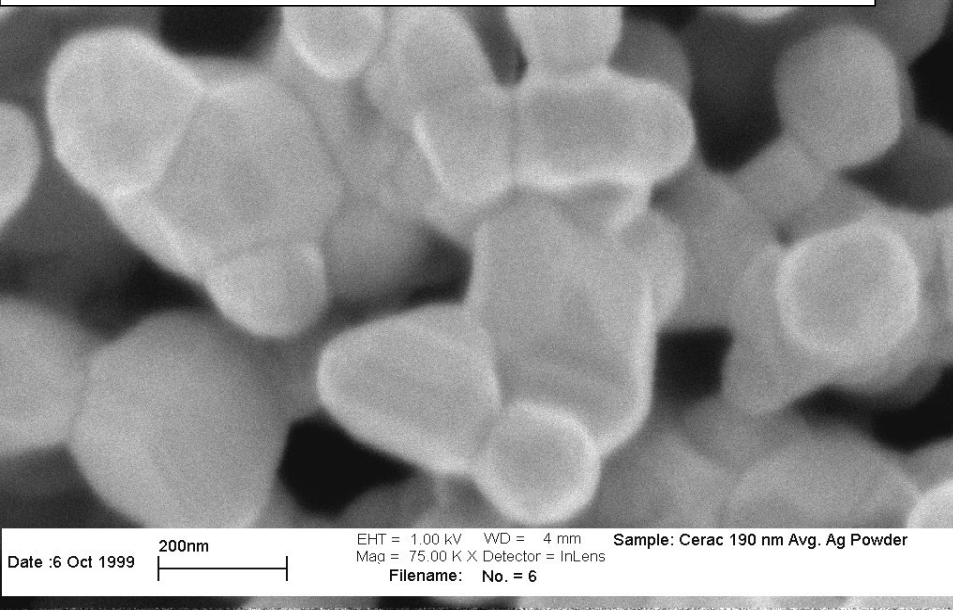
# Artifacts and Examples

- **Thermal damage**
  - **Sample temperature can rise up to 40°C**
- **Surface Etching**
- **Coating Film Adhesion**
  - **Not sticking, temperature sensitive**
- **Contamination**
- **Morphology Modification**

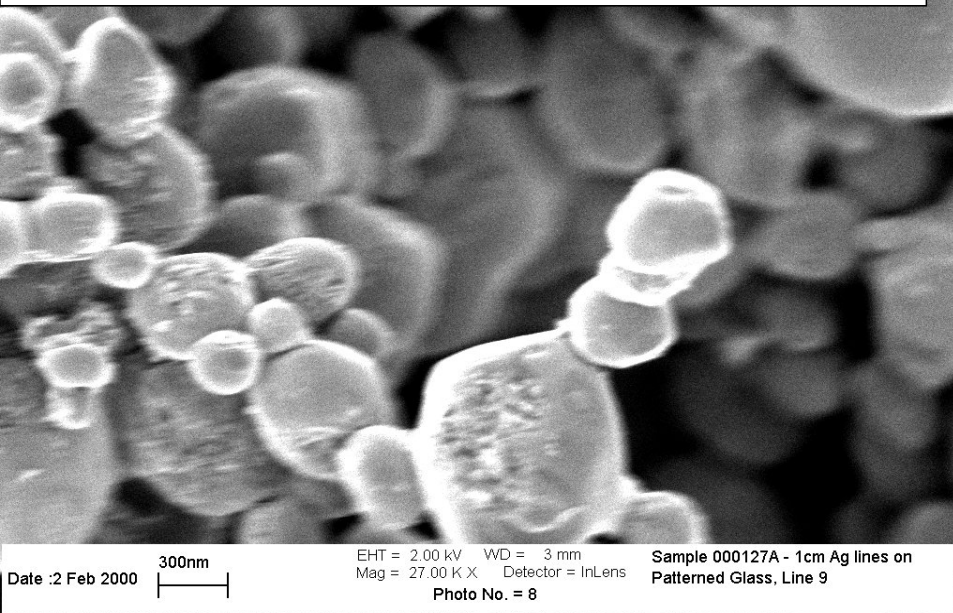
# Ag powders written by a laser on glass



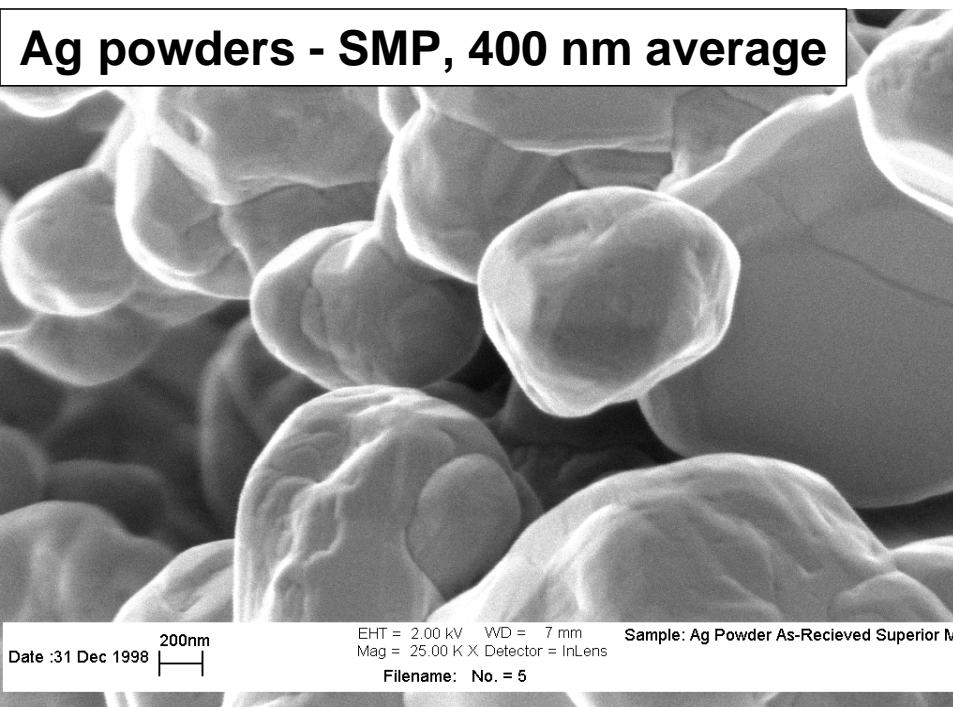
# Ag powders - Cerac, 190 nm average



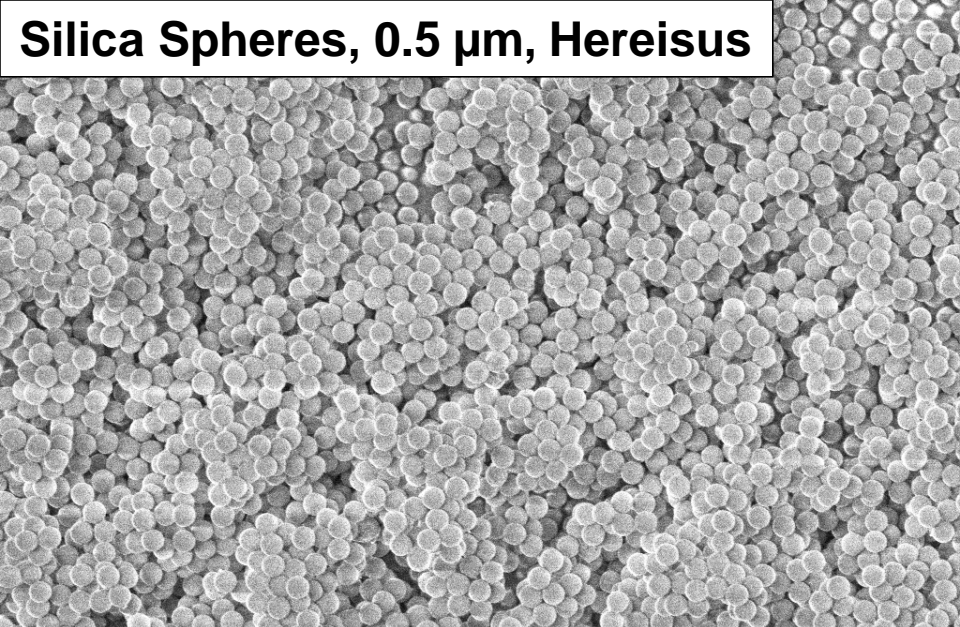
# Ag powders written by a laser on glass



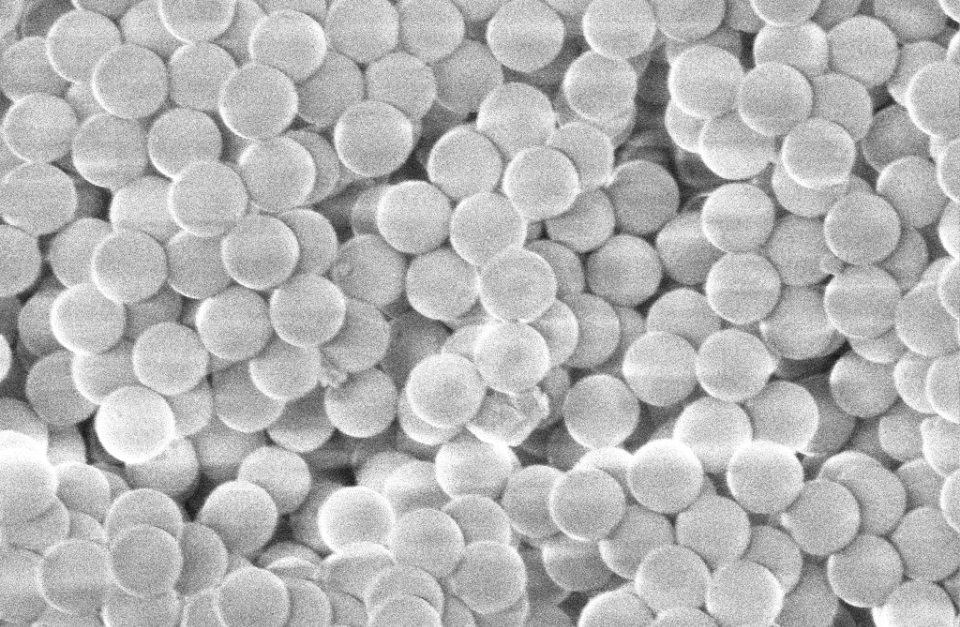
# Ag powders - SMP, 400 nm average



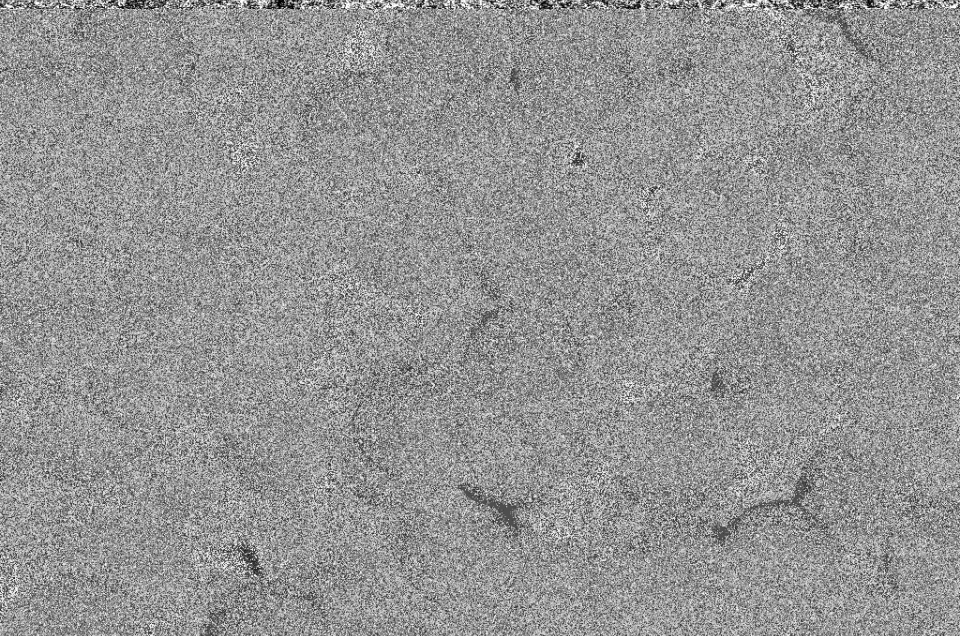
# Silica Spheres, 0.5 $\mu\text{m}$ , Hereisus



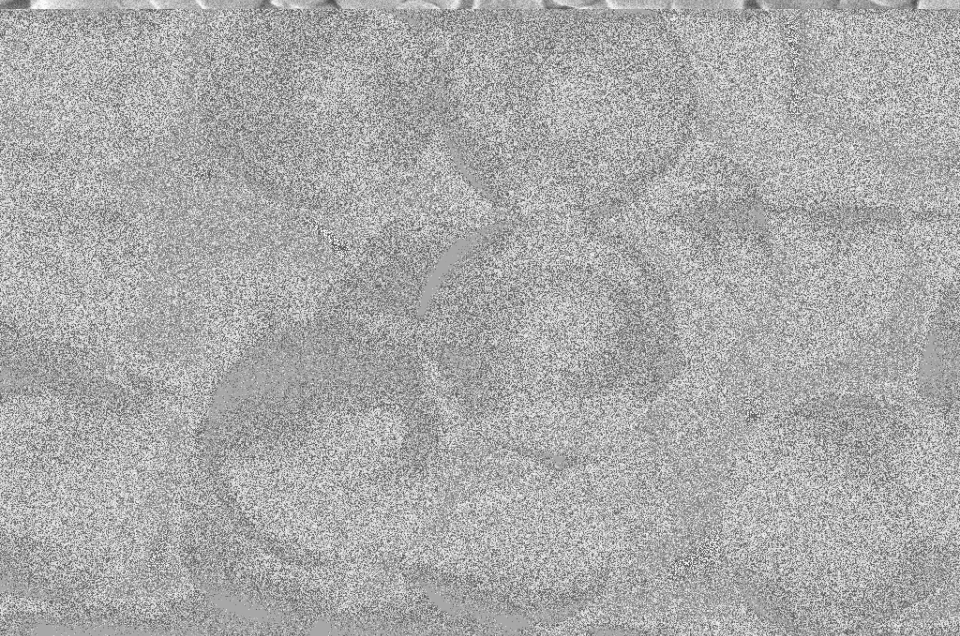
Date :17 Feb 1999  $1\mu\text{m}$  EHT = 2.00 kV WD = 4 mm Sample: Silica Her. 0.5 micron  
Mag = 5.00 K X Detector = InLens  
Filename: No. = 1



Date :17 Feb 1999  $1\mu\text{m}$  EHT = 2.00 kV WD = 4 mm Sample: Silica Her. 0.5 micron  
Mag = 15.00 K X Detector = InLens  
Filename: No. = 2

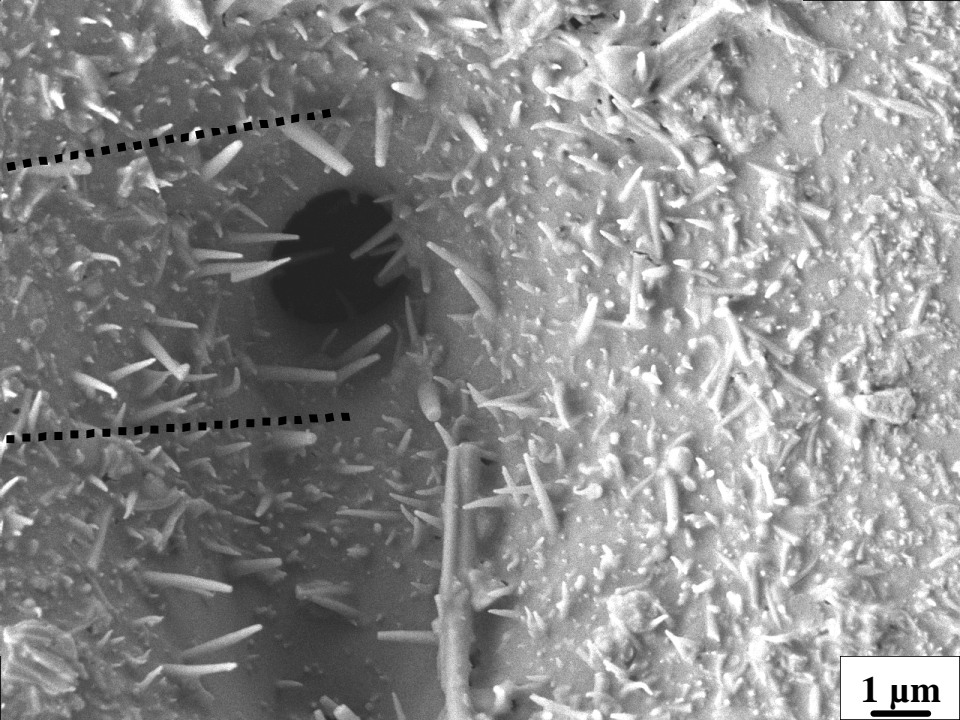
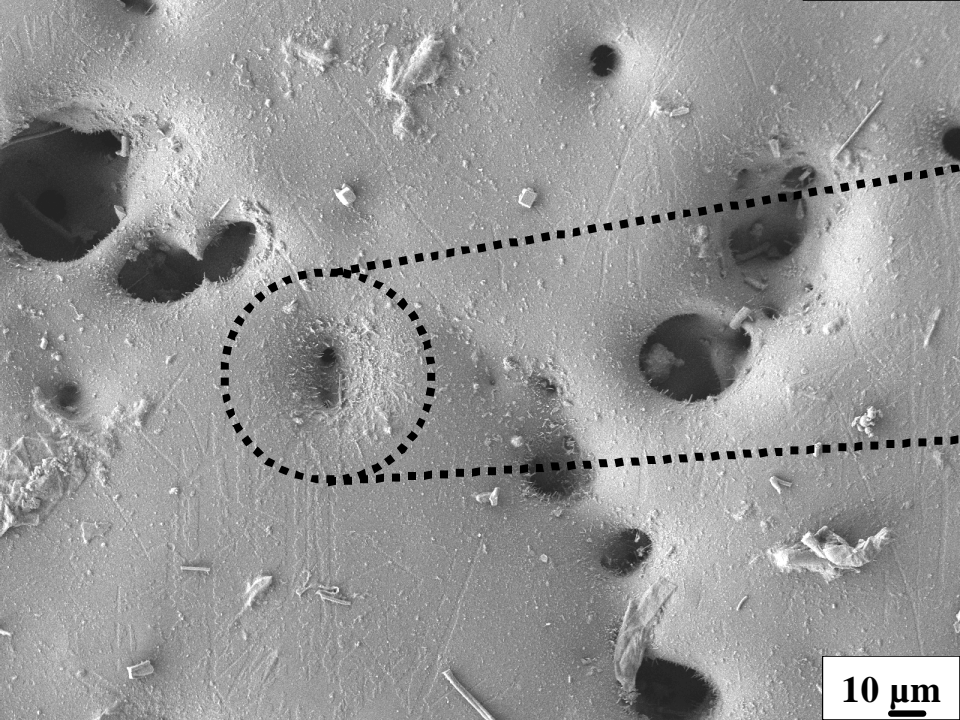
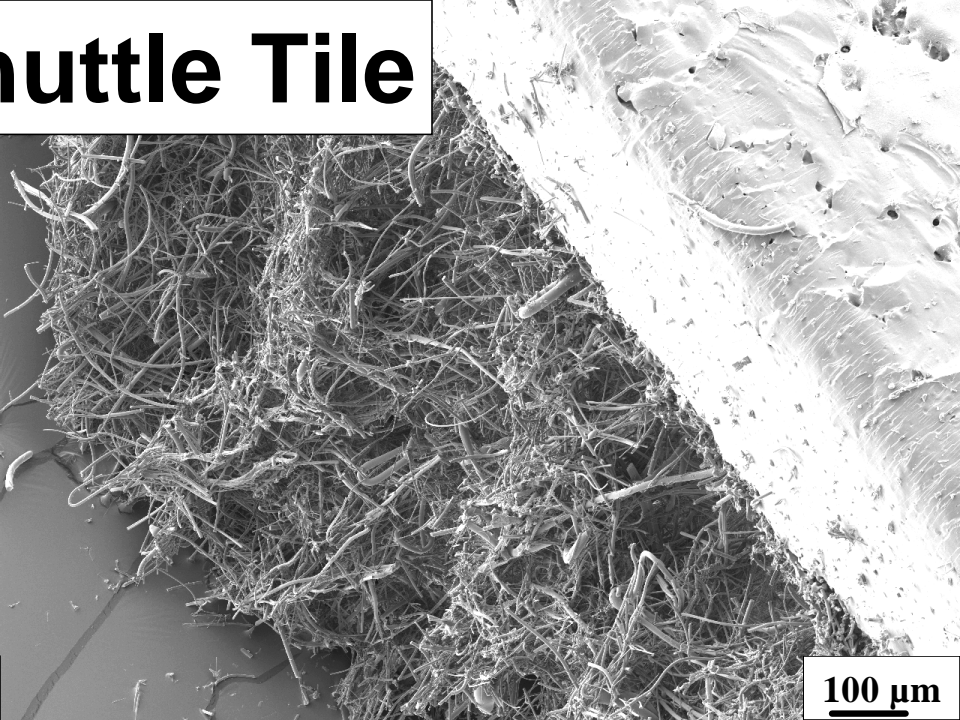
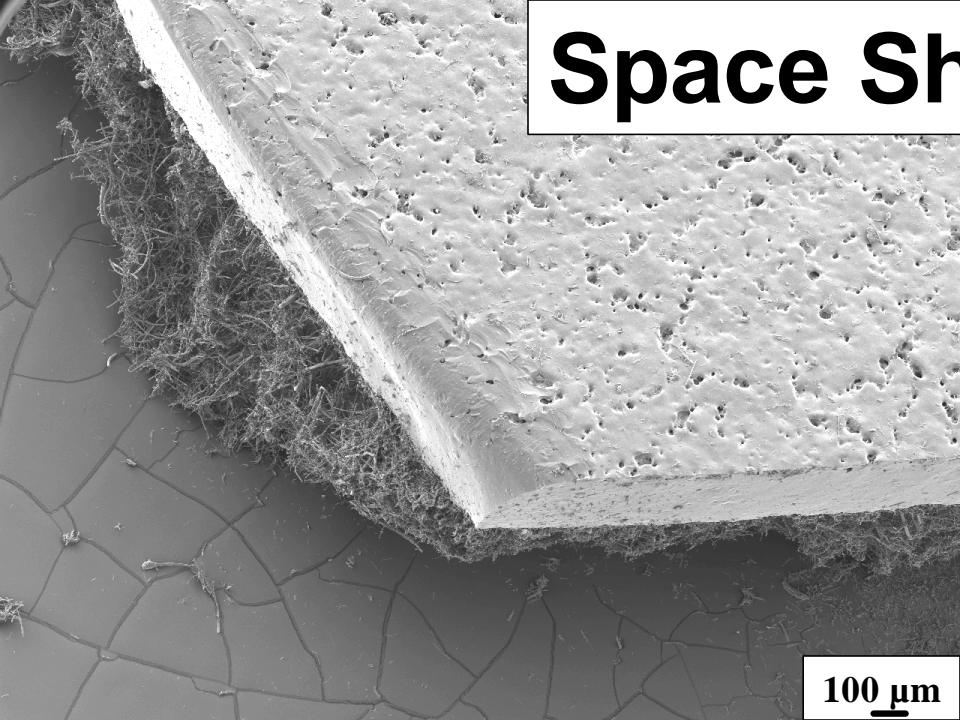


Date :17 Feb 1999  $200\text{nm}$  EHT = 2.00 kV WD = 4 mm Sample: Silica Her. 0.5 micron  
Mag = 25.00 K X Detector = InLens  
Filename: No. = 3

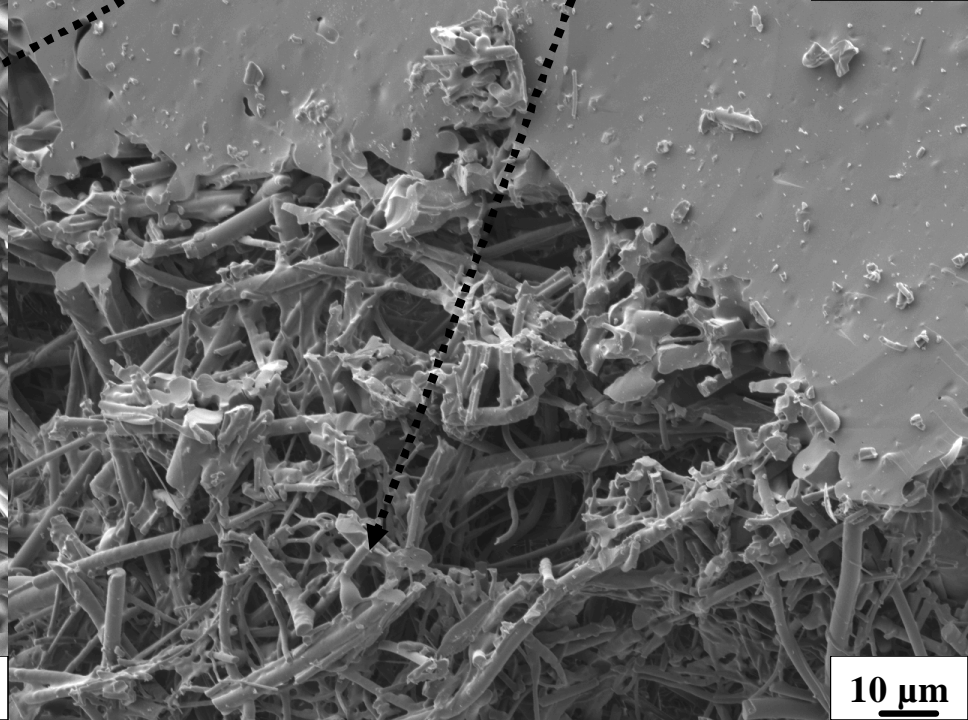
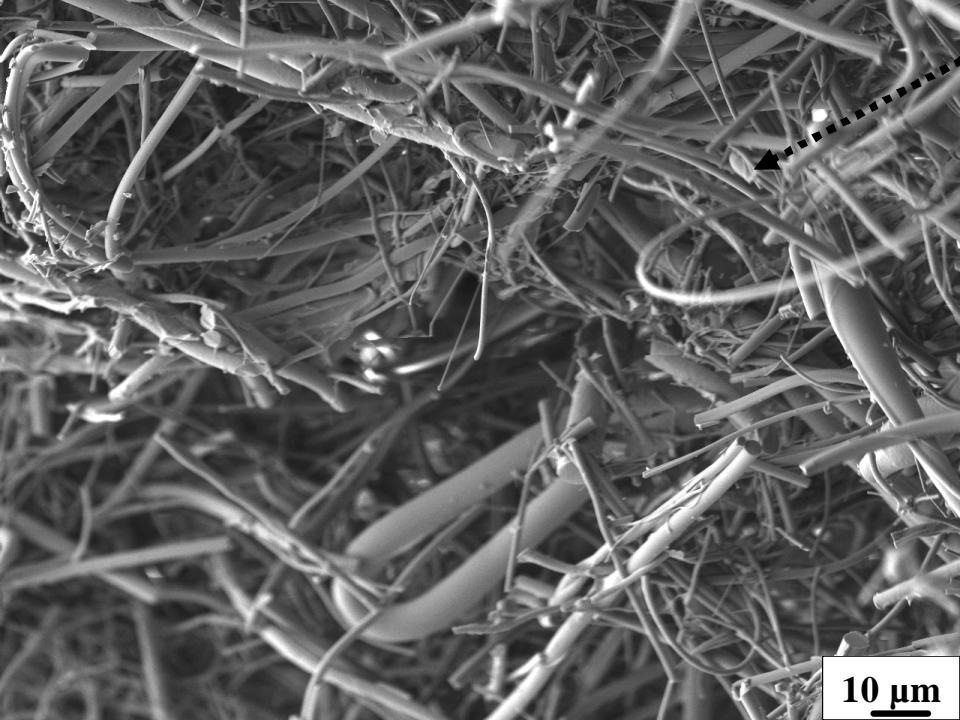
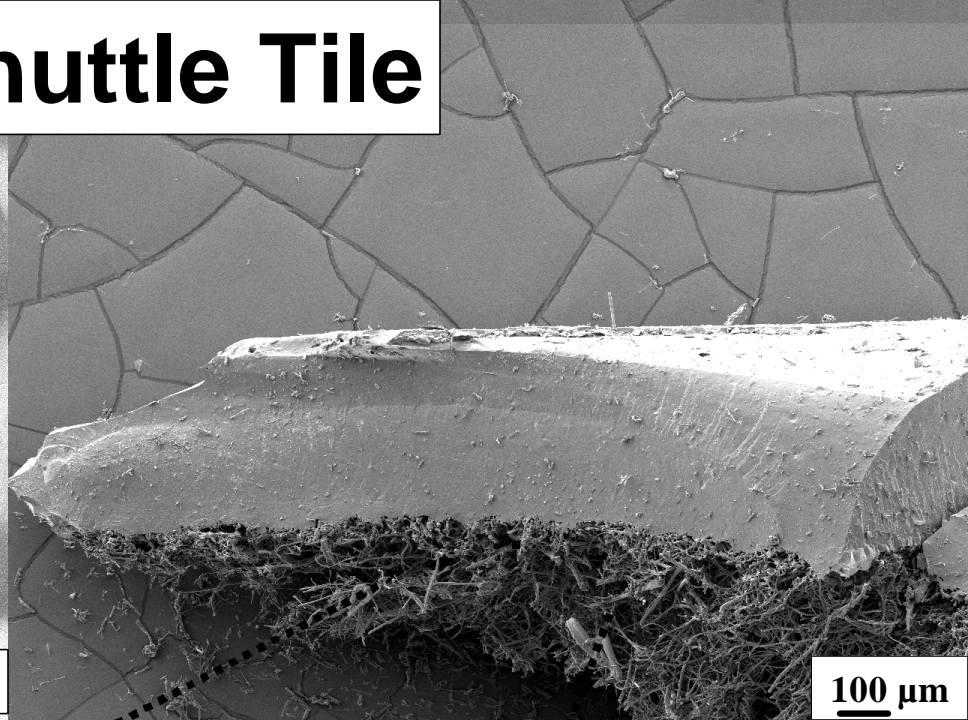
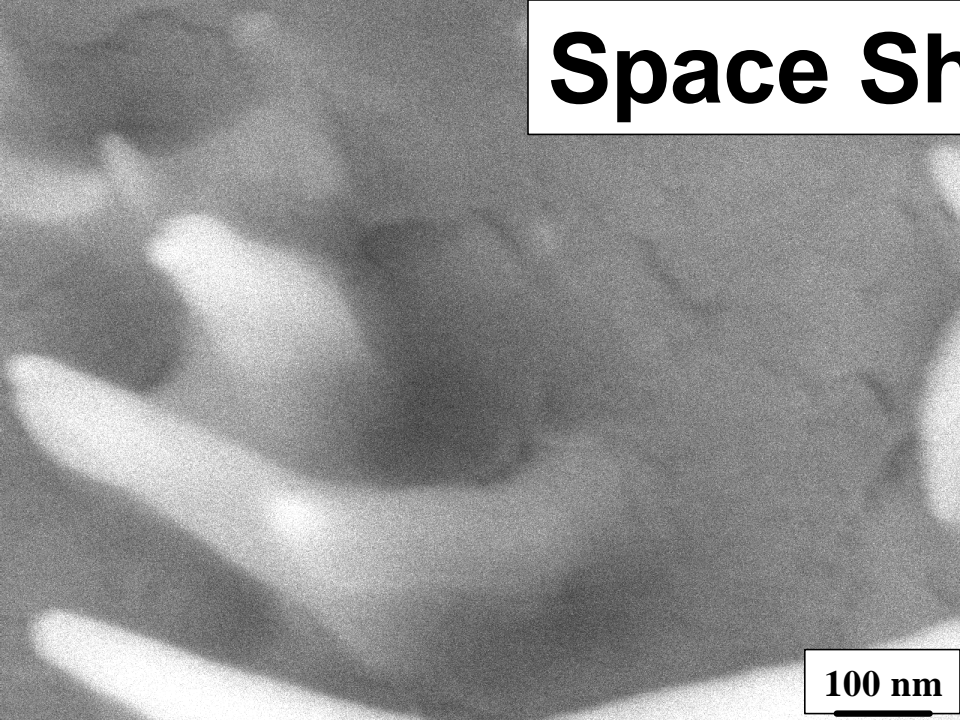


Date :17 Feb 1999  $200\text{nm}$  EHT = 2.00 kV WD = 4 mm Sample: Silica Her. 0.5 micron  
Mag = 50.00 K X Detector = InLens  
Filename: No. = 4

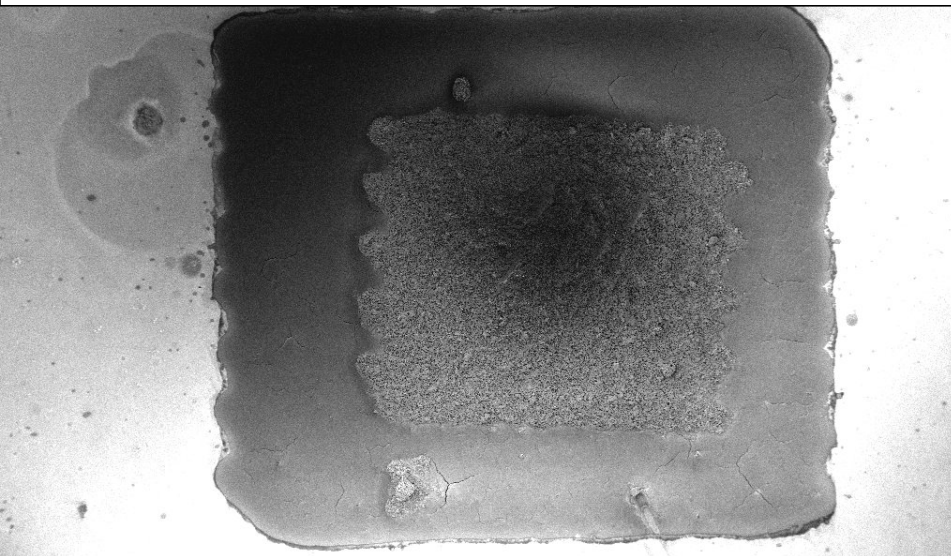
# Space Shuttle Tile



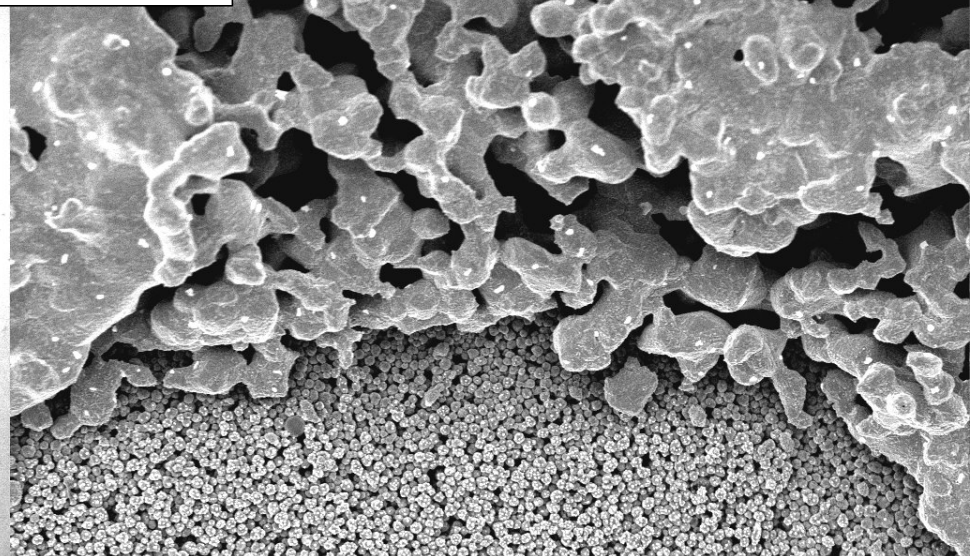
# Space Shuttle Tile



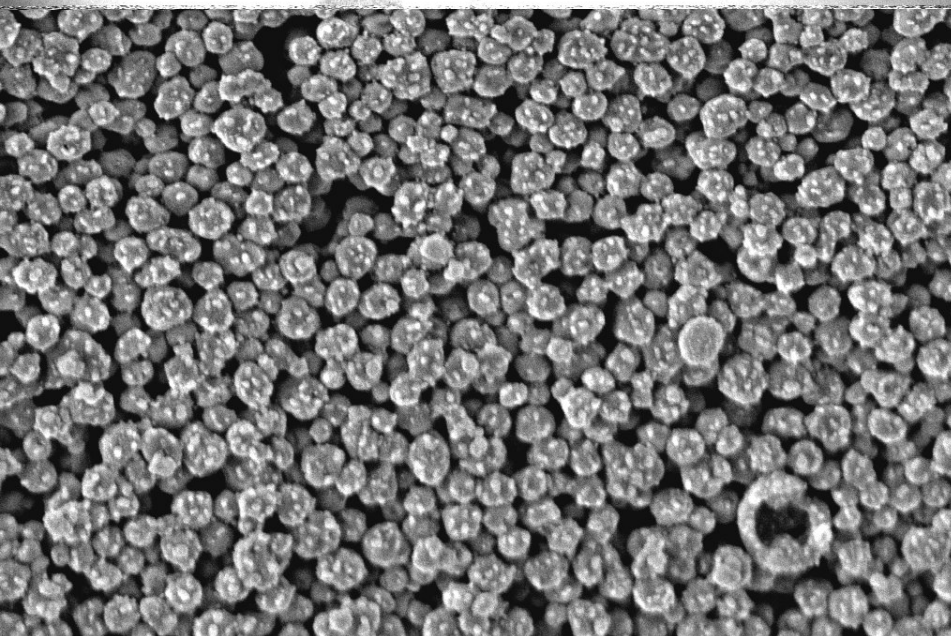
# BTO and Ag powders written by a laser on Kapton



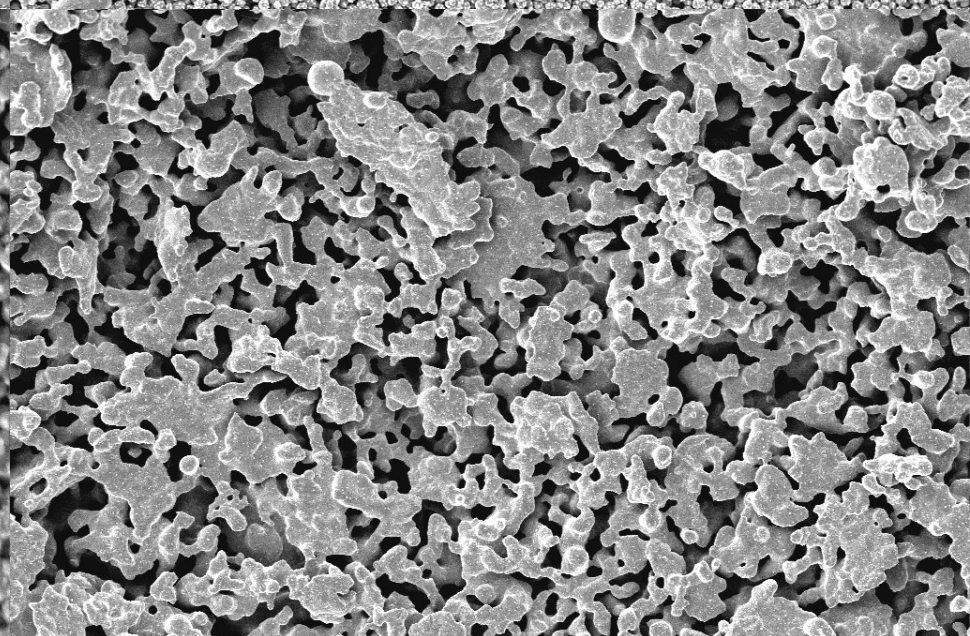
Date : 19 Apr 2000    100µm    EHT = 2.00 kV    WD = 3 mm    Sample: 0003-31B, PP CAP  
Mag = 150 X    Detector = InLens  
Photo No. = 2



Date : 19 Apr 2000    1µm    EHT = 2.00 kV    WD = 3 mm    Sample: 0003-31B, PP CAP, AREA 2  
Mag = 15.00 K X    Detector = InLens  
Photo No. = 5



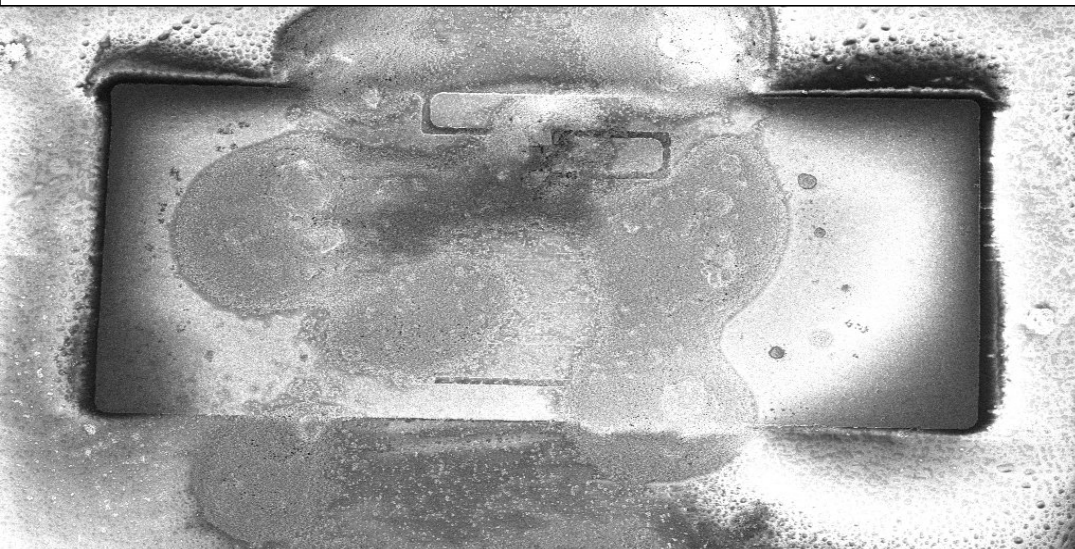
Date : 19 Apr 2000    200nm    EHT = 2.00 kV    WD = 3 mm    Sample: 0003-31B, PP CAP, AREA 2  
Mag = 50.00 K X    Detector = InLens    BTO W/WHIT DOTS  
Photo No. = 6



Date : 19 Apr 2000    2µm    EHT = 2.00 kV    WD = 3 mm    Sample: 0003-31B, PP CAP, AREA 3  
Mag = 5.00 K X    Detector = InLens    AG HAT  
Photo No. = 7

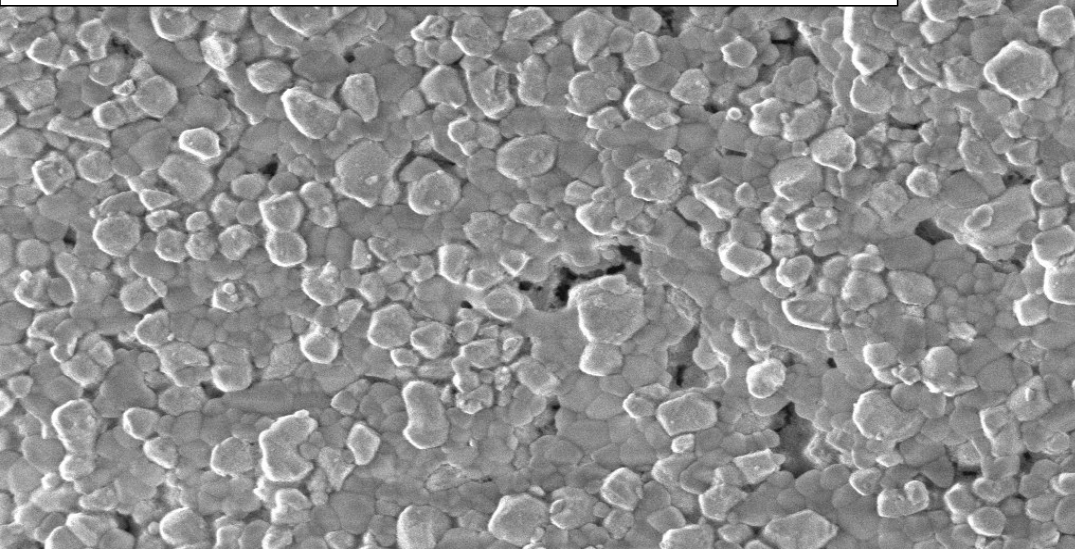


# BTO and Ag powders written by a laser on Kapton



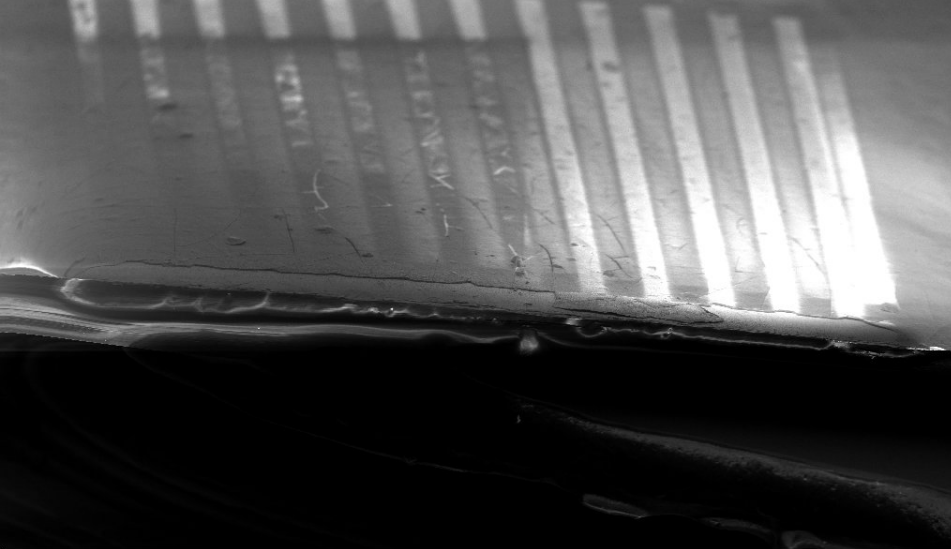
Date :26 May 1999 | 20µm | EHT = 2.00 kV WD = 2 mm Sample: Interdig Prico  
Mag = 304 X Detector = InLens  
Filename: No. = 1

# BTO sol gel matrix with BTO powders

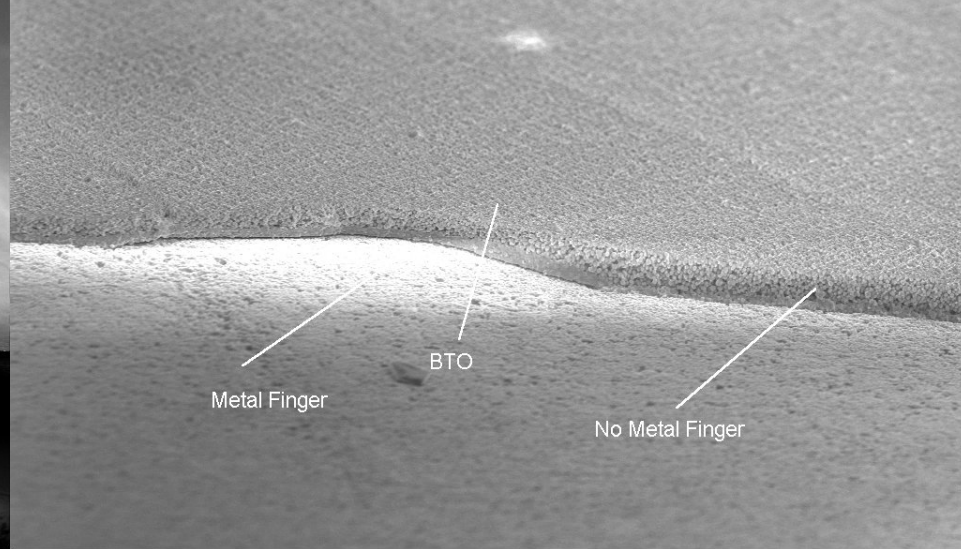


Date :26 May 1999 | 1µm | EHT = 2.00 kV WD = 2 mm Sample: Interdig Prico  
Mag = 10.00 K X Detector = InLens  
Filename: No. = 3

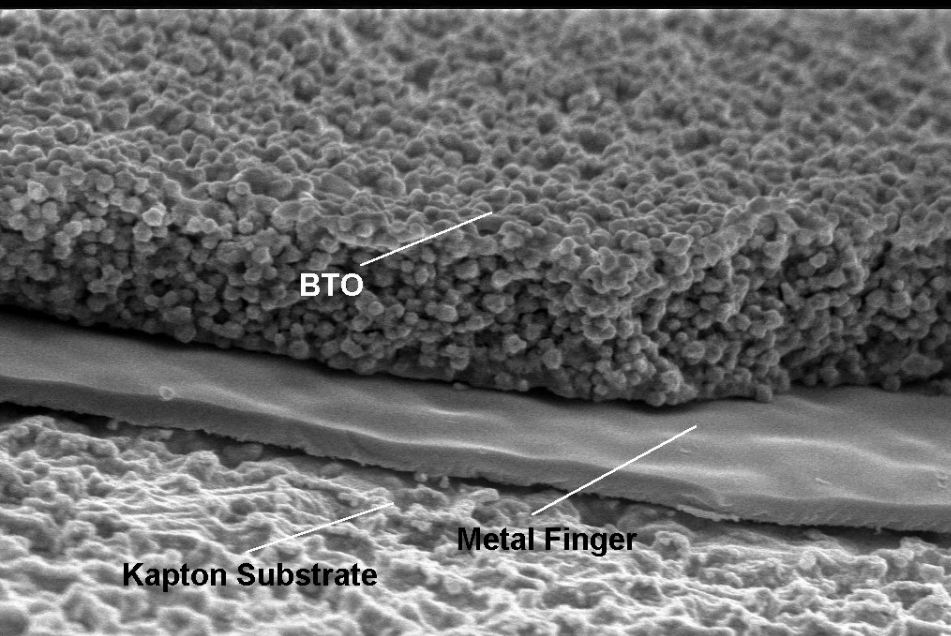
# BTO and Ag powders written by a laser on Kapton



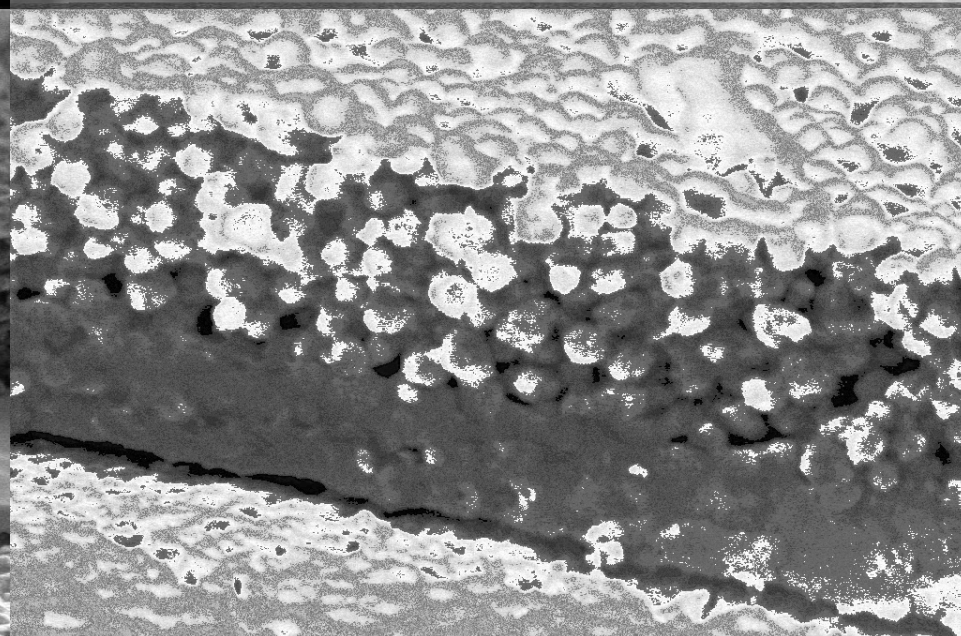
Date :10 May 2000 | 100µm | EHT = 4.00 kV WD = 4 mm Sample: 000508C, BTO/Kapton  
Mag = 125 X Detector = SE2 ID Caps, PAD 1, SMP QCR081032E



Date :10 May 2000 | 2µm | EHT = 4.00 kV WD = 4 mm Sample: 000508C, BTO/Kapton  
Mag = 2.50 K X Detector = SE2 ID Caps, PAD 1, SMP QCR081032E



Date :10 May 2000 | 1µm | EHT = 4.00 kV WD = 5 mm Sample: 000508C, BTO/Kapton  
Mag = 10.00 K X Detector = SE2 ID Caps, PAD 1, SMP QCR081032E



Date :10 May 2000 | 300nm | EHT = 4.00 kV WD = 4 mm Sample: 000508C, BTO/Kapton  
Mag = 25.00 K X Detector = SE2 ID Caps, PAD 1, SMP QCR081032E

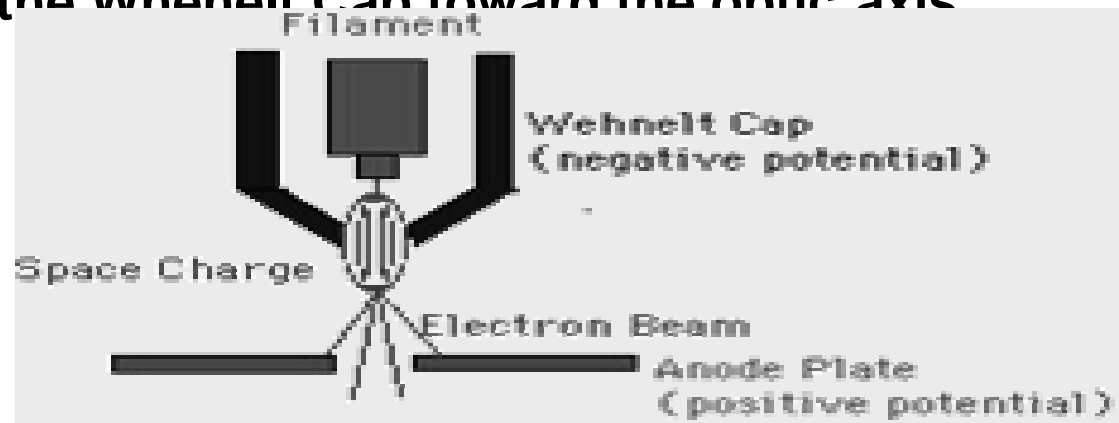
# Summary

- **The basic layout and operation of an SEM.**
- **The different imaging modes of the SEM.**
- **How an image is obtained**
- **Sample preparation**
- **Materials**
- **Vacuum levels**
- **Information: Imaging, EDS, WDS, EBL, beam current**

# Typical Thermionic Gun

A Thermionic Electron Gun functions in the following manner:

- 1) A positive electrical potential is applied to the anode.
- 2) The filament (cathode) is heated until a stream of electrons is produced.
- 3) The electrons are then accelerated by the positive potential down the column.
- 4) A negative electrical potential ( $\sim 500$  V) is applied to the Wehnelt Cap.
- 5) As the electrons move toward the anode any ones emitted from the filament's side are repelled by the Wehnelt Cap toward the optic axis (horizontal center).



# Typical Thermionic Gun

## Cont'd

6. A collection of electrons occurs in the space between the filament tip and Whenelt Cap. This collection is called a space charge
7. Those electrons at the bottom of the space charge (nearest to the anode) can exit the gun area through the small (<1 mm) hole in the Whenelt Cap
8. These electrons then move down the column to be later used in imaging

*This process insures several things:*

- That the electrons later used for imaging will be emitted from a nearly perfect point source (the space charge)
- The electrons later used for imaging will all have similar energies (monochromatic)
- Only electrons nearly parallel to the optic axis will be allowed out of the gun area

# Other Bulk Specimen Interactions

## X-rays

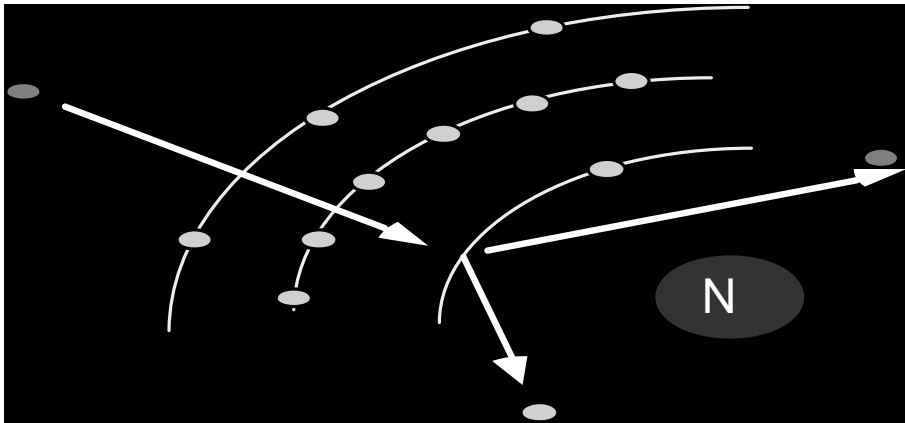
### *Source*

Caused by the de-energization of the specimen atom after a secondary electron is produced. Since a lower (usually K-shell) electron was emitted from the atom during the secondary electron process an inner (lower energy) shell now has a vacancy. A higher energy electron can "fall" into the lower energy shell, filling the vacancy. As the electron "falls" it emits energy, usually X-rays to balance the total energy of the atom.

### *Utilization*

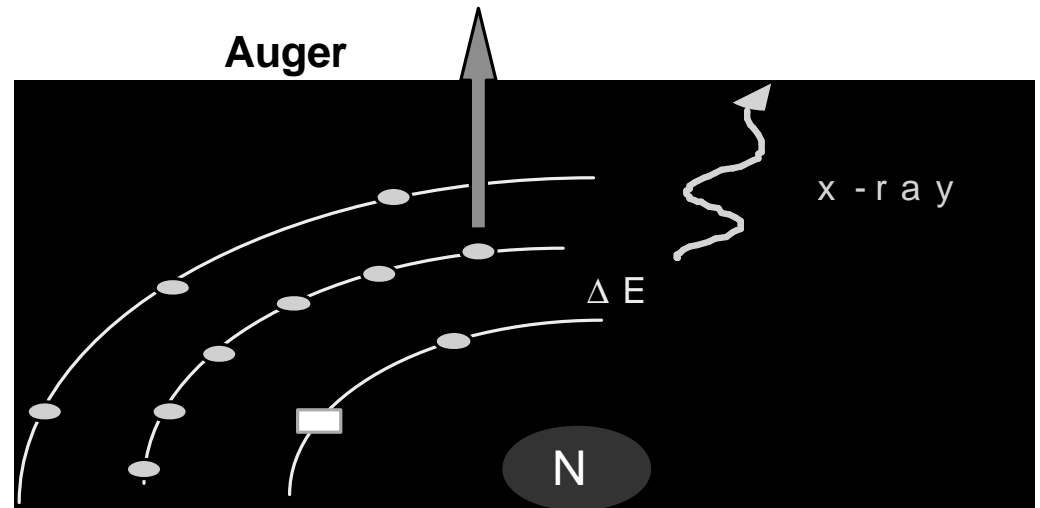
X-rays or Light emitted from the atom will have a characteristic energy which is unique to the element from which it originated. These signals are collected and sorted according to energy to yield micrometer diameter) of bulk specimens limiting the point-to-point comparisons available (see Interaction Volume).

# Characteristic X-rays and Auger Electrons



We've already discussed how an electron collision can create an x-ray that is characteristic of the sample.

This same process also produces Auger electrons



Generally, in the SEM we are only concerned with the characteristic x-ray signal.

# Other Bulk Specimen Interactions

## Auger Electrons

### *Source*

Caused by the de-energization of the specimen atom after a secondary electron is produced. Since a lower (usually K-shell) electron was emitted from the atom during the secondary electron process, an inner (lower energy) shell now has a vacancy.

*A higher energy electron from the same atom can "fall" to a lower energy, filling the vacancy. This creates an energy surplus in the atom which can be corrected by emitting an outer (lower energy) electron; an Auger Electron.*

### *Utilization*

Auger Electrons have a characteristic energy, unique to each element from which it was emitted from. These electrons are collected and sorted according to energy to give compositional information about the specimen. Since Auger Electrons have relatively low energy they are only emitted from the bulk specimen from a depth of <3 nm (see Interaction volume).



# Thin Specimen Interactions

## Unscattered Electrons

### *Source*

- Incident electrons which are transmitted through the thin specimen without any interaction occurring inside the specimen.

### *Utilization*

- The transmission of unscattered electrons is inversely proportional to the specimen thickness.
- Areas of the specimen that are thicker will have fewer transmitted unscattered electrons and so will appear darker.
- Conversely the thinner areas will have more transmitted and thus will appear lighter.

# Thin Specimen Interactions

## Elastically Scattered electrons

### *Source*

Incident electrons that are scattered (deflected from their original path) by atoms in the specimen in an elastic fashion (no loss of energy). These scattered electrons are then transmitted through the remaining portions of the specimen.

### *Utilization*

All electrons follow Bragg's Law and thus are scattered according to  $\text{Wavelength} = 2 \times \text{Space between the atoms in the specimen} \times \sin(\text{angle of scattering})$ . All incident electrons have the same energy (thus wavelength) and enter the specimen normal to its surface. All incidents that are scattered by the same atomic spacing will be scattered by the same angle. These "similar angle" scattered electrons can be collated using magnetic lenses to form a pattern of spots; each spot corresponding to a specific atomic spacing (a plane). This pattern can then yield information about the orientation, atomic arrangements and phases present in the area being examined.

# Thin Specimen Interactions

## Inelastically Scattered Electrons

### *Source*

Incident electrons that interact with specimen atoms in an inelastic fashion, losing energy during the interaction. These electrons are then transmitted through the rest of the specimen

### *Utilization*

Inelastically scattered electrons can be utilized two ways

- **Electron Energy Loss Spectroscopy:** The inelastic loss of energy by the incident electrons is characteristic of the elements that were interacted with. These energies are unique to each bonding state of each element and thus can be used to extract both compositional and bonding (i.e. oxidation state) information on the specimen region being examined.
- **Kikuchi Bands:** Bands of alternating light and dark lines that are formed by inelastic scattering interactions that are related to atomic spacings in the specimen. These bands can be either measured (their width is inversely proportional to atomic spacing) or "followed" like a roadmap to the "real" elasticity scattered electron pattern.

# Characteristic X-rays

- 1) Core electron is knocked out
- 2) Outer shell electron drops down
- 3) Energy difference is emitted as an X-Ray
- 4) Electron shells have discrete energy levels
- 5) The X-Ray energy is characteristic of the particular element in question (energy levels)

**X-Ray Detection: X-Ray intensity is  $\mu Z$**

• Two types of detectors:

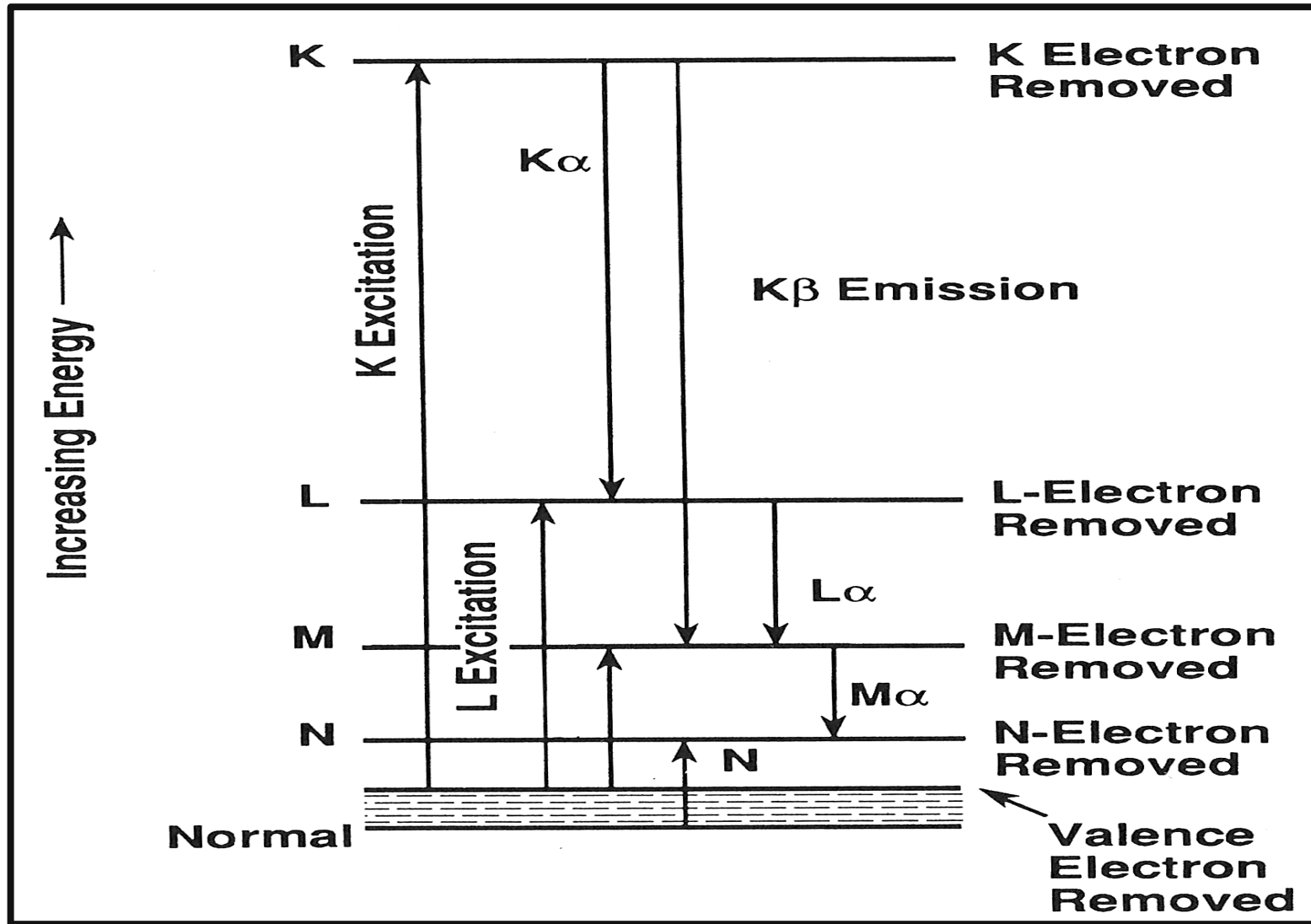
a) Wavelength dispersive:

- Separates X-Rays by wavelength

b) Energy Dispersive (EDXA, EDS, EDX):

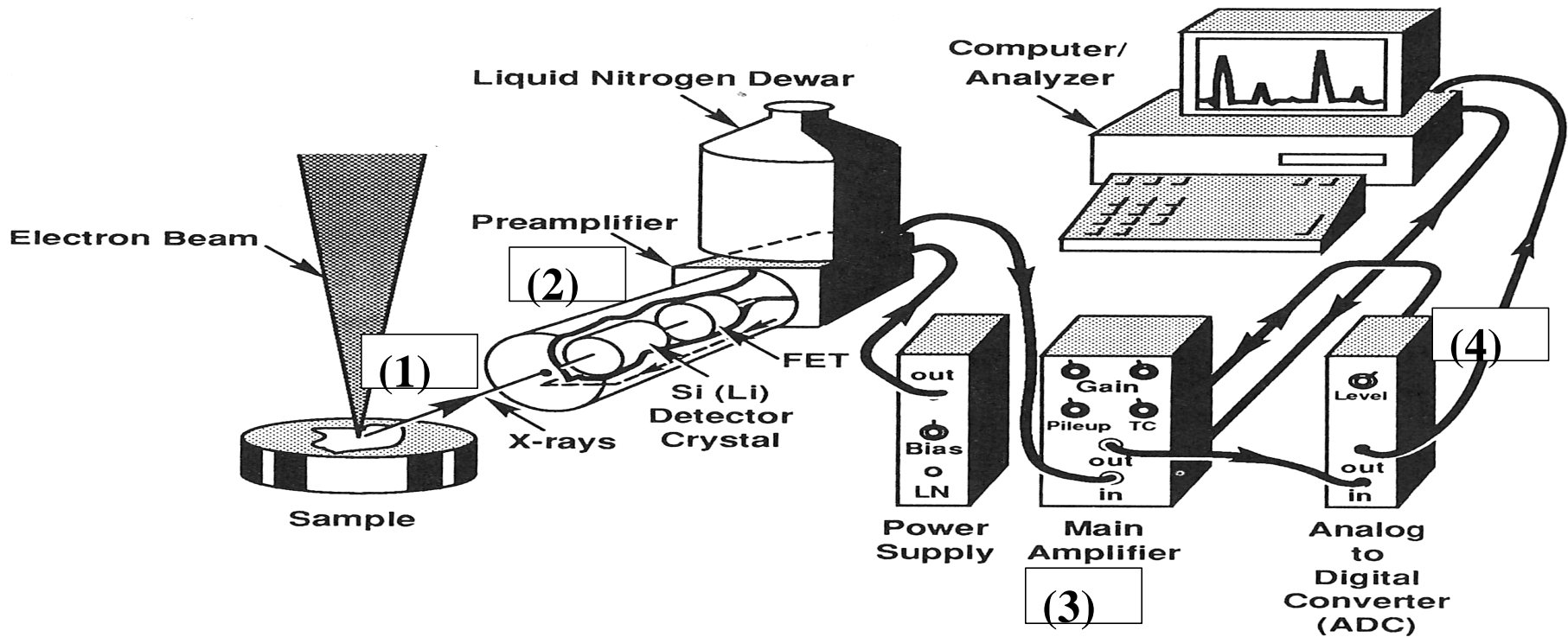
- Separates X-Rays by energy
- Can monitor one energy

# Characteristic X-rays



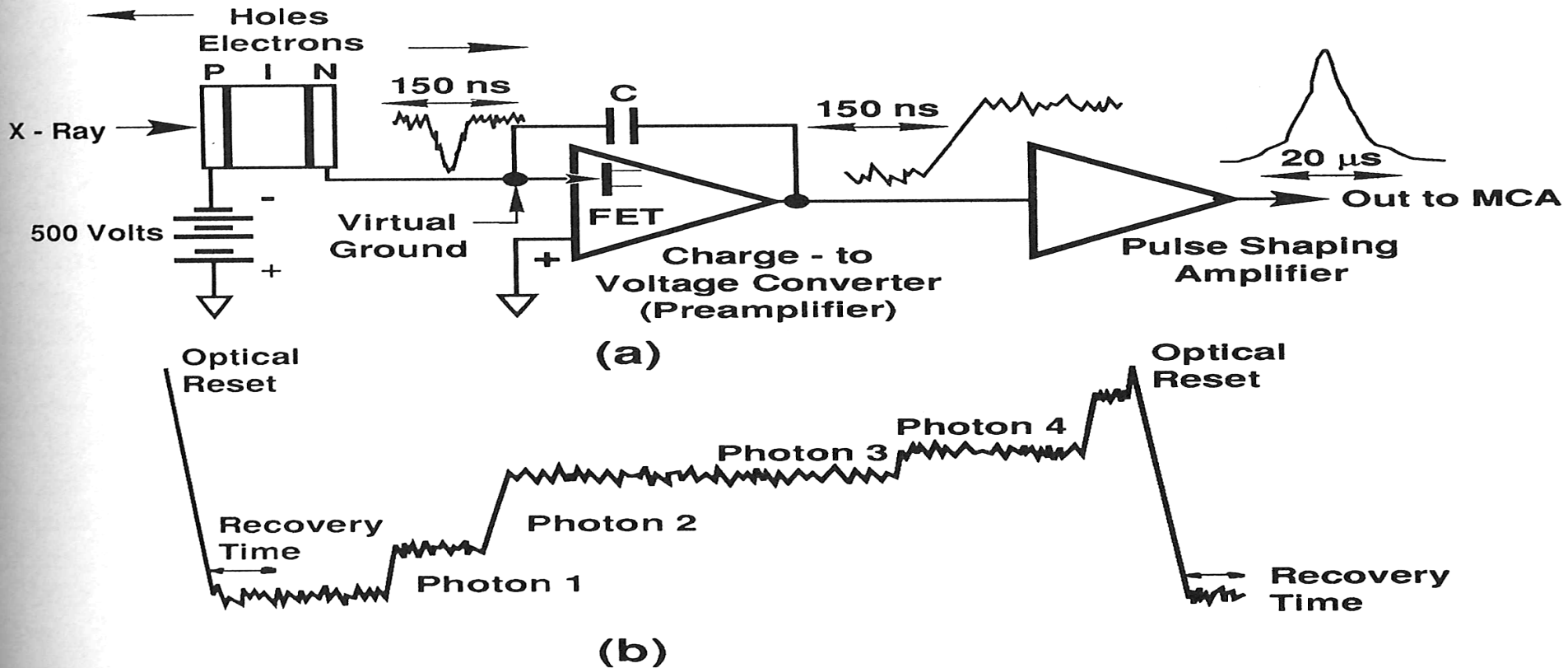
Energy-level diagram for an atom, illustrating the excitation of the K, L, M, and N shells and the formation of  $K_{\alpha}$ ,  $K_{\beta}$ ,  $L_{\alpha}$ , and  $M_{\alpha}$  x-rays.

# X-ray Detection I



1. X-rays enter the detector and strike a Si crystal, creating  $e^- / \text{\AA}$  pairs. The number of  $e^- / \text{\AA}$  pairs is related to the energy of the incoming x-ray.
2. A potential is placed across the crystal to cause the  $e^-$  to move to the + side and the  $\text{\AA}$  to move to the - side, producing a charge pulse.
3. The charge pulse is amplified and sorted. The size of the pulse is again a function of the energy of the incoming x-ray.
4. The number pulses are plotted as a function of pulse size (which = energy of the x-ray).

# Charge to Voltage Conversion Process



(a) Representation of the detector charge-to-voltage converter, and pulse-shaping linear amplifier from an electronic perspective. (b) Output of the charge-to-voltage converter after the detection of a series of x-ray photons.