## Appendix A: The Thermodynamics of Oxygen in Silicon

The formation of oxide precipitates within a silicon crystal lattice has been treated in general; however it is instructive to consider actual quantitative data. For convenience, an oxide precipitate is assumed to be a small, spherical particle of $\mathrm{SiO}_{2}$ embedded within an otherwise perfect silicon crystal. Of course, the standard formation reaction for silicon dioxide is as follows:

$$
\mathrm{Si}(s)+\mathrm{O}_{2}(g) \rightarrow \mathrm{SiO}_{2}
$$

As a matter of chemistry, a standard formation reaction defines a process for which one mole of some material (in this case, $\mathrm{SiO}_{2}$ ) is formed from corresponding elements in standard state. (Here, this is crystalline silicon and oxygen gas.) Extensive tabulations of thermodynamic data for formation reactions have been compiled and are summarized for silicon, oxygen, and silicon dioxide in the following table:

Thermodynamic Potentials of Silicon, Oxygen, and Silicon Dioxide:

|  | $\Delta H_{f}^{\circ *}$ | $\Delta G_{f}^{\circ *}$ | $S^{\circ * *}$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{Si}(s)$ | 0 | 0 | $\begin{gathered} 18.81 \\ 4.50 \end{gathered}$ |
| $\mathrm{O}_{2}(\mathrm{~g})$ | 0 | 0 | $\begin{gathered} 205.152 \\ 49.03 \end{gathered}$ |
| $\mathrm{SiO}_{2}$ (quartz) | $\begin{aligned} & -910.7 \\ & -217.7 \end{aligned}$ | $\begin{aligned} & -856.3 \\ & -204.7 \end{aligned}$ | $\begin{gathered} 41.46 \\ 9.909 \end{gathered}$ |
| $\mathrm{SiO}_{2}(\text { cristobalite })^{\dagger}$ | $\begin{aligned} & -909.5 \\ & -217.37 \end{aligned}$ | $\begin{aligned} & -855.5 \\ & -204.46 \end{aligned}$ | $\begin{array}{r} 42.68 \\ 10.20 \end{array}$ |
| $\mathrm{SiO}_{2}\left(\right.$ tridymite) ${ }^{\dagger}$ | $\begin{aligned} & -909.1 \\ & -217.27 \end{aligned}$ | $\begin{aligned} & -855.3 \\ & -204.42 \end{aligned}$ | $\begin{array}{r} 43.5 \\ 10.4 \end{array}$ |
| $\mathrm{SiO}_{2}\left(\right.$ quartz glass) ${ }^{\dagger}$ | $\begin{aligned} & -903.5 \\ & -215.94 \end{aligned}$ | $\begin{aligned} & -850.7 \\ & -203.33 \end{aligned}$ | $\begin{array}{r} 46.9 \\ 11.2 \end{array}$ |

${ }^{*} \mathrm{~kJ} /$ mole (italics: $\mathrm{kcal} / \mathrm{mole}$ ); ${ }^{* *} \mathrm{~J} /$ mole ${ }^{\circ} \mathrm{K}$ (italics: cal/mole ${ }^{\circ} \mathrm{K}$ )
${ }^{\dagger}$ data taken from Handbook of Chemistry and Physics- $1^{\text {st }}$ Student Ed. (1988), all other data taken from CODATA Key Values for Thermodynamics

Of course, standard conditions are defined as $298.15^{\circ} \mathrm{K}$ and an ambient pressure of one atmosphere. Clearly, the thermodynamic potentials for all forms of silicon dioxide (quartz, cristobalite, tridymite, and glass) are quite similar.

However, within the silicon lattice, oxygen is not in gaseous diatomic form. Therefore, to be applicable to oxide precipitation, the formation reaction must be modified as follows:

$$
\mathrm{Si}(s)+2 \mathrm{O}(i n t) \rightarrow \mathrm{SiO}_{2}
$$

Here, $\mathrm{O}($ int $)$ denotes oxygen atoms occupying interstitial sites within the silicon crystal lattice. Clearly, this reaction and the standard formation reaction are related by a third reaction that represents dissolution of oxygen gas in the silicon lattice and which can be formally written as follows:

$$
\mathrm{O}_{2}(\mathrm{~g}) \rightarrow 2 \mathrm{O}(\text { int })
$$

If the standard free energy of this reaction can be found, then the standard free energy of the previous reaction is easily determined.

For this purpose, it is useful to consider the dissolution of oxygen in solid silicon as a microscopic process. Obviously, oxygen molecules must react with the silicon lattice to form oxygen interstitials. This is represented schematically below:


Of course, this is just an alternative representation of the dissolution reaction appearing above, however, the silicon lattice is included explicitly. Clearly, the overall enthalpy change for this process must include contributions from strain energy associated with an oxygen interstitial, binding energy of an oxygen molecule, and binding energy of an oxygen atom within the silicon crystal. One expects the first two of these contributions to be positive and the last one to be negative. However, with the exception of the binding energy of molecular oxygen, these contributions are not readily determined. In contrast, the entropy change can be represented as the difference of the configurational entropy change due to random distribution of oxygen atoms in interstitial sites, $\Delta S_{\mathrm{O}}^{C}$, and the standard entropy of oxygen gas, $S_{\mathrm{O}_{2}}(T)$ :

$$
\Delta S=\Delta S_{\mathrm{O}}^{C}-\frac{N_{\mathrm{O}}}{2 N_{A}} S_{\mathrm{O}_{2}}(T)
$$

Here, $N_{\mathrm{O}}$ is the number of oxygen interstitials and $N_{A}$ is Avogadro's number. Of course, a standard entropy is also associated with the silicon lattice itself; however, if the lattice is not disrupted by oxygen interstitials, this entropy can be taken to be unchanged when oxygen interstitials are introduced into the lattice and therefore makes no contribution to $\Delta S$. (One should note here that $\Delta S$ is defined as the entropy change associated with the formation of $N_{\mathrm{O}}$ oxygen interstitials.)

Naturally, if $N$ is defined as the number of interstitial sites in the crystal, then the configurational entropy change is easily represented as a binomial coefficient:

$$
\Delta S=k \ln \left(\frac{N!}{\left(N-N_{\mathrm{O}}\right)!N_{\mathrm{O}}!}\right)-\frac{N_{\mathrm{O}}}{2 N_{A}} S_{\mathrm{O}_{2}}(T)
$$

As is usual, one applies Stirling's approximation to obtain:

$$
\Delta S=k\left(N \ln N-\left(N-N_{\mathrm{O}}\right) \ln \left(N-N_{\mathrm{O}}\right)-N_{\mathrm{O}} \ln N_{\mathrm{O}}\right)-\frac{N_{\mathrm{O}}}{2 N_{A}} S_{\mathrm{O}_{2}}(T)
$$

One formally adds and subtracts $N_{\mathrm{O}} \ln N$ to the quantity within the parenthesis, from which it follows that:

$$
\Delta S=-k\left(N_{\mathrm{O}} \ln \frac{N_{\mathrm{O}}}{N}+\left(N-N_{\mathrm{O}}\right) \ln \left(1-\frac{N_{\mathrm{O}}}{N}\right)\right)-\frac{N_{\mathrm{O}}}{2 N_{\mathrm{A}}} S_{\mathrm{O}_{2}}(T)
$$

Clearly, one expects that $N$ will be much larger than $N_{\mathrm{O}}$, hence the second logarithmic term can be ignored, thus:

$$
\Delta S=-k N_{\mathrm{O}} \ln \frac{N_{\mathrm{O}}}{N}-\frac{N_{\mathrm{O}}}{2 N_{A}} S_{\mathrm{O}_{2}}(T)
$$

Conceptually, it is convenient to replace absolute numbers of oxygen interstitials and interstitial sites, $N_{\mathrm{O}}$ and $N$, by corresponding concentrations, $C_{\mathrm{o}}$ and $C$. Furthermore, $\Delta S$ can be recast as a molar quantity if one rescales the right hand side by the ratio, $N_{A} / N_{\mathrm{O}}$. Thus, it follows that:

$$
\Delta S=-k N_{A} \ln \frac{C_{\mathrm{O}}}{C}-\frac{S_{\mathrm{O}_{2}}(T)}{2}
$$

Since there are five interstitial sites per diamond cubic unit cell, it follows that $C$ is just $5 / a^{3}$, such that $a$ is just the lattice parameter for silicon. Therefore, one finds that $C$ is approximately $3.123\left(10^{22}\right) \mathrm{cm}^{-3}$. Furthermore, $k N_{A}$ is the ordinary ideal gas constant, $R$, which has a nominal value of $8.31441 \mathrm{~J} / \mathrm{mole}^{\circ} \mathrm{K}$.

The standard entropy of oxygen gas at any temperature and one atmosphere pressure can be obtained from the standard entropy at $298^{\circ} \mathrm{K}$ by means of the integral formula:

$$
S_{\mathrm{O}_{2}}(T)=\int_{298}^{T} \frac{C_{p}}{T} d T+S_{\mathrm{O}_{2}}\left(298^{\circ} \mathrm{K}\right)
$$

Here, $C_{p}$ is the molar heat capacity at a constant pressure of one atmosphere. If one assumes that oxygen is an ideal diatomic gas, then $C_{p}$ has the value of $7 R / 2$, hence it follows that:

$$
S_{\mathrm{O}_{2}}(T)=\frac{7 R}{2}(\ln T-\ln 298)+S_{\mathrm{O}_{2}}\left(298^{\circ} \mathrm{K}\right)
$$

Alternatively, $S_{\mathrm{O}_{2}}(T)$ can be determined more accurately from published curve fits for the temperature dependence of constant pressure heat capacity, $C_{p}(T)$.

$$
C_{p}(T)=a+b T+c T^{2}+\frac{d}{T^{2}}
$$

Here, $a, b, c$, and $d$, are empirical coefficients. Thus, one obtains:

$$
S_{\mathrm{O}_{2}}(T)=\int_{298}^{T}\left(\frac{a}{T}+b+c T+\frac{d}{T^{3}}\right) d T+S_{\mathrm{O}_{2}}\left(298^{\circ} \mathrm{K}\right)
$$

For convenience, an aggregate coefficient, $B$, can be defined in terms of $a, b, c, d$, and the standard entropy of oxygen gas:

$$
B=a \ln 298+298 b+(298)^{2} \frac{c}{2}-\frac{d}{2(298)^{2}}-S_{\mathrm{O}_{2}}\left(298^{\circ} \mathrm{K}\right)
$$

Hence, $S_{\mathrm{O}_{2}}(T)$ has the following form:

$$
S_{\mathrm{O}_{2}}(T)=a \ln T+b T+\frac{c T^{2}}{2}-\frac{d}{2 T^{2}}-B
$$

For oxygen gas, published values for $a, b, c$, and $d$ are $34.602 \mathrm{~J} / \mathrm{mole}{ }^{\circ} \mathrm{K}, 1.0795\left(10^{-3}\right)$ $\mathrm{J} / \mathrm{mole}^{\circ} \mathrm{K}^{2}, 0 \mathrm{~J} / \mathrm{mole}^{\circ} \mathrm{K}^{3}$, and $-785377 \mathrm{~J}^{\circ} \mathrm{K} / \mathrm{mole}$, respectively. From these values, one finds $B$ equal to $-3.2777 \mathrm{~J} / \mathrm{mole}^{\circ} \mathrm{K}$.

Obviously, it follows from the fundamental definition of Gibbs free energy that for the oxygen dissolution reaction:

$$
\Delta G(T)=\Delta H(T)+R T \ln \frac{C_{\mathrm{O}}}{C}+\frac{T S_{\mathrm{O}_{2}}(T)}{2}
$$

Obviously, $\Delta H$ is unknown. However, an method for the determination of $\Delta H$ is afforded by the oxygen solubility equilibrium. Of course, $\Delta G$ vanishes if dissolved oxygen is in equilibrium with ambient oxygen gas. Therefore, it follows that:

$$
\Delta H(T)=-R T \ln \frac{C_{\mathrm{O}}^{s a t}(T)}{C}-\frac{T S_{\mathrm{O}_{2}}(T)}{2}
$$

Here, $C_{\mathrm{O}}^{\text {sat }}(T)$ is saturated interstitial oxygen concentration at an absolute temperature, $T$. As shown in the following figure, this quantity has been determined experimentally over the temperature range $1000-1300^{\circ} \mathrm{C}$ :


Fig. A1: Interstital oxygen solubility as a function of temperature
Clearly, the saturated oxygen interstitial concentration varies from $10^{17}$ to $10^{18} \mathrm{~cm}^{-3}$ over the given temperature range. This is consistent with typical oxygen concentrations observed in CZ silicon wafers. Furthermore, it seems clear from the trend, that the solid solubility of oxygen in silicon should further decrease when extrapolated to lower temperatures. This effect is likely the result of an increased strain energy contribution to enthalpy due to increased lattice rigidity at lower temperatures. Accordingly, if these concentrations are used to determine $\Delta H(T)$, one finds that resulting values are negative, but relatively small. Of course, negative values imply that energy is released when oxygen dissolves in a silicon crystal. This can be rationalized if one considers experimentally measured binding energies. In particular, $\mathrm{Si}-\mathrm{Si}$ and $\mathrm{O}-\mathrm{O}$ binding energies are observed to be $326.8 \mathrm{~kJ} / \mathrm{mole}$ ( $78.1 \mathrm{kcal} /$ mole) and $498.34 \mathrm{~kJ} / \mathrm{mole}$ ( 119.106 $\mathrm{kcal} /$ mole), respectively. These can be compared to the Si-O binding energy, which is found to be $798.7 \mathrm{~kJ} / \mathrm{mole}$ ( $190.9 \mathrm{kcal} /$ mole). Clearly, the formation of $\mathrm{Si}-\mathrm{O}$ bonds from $\mathrm{Si}-\mathrm{Si}$ and $\mathrm{O}-\mathrm{O}$ bonds is strongly exothermic. (This is also clear just from the large, negative formation enthalpy of silicon dioxide.) However, lattice strain largely offsets this so that the magnitude of the enthalpy of formation for oxygen interstitials is fairly
small. Calculated values for the enthalpy of formation for oxygen interstitials is given in the following figure:


Fig. A2: Enthalpy of formation of interstitial oxygen as a function of temperature (Heavy line: $C_{p}$ obtained from empirical curve fit; Light line: $C_{p}$ taken as $7 R / 2$ )

Here, $C_{p}$ has been estimated both from an empirical curve fit (heavy line) or as a constant, $7 R / 2$ (light line). The difference is found to be only about $3 \mathrm{~kJ} / \mathrm{mole}$ and a simple linear fit is quite sufficient to describe the temperature dependence of both results. Hence, $\Delta H(T)$ can be represented by the empirical linear expression:

$$
\Delta H(T)=c_{p} T+\Delta H_{0}
$$

From the curve fit data, the parameters, $c_{p}$ and $\Delta H_{0}$ are found to be $-0.12459 \mathrm{~kJ} / \mathrm{mole}{ }^{\circ} \mathrm{K}$ and $124.87 \mathrm{~kJ} / \mathrm{mole}$, respectively. Similarly, for $C_{p}$ taken as $7 R / 2, c_{p}$ and $\Delta H_{0}$ are found to be $-0.11873 \mathrm{~kJ} / \mathrm{mole}^{\circ} \mathrm{K}$ and $120.52 \mathrm{~kJ} /$ mole, respectively. This expression may be substituted into the expression for $\Delta G$ to obtain the empirical result:

$$
\Delta G(T)=\Delta H_{0}+T\left(c_{p}+R \ln \frac{C_{\mathrm{O}}}{C}+\frac{S_{\mathrm{O}_{2}}(T)}{2}\right)
$$

This is the Gibbs free energy of formation per mole of oxygen interstitials for an elemental silicon crystal having an oxygen interstitial concentration of $C_{0}$.

Conventional enthalpy of formation of $\mathrm{SiO}_{2}$ as a function of temperature is readily obtained by integrating over the heat capacity for $\mathrm{SiO}_{2}$ as follows:

$$
\Delta H_{f_{\mathrm{SiO}_{2}}}(T)=\int_{298}^{T} C_{p}(T) d T+\Delta H_{f_{\mathrm{SiO}_{2}}}\left(298^{\circ} \mathrm{K}\right)
$$

Again, as in the case of elemental oxygen, $C_{p}(T)$ for $\mathrm{SiO}_{2}$ as is expressed as an empirical curve fit. This result is then combined with $\Delta H(T)$ for oxygen dissolution to obtain the quantity, $\Delta H_{\mathrm{SiO}_{2}}$, hence:

$$
\Delta H_{\mathrm{SiO}_{2}}(T)=\Delta H_{f_{\mathrm{SiO}_{2}}}(T)-2\left(c_{p} T+\Delta H_{0}\right)
$$

This is represented by the figure:


Fig. A3: Enthalpy of formation of bulk silicon dioxide in silicon as a function of temperature

Here, $\Delta H_{\mathrm{SiO}_{2}}$, defined as enthalpy of formation of bulk silicon dioxide from elemental silicon and oxygen interstitials (ignoring surface and strain energies associated with precipitates), corresponds to the heavy plot. This is contrasted with ordinary enthalpy of formation of $\mathrm{SiO}_{2}$ (corresponding to the light plot). Clearly, over the temperature range $1000-1300^{\circ} \mathrm{K}, \Delta H_{\mathrm{SiO}_{2}}$ varies only by about $120 \mathrm{~kJ} / \mathrm{mole}$.

Of course, entropies for silicon, oxygen, and silicon dioxide can be determined as a function of temperature in an entirely analogous fashion. These quantities can be combined with the enthalpy of formation obtained previously to obtain the conventional Gibbs free energy of formation of $\mathrm{SiO}_{2}$ as a function of temperature. Naturally, the resulting Gibbs free energy of formation is then combined with $\Delta G(T)$ for oxygen dissolution to obtain the quantity, $\Delta G_{\mathrm{SiO}_{2}}$, hence:

$$
\Delta G_{\mathrm{SiO}_{2}}(T)=\Delta G_{f_{\mathrm{SiO}_{2}}}(T)-2 \Delta H_{0}-T\left(2 c_{p}+2 R \ln \frac{C_{\mathrm{O}}}{C}+S_{\mathrm{O}_{2}}(T)\right)
$$

As before, this relation is represented figuratively as follows:


Fig. A4: Gibbs free energy of formation of bulk silicon dioxide in silicon as a function of temperature black: $C_{0}=10^{14} \mathrm{~cm}^{-3}$; blue: $C_{0}=10^{15} \mathrm{~cm}^{-3}$; red: $C_{0}=10^{16} \mathrm{~cm}^{-3}$; green: $C_{0}=10^{17} \mathrm{~cm}^{-3}$; brown: $C_{0}=10^{18} \mathrm{~cm}^{-3}$

Of course, $\Delta G_{\mathrm{SiO}_{2}}$ is Gibbs free energy of formation of bulk silicon dioxide from elemental silicon and oxygen interstitials (again, ignoring any precipitate surface and strain energies). Clearly, the Gibbs free energy has a much larger temperature variation than enthalpy. Furthermore, unlike enthalpy, Gibbs free energy is a function of interstitial oxygen concentration. This is illustrated by the heavy plots of various colors in the preceding figure. As one would expect, when interstitial oxygen concentration decreases, $\Delta G_{\mathrm{SiO}_{2}}$ becomes more positive (i.e., oxide formation from interstitial oxygen is less favored.) As for enthalpy, the light plot corresponds to the conventional Gibbs free energy of formation of $\mathrm{SiO}_{2}$. Clearly, if at some temperature the ordinary Gibbs free energy of formation and $\Delta G_{\mathrm{SiO}_{2}}$ are equal (i.e., corresponding plots intersect), then the associated concentration of oxygen interstitials, $C_{0}$, can be identified with the solubility limit, i.e., $\Delta G(T)$, as defined previously, exactly vanishes. Furthermore, since elemental oxygen gas is no longer a formal reactant for formation of bulk silicon dioxide from interstitial oxygen, all reactant and product phases can be considered condensed. Therefore, enthalpy, $\Delta H_{\mathrm{SiO}_{2}}$, is equivalent to internal energy, $\Delta E_{\mathrm{SiO}_{2}}$, and, likewise Gibbs free energy, $\Delta G_{\mathrm{SiO}_{2}}$, is equivalent to Helmholtz free energy, $\Delta A_{\mathrm{SiO}_{2}}$. Along with appropriate expressions for surface and strain energies, these quantities can be used to describe oxygen precipitation in silicon.

