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# The Laplacian and Mean and Extreme Values

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**Abstract.** The Laplace operator is pervasive in many important mathematical models, and fundamental results such as the Mean Value Theorem for harmonic functions, and the Maximum Principle for super-harmonic functions are well-known. Less well-known is how the Laplacian and its powers appear naturally in a series expansion of the mean value of a function on a ball or sphere. This result is proven here using Taylor's Theorem and explicit values for integrals of monomials on balls and spheres. This result allows for non-standard proofs of the Mean Value Theorem and the Maximum Principle. Connections are also made with the discrete Laplacian arising from finite difference discretization.

**1. THE SERIES EXPANSION OF THE MEAN VALUE.** The Laplace operator (or Laplacian) in  $d$  dimensions,

$$\Delta = \frac{\partial^2}{\partial x_1^2} + \cdots + \frac{\partial^2}{\partial x_d^2}, \quad (1)$$

appears in a variety of important mathematical models, such as the heat/diffusion equation, the wave equation, the Schrödinger equation and some forms of the Navier-Stokes equations. The Mean Value Theorem for harmonic functions ( $\Delta u = 0$ ) on a ball or sphere, and the Maximum and Minimum Principles for super-harmonic ( $\Delta u \geq 0$ ) and sub-harmonic ( $\Delta u \leq 0$ ) functions on bounded domains, are well-known. It is perhaps less well-known that the Laplacian and its powers appear very naturally when the mean value of  $u$  is expanded in terms of a series with respect to the radius of the ball (or sphere). Our key result, Theorem 3, can be found, for example, in [2] for balls and spheres in  $\mathbb{R}^3$ . The expansion given in Theorem 3 is sometimes referred to as a Pizzetti series (cf. [1]), in reference to its earliest known occurrence [4]. The proof given in [2] uses Green's second identity, the fundamental solution for the Laplacian, and a clever recursion. In this section we provide a more straightforward proof of the result in  $\mathbb{R}^d$ , employing Taylor's Theorem and formulas for integrals of monomials on balls and spheres.

**Definition (Multi-index Notation and Taylor Polynomials).** Let  $\mathbb{N}_0 = \mathbb{N} \cup \{0\}$ . For a multi-index  $\alpha = (\alpha_1, \dots, \alpha_d) \in [\mathbb{N}_0]^d$ , we define

$$|\alpha| = \alpha_1 + \cdots + \alpha_d, \quad \alpha! = \alpha_1! \cdots \alpha_d!. \quad (2)$$

The  $\alpha$  partial derivative of a function  $u$ ,  $D^\alpha u$ , and the monomial  $(x - \bar{x})^\alpha$  for fixed  $\bar{x} = (\bar{x}_1, \dots, \bar{x}_d)$  and variable  $x = (x_1, \dots, x_d)$ , are given by

$$D^\alpha u = \frac{\partial^{|\alpha|} u}{\partial x_1^{\alpha_1} \cdots \partial x_d^{\alpha_d}}, \quad (x - \bar{x})^\alpha = (x_1 - \bar{x}_1)^{\alpha_1} \cdots (x_d - \bar{x}_d)^{\alpha_d}. \quad (3)$$

Assuming that  $D^\alpha u(\bar{x})$  is defined for all multi-indices  $|\alpha| \leq m$ , the  $m^{\text{th}}$ -order Taylor Polynomial of  $u$ , centered at  $\bar{x}$ , is

$$T_m(x) = \sum_{|\alpha| \leq m} \frac{D^\alpha u(\bar{x})}{\alpha!} (x - \bar{x})^\alpha. \quad (4)$$

**Theorem 1 (Taylor's Theorem).** Suppose  $\Omega \subset \mathbb{R}^d$  is open and convex, and  $x, \bar{x} \in \Omega$ . If  $u \in C^m(\Omega)$ , then

$$u(x) = T_m(x) + o(|x - \bar{x}|^m) \text{ as } |x - \bar{x}| \rightarrow 0, \quad (5)$$

where  $T_m$  is given by (4).

Here, and elsewhere, we overload the notation  $|\cdot|$ , using it to represent the order of a multi-index, the Euclidean norm of a point/vector in  $\mathbb{R}^d$ , and the volume of a ball,  $|B|$ , or surface area of a sphere,  $|S|$ .

**Definition (Balls and Spheres).** Given a point  $\bar{x} \in \mathbb{R}^d$  and a radius  $h > 0$ , we have the ball and sphere

$$B = B(\bar{x}, h) = \{x \in \mathbb{R}^d : |x - \bar{x}| < h\}, \quad (6)$$

$$S = S(\bar{x}, h) = \{x \in \mathbb{R}^d : |x - \bar{x}| = h\}. \quad (7)$$

An elegant proof of the following result concerning integration of monomials on balls and spheres is given in [3].

**Theorem 2 (Integration of Monomials).** It holds that

$$\int_B (x - \bar{x})^\alpha dx = \begin{cases} \frac{2\Gamma(\beta_1)\cdots\Gamma(\beta_d)}{\Gamma(\beta_1+\cdots+\beta_d)} \frac{h^{|\alpha|+d}}{|\alpha|+d} & \text{if all } \alpha_j \text{ are even,} \\ 0 & \text{otherwise,} \end{cases} \quad (8)$$

$$\int_S (x - \bar{x})^\alpha ds = \begin{cases} \frac{2\Gamma(\beta_1)\cdots\Gamma(\beta_d)}{\Gamma(\beta_1+\cdots+\beta_d)} h^{|\alpha|+d-1} & \text{if all } \alpha_j \text{ are even,} \\ 0 & \text{otherwise,} \end{cases} \quad (9)$$

where  $\beta_j = (\alpha_j + 1)/2$ , and  $\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt$  is the Gamma-function.

From the case  $\alpha = (0, \dots, 0)$ , we quickly obtain the well-known formulas for the volume of  $B$  and the surface area of  $S$ ,

$$|B| = \frac{2\pi^{d/2}}{d\Gamma(d/2)} h^d, \quad |S| = \frac{2\pi^{d/2}}{\Gamma(d/2)} h^{d-1}. \quad (10)$$

Let  $\alpha = 2\sigma$ , so all components are even, and take  $|\alpha| = 2k$ . After some simplification, using  $\Gamma(z+1) = z\Gamma(z)$ ,  $\Gamma(1/2) = \sqrt{\pi}$  and (10), we see that

$$\frac{1}{\alpha!} \int_B (x - \bar{x})^\alpha dx = \frac{|B| h^{2k}}{2^k \sigma! \prod_{j=1}^k (d+2j)}, \quad (11)$$

$$\frac{1}{\alpha!} \int_S (x - \bar{x})^\alpha ds = \frac{|S| h^{2k}}{2^k \sigma! \prod_{j=0}^{k-1} (d+2j)}. \quad (12)$$

We are now ready to state our main result concerning the series expansion of the mean values of  $u$  on  $B$  and  $S$ .

**Theorem 3 (Series Expansion of Mean Values).** Suppose that  $u \in C^{2p}(\Omega)$  for some open set  $\Omega \subset \mathbb{R}^d$  containing  $\bar{B}$ . It holds that

$$u_{ave}(B) = \frac{1}{|B|} \int_B u(x) dx = \sum_{k=0}^p \frac{\Delta^k u(\bar{x})}{2^k k! \prod_{j=1}^k (d+2j)} h^{2k} + o(h^{2p}), \quad (13)$$

$$u_{ave}(S) = \frac{1}{|S|} \int_S u(x) ds = \sum_{k=0}^p \frac{\Delta^k u(\bar{x})}{2^k k! \prod_{j=0}^{k-1} (d+2j)} h^{2k} + o(h^{2p}), \quad (14)$$

as  $h \rightarrow 0$ .

*Proof.* Since the proof for  $S$  is essentially identical, we only prove the result for  $B$ . In particular, we see that

$$\int_B T_{2p}(x) dx = \sum_{k=0}^p \frac{|B| \Delta^k u(\bar{x})}{2^k k! \prod_{j=1}^k (d+2j)} h^{2k}.$$

Here we have used three basic results: only multi-indices of the form  $\alpha = 2\sigma$  survive the integration, the identity (11), and the fact that

$$\Delta^k u = \left( \frac{\partial^2}{\partial x_1^2} + \cdots + \frac{\partial^2}{\partial x_d^2} \right)^k u = \sum_{\substack{|\alpha|=2k \\ \alpha=2\sigma}} \frac{k!}{\sigma!} D^\alpha u, \quad (15)$$

which follows from the multinomial expansion of  $\Delta^k$ . Noting that

$$\int_B u dx = \int_B T_{2p} dx + |B| o(h^{2p}),$$

and dividing by  $|B|$ , completes the proof. ■

For convenience, we give the first three terms of the series for both mean values.

$$u_{ave}(B) = u(\bar{x}) + \frac{\Delta u(\bar{x})}{2(d+2)} h^2 + \frac{\Delta^2 u(\bar{x})}{8(d+2)(d+4)} h^4 + \cdots, \quad (16)$$

$$u_{ave}(S) = u(\bar{x}) + \frac{\Delta u(\bar{x})}{2d} h^2 + \frac{\Delta^2 u(\bar{x})}{8d(d+2)} h^4 + \cdots. \quad (17)$$

The series expansions of the mean values can be used to provide a simple proof of the Mean Value Theorem for harmonic functions. More specifically, if  $\Delta u = 0$  in some open set  $\Omega \supset \bar{B}$ , then  $\Delta^k u(\bar{x}) = 0$  for all  $k > 0$  (we note that  $u$  is analytic on  $\Omega$ ). So  $u_{ave}(B) = u_{ave}(S) = u(\bar{x})$ . We have proved

**Theorem 4 (Mean Value Theorem for Harmonic Functions).** *If  $\Delta u = 0$  in some open set  $\Omega \supset \bar{B}$ , then  $u_{ave}(B) = u_{ave}(S) = u(\bar{x})$ .*

**Remark.** Re-expressing the results of Theorem 3 in the case  $p = 1$ , we see that that

$$\Delta u(\bar{x}) = \lim_{h \rightarrow 0} 2(d+2) \frac{u_{ave}(B) - u(\bar{x})}{h^2} = \lim_{h \rightarrow 0} 2d \frac{u_{ave}(S) - u(\bar{x})}{h^2}. \quad (18)$$

It follows simply from this that, for  $u \in C^2$  in some open set  $\Omega$ , if the average value of  $u$  on any ball (sphere) in  $\Omega$  is equal to the value of  $u$  at the center of the ball (sphere), then  $u$  is harmonic in  $\Omega$ .

**2. THE MAXIMUM AND MINIMUM PRINCIPLES.** We provide a proof of the Maximum Principle that applies Theorem 3 in the case  $p = 1$  at a key point in the argument.

**Theorem 5 (Maximum and Minimum Principles).** *Let  $\Omega \subset \mathbb{R}^d$  be a bounded domain with boundary  $\partial\Omega$ , and let  $u \in C^2(\Omega) \cap C(\overline{\Omega})$ .*

1. *If  $\Delta u \geq 0$  in  $\Omega$ , then  $\max_{x \in \overline{\Omega}} u(x) = \max_{x \in \partial\Omega} u(x)$ .*
2. *If  $\Delta u \leq 0$  in  $\Omega$ , then  $\min_{x \in \overline{\Omega}} u(x) = \min_{x \in \partial\Omega} u(x)$ .*

*Proof.* We prove only the first of these, by contradiction. Let  $M_0 = \max_{x \in \overline{\Omega}} u(x)$  and  $M = \max_{x \in \partial\Omega} u(x)$ . Suppose that  $M < M_0$ , and let  $x_0 \in \Omega$  be a point for which  $u(x_0) = M_0$ . Choose  $R > 0$  large enough so that  $B(x_0, R) \supset \overline{\Omega}$ , and define

$$v = u + \frac{M_0 - M}{2dR^2} |x - x_0|^2. \quad (19)$$

Note that  $\Delta v = \Delta u + (M_0 - M)/R^2 > 0$ ; there is no need to introduce  $v$  in the argument if we assume that  $\Delta u > 0$  in  $\Omega$ . We also see that  $v(x_0) = u(x_0) = M_0$  and

$$v(x) \leq u(x) + \frac{M_0 - M}{2d} \leq M + \frac{M_0 - M}{2d} < M_0 \text{ for } x \in \partial\Omega, \quad (20)$$

so  $v$  must attain its maximum over  $\overline{\Omega}$  at some interior point  $x_1 \in \Omega$ . Since  $\Delta v(x_1) > 0$ , employing Theorem 3 for  $p = 1$  and  $h > 0$  sufficiently small, we see that  $v_{ave}(B) > v(x_1)$  for  $B = B(x_1, h)$ . But this contradicts the fact that  $v(x_1)$  is the maximal value of  $v$ ; the mean value on any ball cannot exceed the maximal value within that ball. ■

**3. CONNECTIONS WITH THE DISCRETE LAPLACIAN.** Recalling (18), we here consider the case in which we fix some (small)  $h$  instead of taking the limit, and replace the true average on the sphere with a discrete average at the points on the sphere that intersect the cartesian axes. Letting  $e_i$  denote the standard coordinate vectors in  $\mathbb{R}^d$ , a simple consequence of Taylor's Theorem in  $\mathbb{R}$  is that

$$u(\bar{x} - he_i) + u(\bar{x} + he_i) = 2u(\bar{x}) + \frac{\partial^2 u(\bar{x})}{\partial x_i^2} h^2 + o(h^2), \quad (21)$$

as  $h \rightarrow 0$ , for any  $u \in C^2$  in a neighborhood of  $\bar{x}$ . Summing both sides for  $1 \leq i \leq d$ , and then dividing by the number of "boundary points", namely  $2d$ , we obtain

$$\frac{1}{2d} \sum_{i=1}^d [u(\bar{x} - he_i) + u(\bar{x} + he_i)] = u(\bar{x}) + \frac{h^2}{2d} \Delta u(\bar{x}) + o(h^2), \quad (22)$$

which is obviously a "discrete" version of (14) when  $p = 1$ , in which the true average is replaced by the standard discrete average of these  $2d$  values.

The (finite difference) discrete Laplacian is defined by

$$\Delta_h u(\bar{x}) = \frac{1}{h^2} \sum_{i=1}^d [u(\bar{x} - he_i) - 2u(\bar{x}) + u(\bar{x} + he_i)], \quad (23)$$

and Taylor's Theorem guarantees that  $\Delta_h u(\bar{x}) = \Delta u(\bar{x}) + o(1)$  as  $h \rightarrow 0$ . Denoting the lefthand-side of (22) by  $\hat{u}_{ave}(S)$ , we have

$$\hat{u}_{ave}(S) = u(\bar{x}) + \frac{h^2}{2d} \Delta_h u(\bar{x}), \quad (24)$$

a result that is commonly used to prove a Discrete Maximum Principle in the analysis of finite difference methods.

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